VLA and ALMA Imaging of Intense, Galaxy-Wide Star Formation in $z \sim 2$ Galaxies

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We present 12-resolution extinction-independent distributions of star formation and dust in 11 star-forming galaxies (SFGs) at $z = 1.3-3.0$. These galaxies are selected from sensitive blank-field surveys of the $2' \times 2'$ Hubble Ultra-Deep Field at $\lambda = 5$ cm and 1.3 mm using the Karl G. Jansky Very Large Array and Atacama Large Millimeter/submillimeter Array. They have star formation rates (SFRs), stellar masses, and dust properties representative of massive main-sequence SFGs at $z \sim 2$. Morphological classification performed on spatially resolved stellar mass maps indicates a mixture of disk and morphologically disturbed systems; half of the sample harbor X-ray active galactic nuclei (AGNs), thereby representing a diversity of $z \sim 2$ SFGs undergoing vigorous mass assembly. We find that their intense star formation most frequently occurs at the location of stellar-mass concentration and extends over an area comparable to their stellar-mass distribution, with a median diameter of 4.2 $\pm$ 1.8 kpc. This provides direct evidence of galaxy-wide star formation in distant blank-field-selected main-sequence SFGs. The typical galactic-average SFR surface density is $2.5 M_\odot$ yr$^{-1}$ kpc$^{-2}$, sufficiently high to drive outflows. In X-ray-selected AGN where radio emission is enhanced over the level associated with star formation, the radio excess pinpoints the AGNs, which are found to be cospatial with star formation. The median extinction-independent size of main-sequence SFGs is two times larger than those of bright submillimeter galaxies, whose SFRs are 3–8 times larger, providing a constraint on the characteristic SFR ($\sim 300 M_\odot$ yr$^{-1}$) above which a significant population of more compact SFGs appears to emerge.

Key words: galaxies: active – galaxies: evolution – galaxies: high-redshift – galaxies: star formation

1. INTRODUCTION

Numerical simulations and observational inferences suggest that typical star-forming galaxies (SFGs) at the peak of galaxy assembly activity, $z \simeq 1-3$, assembled most of their stellar mass via accretion of cold gas, which led to gas-rich unstable disks and disk-wide star formation (e.g., Kereš et al. 2005; Bournaud & Elmegreen 2009; Dekel et al. 2009). The isolated in situ assembly of typical SFGs is inferred from a relationship between star formation and stellar mass, the “main sequence” (e.g., Noeske et al. 2007; Whitaker et al. 2012; Speagle et al. 2014), and from the rarity of compact starbursts at $z \sim 2$, as indicated by their specific star formation rate (SFR) distribution and infrared color (Elbaz et al. 2011; Rodighiero et al. 2011). At higher SFRs ($>$ 500 $M_\odot$ yr$^{-1}$ at $z \sim 2$), there is no theoretical consensus on whether mergers or continuous accretion are the dominant triggering mechanism of intense star formation (Davé et al. 2010; Hopkins et al. 2010; Hayward et al. 2013; Narayanan et al. 2015); the relative contribution of
in situ versus merger modes at this SFR regime have not been constrained observationally.

Contrary to the inferred in situ assembly of SFGs, Hubble Space Telescope (HST) observations of main-sequence SFGs at \( z \sim 2 \) commonly show galactic substructures and asymmetry that are signposts of mergers that drive intense star formation in the local universe (Lotz et al. 2004; Kartaltepe et al. 2012). However, some substructures, such as optically bright star-forming clumps with SFR \( \simeq 1-30 \) of \( M_\odot \text{yr}^{-1} \) (e.g., Förster Schreiber et al. 2009; Guo et al. 2015) can also be a natural consequence of gas-rich turbulent disks evolving in isolation. The role of star-forming clumps in assembling the bulk of stellar mass is debated, however (Genel et al. 2012; Bournaud et al. 2014), and optically selected clumps altogether contain <10%–20% of the total star formation of their host (Guo et al. 2012, 2015; Wyots et al. 2012). Directly observing the distribution of the bulk of star formation in galaxies at \( z \sim 2 \) is therefore key to establishing the relative contribution between modes of star and bulge formation.

Significant progress in directly imaging the distribution of star formation at \( z > 1 \) has been made with spatially resolved H\( \alpha \) spectroscopy, e.g., SINS (Förster Schreiber et al. 2009), KMOS3D (Wisnioski et al. 2015), and KROSS (Stott et al. 2016). Average H\( \alpha \) maps from 3D-HST, the largest sample thus far of resolved star formation at \( z = 1.5-2.5 \), show that the star formation surface density, \( S_{\text{SFR}} \), on average peaks near the centers of massive galaxies (Nelson et al. 2015). However, the dust extinction also peaks at the center, such that a factor of 6 correction to the inferred H\( \alpha \) SFR is required in the central kpc (and more than a factor of 10 close to the center; Nelson et al. 2016). Above SFRs as low as \( 20 M_\odot \text{yr}^{-1} \), which is 0.2\( \times \) the typical rate for main-sequence SFGs at \( z > 1 \), galaxies become so dust enshrouded that almost no light emerges in rest-frame ultraviolet observations (Reddy et al. 2010). The question of where exactly new stars form within typical \( z \sim 2 \) SFGs is hence a deceptively simple one that is challenging to address.

Breakthroughs in this area require star formation tracers of subarcsecond resolution that are extinction independent for main-sequence SFGs at \( z > 1 \), which are now available with the Karl G. Jansky Very Large Array (VLA) and the Atacama Large Millimeter/submillimeter Array (ALMA). To this end, we conduct two sensitive blank-field imaging surveys of the Hubble Ultra-Deep Field (HUDF, H\( \alpha \), \( \alpha = 03^h32^m, \delta = -27^\circ47' \)) using the VLA and ALMA at \( \lambda = 5 \) cm and 1.3 mm, respectively, to make \( 0^\prime\!4 \) resolution images of SFGs at \( z \sim 2 \). The 5 cm continuum traces star formation through the synchrotron emission from supernova remnants, but can be affected by active galactic nucleus (AGN) emission; whereas the 1.3 mm continuum traces cold dust associated with star formation, but requires uncertain assumptions about the shapes of spectral energy distributions (SEDs) to estimate the SFR. The combination of the two surveys therefore provides complementary strengths, especially in the HUDF where the wealth of ancillary data can help, e.g., identify AGN. By establishing that the VLA and ALMA trace the common extent of star formation, ALMA can serve as a morphological tracer of “pure” star formation in AGN hosts because the 1.3 mm dust continuum is neither contaminated by AGN torus emission (Elvis et al. 1994; Mullaney et al. 2011; Mor & Netzer 2012) nor by synchrotron emission from the jets.\(^{26}\) Furthermore, because ALMA significantly gains in sensitivity to star formation with the help of negative \( K \)-corrections at \( z > 2.5 \), whereas the VLA gradually loses star formation sensitivity beyond this redshift, the combination of VLA and ALMA yields a sensitive probe of the morphology of star formation over the entire range of \( z = 1-3 \).

In this paper, we present first results from combining the VLA and ALMA HUDF surveys, focusing on the extinction-independent distributions of the star formation in SFGs at \( z = 1-3 \). We discuss the VLA and ALMA surveys and ancillary data in Section 2 and present the size and location of star formation in SFGs selected from VLA and ALMA, along with their implications, in Section 3. We adopt a \( \Lambda \)CDM cosmology with \( \Omega_M = 0.3, \Omega_\Lambda = 0.7, H_0 = 70 \text{ km s}^{-1} \text{ Mpc}^{-1} \), and the Chabrier (2003) IMF.

2. OBSERVATIONS AND ANCILLARY DATA

To image SFGs at \( z \sim 2 \) with the VLA, the radio continuum at 5 cm offers a good balance between angular resolution and the expected flux density, given the synchrotron radio spectrum, \( S_{\nu} \propto \nu^{-0.7} \). We observed the HUDF for 177 hr in the A, B, and C configurations during 2014 March–2015 September using the full C-band bandwidth of 4–8 GHz (\( \lambda = 7.5-3.7 \) cm). The observations comprised 42 dynamically scheduled sessions of 2.5–5.5 hr. Each session observed 3C48 for flux and bandpass calibrations; J0402-3147 was observed for phase calibration every 25 minutes. Data reduction was carried out with CASA (McMullin et al. 2007) using the following steps: (1) standard calibration using the VLA Data Reduction Pipeline (C. J. Chandler et al. 2016, in preparation); (2) removal of any portions of the data corrupted by strong radio frequency interference; and (3) imaging with the task Tclean. The imaging parameters are the following: MT-MFS deconvolver withntem of 2, 0.0506 pixel size, and Briggs weighting with a robust parameter of 0.5. We imaged the data well beyond the primary beam radius of 3.6\( ^\prime \) (employing the w-projection) to mitigate the imaging artifacts caused by the sidelobes from bright sources far from the pointing center. The final image has a 0.031 \( \times \) 0.061 synthesized beam and rms noise at the pointing center of 0.32 \( \mu \)Jy beam\(^{-1} \), consistent with the theoretical sensitivity (0.30 \( \mu \)Jy beam\(^{-1} \)).

The ALMA HUDF Deep Field is a 1.3 mm survey of the 4.5 arcmin\(^2 \) HUDF using a 45-pointing mosaic during 2014 July–2015 May in 13 sessions, using a total of approximately 20 hr. The rms noise of the naturally weighted ALMA map is 29 \( \mu \)Jy beam\(^{-1} \) and the synthesized beam is 0.037 \( \times \) 0.048. Each session observed J0334-301 for flux and bandpass calibrations, and also for phase calibration during the 2015 sessions (2014 sessions used J0348-2749 for phase calibration). Calibration was carried out with CASA and imaging with the task CLEAN, adopting natural weighting to maximize sensitivity. Deconvolution was not performed as there are no strong sidelobes from 1.3 mm objects in the HUDF. Details of the ALMA observations, data reduction, and source extraction are given in a companion paper on the ALMA HUDF survey (Dunlop et al. 2016, hereafter, D16).

The following physical parameters in this paper are also estimated by D16: (1) SFR\(_{\text{IR}} \) that combines Spitzer (24 \( \mu \)m), Herschel (deblended 70–500 \( \mu \)m), and ALMA measurements; (2) galaxy-integrated stellar mass via SED fitting; (3) spectroscopic redshift compilation, using redshift measurements, e.g., from the VLT/MUSE IFU survey (R. Bacon et al. 2016, in preparation), and photometric redshifts where spectroscopic

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\(^{26}\) Assuming a flat spectral index for AGN, the radio continuum from the brightest radio AGN in the HUDF is less than 1% of their 1.3 mm emission from star formation.
redshifts are not available. We identify X-ray AGN using the 4 Ms Chandra catalog (Xue et al. 2011) with updated redshifts and estimate the 5 cm radio SFR using the Bell (2003) indicator, assuming a spectral slope $S_\nu \propto \nu^{-0.7}$ for K-correction.

HST images of the HUDF reach 29.5–30.3 mag (AB) at 0.4–1.6 μm (Ellis et al. 2013; Koekemoer et al. 2013), from which we construct stellar-mass maps by fitting spatially resolved SEDs using the procedures described by Cibinel et al. (2015). The sums of the stellar mass in these maps agree within 0.1 dex with the integrated stellar-mass estimates that use longer wavelength photometry (e.g., Spitzer/IRAC 3.6–8.0 μm; Cibinel et al. 2015), suggesting that the maps provide a good description of stellar-mass distribution in obscured SFGs. These stellar-mass maps serve as a reference frame to map where star formation is occurring in relation to stellar-mass buildup within each galaxy. We further classify the stellar-mass morphologies as isolated or disturbed by performing the asymmetry ( Conselice 2003; Zamojski et al. 2007) and $M_{20}$ (Lotz et al. 2004) analysis on these stellar-mass maps, with the $M_{20}$ being the second-order moment of the 20% brightest pixels, following Lotz et al. (2004). This classification method is shown by Cibinel et al. (2015) to be capable of identifying galaxy mergers with a smaller fraction of contamination from clumpy disks compared to single-band classifications. Specifically, Cibinel et al. (2015) demonstrated that the asymmetry $-M_{20}$ classification results in $\lesssim$20% contaminations when performed on stellar-mass maps, whereas the contamination fraction is $\sim$50% using the single-band HST/F160W images alone. This is because rest-frame optical images often contain optically bright clumps of star formation in addition to the dominant stellar-mass concentrations; spatially resolved SED fitting is required to distinguish them.

The positional accuracy, $\sigma_{\text{pos}}$, with which we can pinpoint the locations of star formation from the VLA and ALMA images depends on (1) the positional accuracy of the phase calibrators used as the astrometric reference (these errors are $<\!2$ mas for our observations); and (2) the signal-to-noise ratio ($S/N$) of the detection, since $\sigma_{\text{pos}} \approx \theta_{\text{beam}}/\sqrt{2S/N}$, following Condon (1997), corresponding to $\approx\!40$ and 60 mas for VLA and ALMA, respectively, at their detection limits. Comparing optical, radio, and millimeter morphologies further requires accurate astrometric alignment between the wave bands. The positions of emission peaks from the VLA and ALMA images are mutually consistent within 40 mas, corresponding to 0.3 kpc at $z = 2$, indicating good astrometric agreement despite observing different phase calibrators, and neither shows systematic offsets compared to 2MASS positions. The VLA primary beam extends beyond the HUDF to a total area of 61 arcmin$^2$ with an rms sensitivity better than $1 \mu$Jy beam$^{-1}$, comspatial with the CANDELS imaging (Grogin et al. 2011; Koekemoer et al. 2011). This area contains 68 bright point sources ($>\!8\sigma$) detected in both VLA and HST/F160W images that we use to compare the VLA astrometry (and by proxy, that of ALMA) against the HST astrometry. An offset of $\Delta\alpha = -80 \pm 110$ mas, $\Delta\delta = 260 \pm 130$ mas is required to bring HST astrometry into agreement with those of VLA and ALMA. This offset is constant throughout the field (a possible source of offset is discussed in D16). We apply the offset to all HST images for further analysis; the resulting median systematic offset is $<\!10$ mas in both $\alpha$ and $\delta$, with an rms dispersion of 150 mas.

3. RESULTS

Our sample is based on the ALMA-selected SFGs at $z = 1$–3, which D16 has demonstrated are representative of massive main-sequence SFGs at these redshifts. At $z = 1$–3, the D16 sample contains 13 galaxies detected in the ALMA image at $>\!3.5\sigma$, 11 of these are detected at $5$–$30\sigma$ in the VLA image; the remaining two, UDF10 and 15 (see Table 2 of D16), are detected at 2.5–3.0$\sigma$ in the VLA image (they are $3.6$–$4.0\sigma$ in the ALMA image), which we exclude from further morphological analysis. Since their stellar masses and rest-frame optical extents (an upper limit of star formation size) are similar to the rest of the sample, their exclusion will not bias our conclusions. The final sample contains 11 galaxies selected from ALMA and VLA, tabulated in Table 1 and shown in Figure 1. The sample has a median redshift of $z = 2.2$ and SFRs and stellar masses ranging from 40–326 $M_\odot$ yr$^{-1}$ (median = $102 M_\odot$ yr$^{-1}$) and $10.3 \leq \log(M_{\ast}/M_\odot) \leq 11.2$ (median = 10.7), respectively, implying vigorous assembly of stellar mass with a median mass-doubling time of 0.4 Gyr. It is worth noting that there are six VLA-detected galaxies in the HUDF in the same redshift range that are not detected by ALMA, these galaxies are at $z_{\text{median}} = 1.3$ where VLA is sensitive to lower SFRs than ALMA (e.g., at $z = 1.0$, the VLA detection limit is $\approx 10 M_\odot$ yr$^{-1}$, whereas that of ALMA is $\approx 60 M_\odot$ yr$^{-1}$). Likewise, their exclusion will not bias the conclusion for massive galaxies above $10^{10}M_\odot$. The morphological classification performed on stellar-mass maps (Section 2) indicates that five and four galaxies are unambiguously isolated and morphologically disturbed, respectively; the remaining three exhibit substructures in stellar-mass maps that could indicate, e.g., late-stage mergers (classifications indicated in Figure 1).

We first show that individual SFGs follow the far-infrared/radio correlation (Section 3.1), then report on the sizes and locations of their star formation in relation to the stellar-mass buildup of host galaxies (Section 3.2), the dust and gas masses and the implied gas fraction and depletion time (Section 3.3), the spatially resolved $\Sigma_{\text{SFR}}$ and their potential implications on star formation-driven outflows (Section 3.4), as well as discussing our general observations on star formation in AGN hosts (Section 3.5).

3.1. Far-infrared/Radio Correlation at $z \sim 1$–3

The far-infrared/radio correlation (Helou et al. 1985) has been demonstrated to hold in statistical (i.e., stacked) samples out to $z \sim 2$ (Ivison et al. 2010; Magnelli et al. 2015; Pannella et al. 2015), but this has not been established for individual galaxies in the main sequence of SFGs. Yet the correlation is important as a physically motivated rationale that radio and millimeter measurements trace the same extent of star-forming regions. We explore the validity of the correlation for the individual VLA-ALMA-selected SFGs by investigating their $S_{3\text{ cm}}/S_{1.3\text{ mm}}$ flux ratios as a function of redshift (e.g., Carilli & Yun 1999) in comparison with the ratio predicted by the Rieke et al. (2009) infrared SED library. To ascertain that the 5 cm fluxes for galaxies in this test are not contaminated by AGN emission, we select a star-forming subsample by conservatively excluding all objects with $L_{5\text{ cm}}(0.5$–$8 \text{ keV}) \gtrsim 3 \times 10^{42}$ erg s$^{-1}$ following Xue et al. (2011). That is, AGN candidates are excluded from this test on the basis of X-ray luminosity regardless of whether their radio emission is enhanced over the level of star formation.
We find that for all five galaxies in the star-forming subsample, $S_{\text{cm}}/S_{1.3\,\text{mm}}$ at $z \sim 1–3$ is well predicted by that of a local SED template for SFG with infrared luminosity of $10^{11.5} L_{\odot}$ (Figure 2); the rms scatter of the ratio for these five objects around this template is 0.03 dex, which suggests that the far-infrared/radio correlation does hold for these SFGs. That a single SED is a good descriptor of the flux ratio should not come as a surprise because of the small dispersion of dust temperature found in main-sequence SFGs at $z \sim 2$ (Magnelli et al. 2014; Scoville et al. 2016).

On the other hand, all but one of the X-ray-selected AGN also have enhanced $S_{\text{cm}}$ over the level of $S_{1.3\,\text{mm}}$ to varying degrees. X-ray-luminous AGN (UDF1 and 8) tend to have relatively small radio enhancement, having $S_{\text{cm}}$ similar to the levels predicted by the far-infrared/radio correlation given the $S_{1.3\,\text{mm}}$, whereas X-ray-weak AGN candidates are among the most radio-luminous. This may be indicative of an anticorrelation between radiative and mechanical power that is well established in the local universe (Best & Heckman 2012), although analysis of a larger radio and X-ray sample is required to test whether such a dichotomy in the Eddington ratio of AGN exists at high redshift.

We classify as radio AGN those with $S_{\text{cm}}$ enhancement greater than twice the highest level expected from star-forming SED templates, about 0.5 dex higher than the observed $S_{\text{cm}}/S_{1.3\,\text{mm}}$ ratios for the star-forming subsample, which is $\sim 15 \times$ the rms scatter of $S_{\text{cm}}/S_{1.3\,\text{mm}}$. These radio AGN are indicated in Table 1 and in the corresponding VLA images in Figure 1.

In the local universe, e.g., Beck (2007) and Fletcher et al. (2011) have shown that the morphology of SFGs is not strongly dependent on frequency at $\sim 1.5–10$ GHz (20–3 cm). While the rest-frame $\sim 1.7$ cm continuum emission probed by our 5 cm VLA observations at $z \sim 2$ is dominated by synchrotron emission from supernovae, the thermal fraction associated with H II regions that may be present is not well constrained at these redshifts, and could be a source of morphological uncertainties. The agreement between $S_{\text{cm}}/S_{1.3\,\text{mm}}$ and the template prediction suggests that this is not a strong effect. Furthermore, the scale length of synchrotron and inverse-Compton losses decreases rapidly with the increasing $\alpha$ and are less than 1 kpc at the typical $\alpha$ of $\sim 2$ SFGs (Murphy et al. 2006). Hence we expect that for non-AGN, both VLA and ALMA should trace similar extents of star formation, i.e., if both are detected at a comparable S/N, they can be used interchangeably to measure the extinction-free star formation sizes at $z \sim 1–3$. Although we are only able to investigate the correlation for five SFGs with SFRs $\sim 100 M_{\odot}\text{ yr}^{-1}$, the success of this test highlights that it would be worth examining a larger sample of individual SFGs with future ultra-deep VLA and ALMA observations.

### 3.2. Sizes and Locations of Intense Star Formation

In Figure 1, a common picture emerges that intense star formation usually occurs at the location of the stellar-mass concentration. The sizes of star-forming regions, as independently traced by the VLA and ALMA images, extend over the dominant buildup of stellar mass. In this section, we discuss the size and location measurements to quantify these observations. None of these star-forming regions are identifiable in the rest-frame ultraviolet images, which show clumps of unobscured star formation that tend to occur in the areas peripheral to the intense obscured star formation and that comprise 0.1%–5% of the galaxy-integrated SFR (D16).

We measure the extinction-independent effective radius of the star formation distribution, $r_{\text{SFR}}$, from the VLA and ALMA images by means of deconvolved source sizes determined from two-dimensional (2D) elliptical Gaussian fitting (the position, major and minor axes, flux, and position angle of the Gaussians are free parameters); multiple Gaussians are permitted, in which case all Gaussian components from the same flux “island” above 3$\sigma$ are grouped together into a source. This is carried out with PyBDSM. We also construct images at multiple $u$, $v$ tapering scales to accurately measure sizes of extended sources; the scale with the highest $S/N$ for each object is adopted for the size measurement. All VLA sources except UDF1 and all ALMA sources are well described with a

Note. IDs, stellar mass ($M_{\star}$), SFR$_{\text{IR}}$ and redshift ($z$) are from D16. Photometric redshifts are reported with two decimal points (redshifts are measured spectroscopically otherwise). The peak positions of VLA sources ($\alpha_{\text{radio}}$ and $\delta_{\text{radio}}$) are from Gaussian fits (Section 3.2). We do not tabulate the AGN positions here because they are cospatial with the radio positions within 40 mas (Section 2); VLA and ALMA flux densities are tabulated in Tables 2 and 3 of D16. The X-ray and radio AGN definitions are discussed in Section 3.1; for AGN, SFR$_{\text{radio}}$ are conservatively reported as upper limits. The dust mass $M_{\text{dust}}$, gas fraction $f_{\text{gas}}$, and gas-depletion time $t_{\text{dust}}$ estimation assumes $T_{\text{dust}} = 25$ K (Magnelli et al. 2014; Scoville et al. 2016) and a gas-to-dust ratio of 100; uncertainties of these quantities reflect the range of possible values from the assumption of $T_{\text{dust}} = 20–30$ K (details in Section 3.3).

### Table 1

VLA and ALMA-selected Star-forming Galaxies at $z = 1–3$ in the HUDF

<table>
<thead>
<tr>
<th>ID</th>
<th>$\alpha_{\text{radio}}$</th>
<th>$\delta_{\text{radio}}$</th>
<th>$z$</th>
<th>$M_{\star}$</th>
<th>AGN?</th>
<th>SFR$_{\text{radio}}$</th>
<th>SFR$_{\text{IR}}$</th>
<th>$M_{\text{dust}}$</th>
<th>$f_{\text{gas}}$</th>
<th>$t_{\text{dust}}$</th>
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</thead>
<tbody>
<tr>
<td>UDF1</td>
<td>53.18346</td>
<td>−27.77664</td>
<td>3.00</td>
<td>10.7 ± 0.1</td>
<td>X-ray</td>
<td>$&lt;476$</td>
<td>326 ± 83</td>
<td>9.6 ± 0.2</td>
<td>0.70 ± 0.08</td>
<td>0.35 ± 0.24</td>
</tr>
<tr>
<td>UDF2</td>
<td>53.18137</td>
<td>−27.77358</td>
<td>2.794</td>
<td>11.1 ± 0.2</td>
<td>...</td>
<td>311 ± 28</td>
<td>247 ± 76</td>
<td>9.1 ± 0.2</td>
<td>0.51 ± 0.12</td>
<td>0.52 ± 0.16</td>
</tr>
<tr>
<td>UDF3</td>
<td>53.16060</td>
<td>−27.77628</td>
<td>2.543</td>
<td>10.3 ± 0.2</td>
<td>X-ray</td>
<td>$&lt;470$</td>
<td>195 ± 69</td>
<td>9.1 ± 0.1</td>
<td>0.80 ± 0.05</td>
<td>0.60 ± 0.37</td>
</tr>
<tr>
<td>UDF4</td>
<td>53.17092</td>
<td>−27.77544</td>
<td>2.43</td>
<td>10.5 ± 0.2</td>
<td>...</td>
<td>116 ± 23</td>
<td>94 ± 4</td>
<td>8.0 ± 0.16</td>
<td>0.57 ± 0.08</td>
<td>0.45 ± 0.27</td>
</tr>
<tr>
<td>UDF5</td>
<td>53.15402</td>
<td>−27.79090</td>
<td>1.759</td>
<td>10.4 ± 0.2</td>
<td>...</td>
<td>101 ± 7</td>
<td>102 ± 7</td>
<td>8.7 ± 0.16</td>
<td>0.66 ± 0.07</td>
<td>0.48 ± 0.13</td>
</tr>
<tr>
<td>UDF6</td>
<td>53.14350</td>
<td>−27.78328</td>
<td>1.413</td>
<td>10.5 ± 0.1</td>
<td>...</td>
<td>80 ± 5</td>
<td>87 ± 11</td>
<td>8.0 ± 0.16</td>
<td>0.55 ± 0.09</td>
<td>0.44 ± 0.11</td>
</tr>
<tr>
<td>UDF7</td>
<td>53.18052</td>
<td>−27.77971</td>
<td>2.59</td>
<td>10.6 ± 0.1</td>
<td>X-ray</td>
<td>$&lt;770$</td>
<td>56 ± 22</td>
<td>8.5 ± 0.2</td>
<td>0.44 ± 0.08</td>
<td>0.56 ± 0.16</td>
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<tr>
<td>UDF8</td>
<td>53.16558</td>
<td>−27.76989</td>
<td>1.546</td>
<td>11.2 ± 0.2</td>
<td>X-ray</td>
<td>$&lt;87$</td>
<td>149 ± 90</td>
<td>8.5 ± 0.16</td>
<td>0.17 ± 0.04</td>
<td>0.22 ± 0.06</td>
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<tr>
<td>UDF11</td>
<td>53.16688</td>
<td>−27.79885</td>
<td>1.998</td>
<td>10.8 ± 0.1</td>
<td>X-ray</td>
<td>$&lt;202$</td>
<td>162 ± 94</td>
<td>8.5 ± 0.2</td>
<td>0.31 ± 0.10</td>
<td>0.18 ± 0.05</td>
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<tr>
<td>UDF13</td>
<td>53.14616</td>
<td>−27.77995</td>
<td>2.497</td>
<td>10.8 ± 0.1</td>
<td>X-ray</td>
<td>$&lt;166$</td>
<td>68 ± 18</td>
<td>8.4 ± 0.2</td>
<td>0.26 ± 0.10</td>
<td>0.35 ± 0.23</td>
</tr>
<tr>
<td>UDF16</td>
<td>53.17661</td>
<td>−27.78551</td>
<td>1.319</td>
<td>10.9 ± 0.1</td>
<td>...</td>
<td>46 ± 3</td>
<td>40 ± 18</td>
<td>8.4 ± 0.2</td>
<td>0.24 ± 0.07</td>
<td>0.62 ± 0.15</td>
</tr>
</tbody>
</table>

Note: $\alpha$ and $\delta$ are J2000.0 coordinates. The units for the flux density are mJy beam$^{-1}$.
Figure 1. (a) The 11 SFGs detected in both the VLA and ALMA HUDF surveys. From left to right: *HST* i$_{160} - J_115 - H_160$ color composite; unobscured star formation (*HST* /F060W); rest-frame optical morphology (*HST* /F160W); star formation rate surface density, $\Sigma_{\text{SFR}}$ (VLA, details in Section 2); dust mass surface density, $\Sigma_{\text{dust}}$ (ALMA, details in Section 3.3); and stellar-mass surface density, $\Sigma_{\text{M}}$ (from spatially resolved SED fitting using *HST* images). Each image is $4' \times 4'$. North is up, east is on the left. VLA and ALMA synthesized beams are shown in the corresponding columns; the contours are $\Sigma$ from the AGN and hence are marked as $(\cdot)$. (b) We note that the radio $\Sigma_{\text{SFR}}$ maps for UDF7, 11, and 13 that harbor X-ray–selected AGN with radio emission enhanced more than twice the level of star formation emission (Figure 2, tabulated in Table 1) may contain significant contributions from the AGN and hence are marked as "radio AGN".

To compare the star formation size with the size of stellar-mass buildup, we quantify the effective radius for the stellar-mass distribution by a circularized radius that encircles half of the stellar mass, $r_{\text{MS}}$. This is measured non-parametrically from the stellar-mass map by determining the area of the stellar-mass "island" that contains half of the total stellar mass. In the case of an ideal Gaussian, this size measure yields the same radius as Gaussian fitting, which allows us to compare the $r_{\text{MS}}$ and $r_{\text{SF}}$. The $r_{\text{MS}}$ uncertainties are estimated from the rms maps from the spatially resolved SED fitting.

The location of star formation concentration in relation to the hosts’ stellar-mass concentration is quantified by the separation between their barycenters, measured from stellar-mass and star formation maps. The choice of band(s) used to determine the star formation barycenter follows those of $r_{\text{SF}}$ measurements, i.e., we adopt a geometric mean of the barycenters from VLA and ALMA images for non-AGN, and from the ALMA images otherwise (again, UDF8 is treated as a non-AGN for this purpose). The uncertainties are the sum in quadrature of the resolution of the stellar-mass maps ($0.06$ kpc) and the uncertainties of Gaussian fitting ($\theta_{\text{beam}} / 2S/N$), see $\sigma_{\text{pos}}$ discussion in Section 2).

We find a median SFGs $r_{\text{SF}}$ of $2.1 \pm 0.9$ kpc. The separations between the star formation and stellar-mass barycenters (Table 3) are smaller than this radius in all cases except UDF7, which appears to be interacting. On average, the star formation and stellar-mass peaks are separated by 0.53 $r_{\text{MS}}$ (the median separation is 0.36 $r_{\text{MS}}$, corresponding to 0.9 kpc), indicating that star formation occurs nearly cospatially with the stellar-mass concentration. Individually, the $r_{\text{SF}}$ broadly follows the $r_{\text{MS}}$ (median $r_{\text{MS}} = 2.6 \pm 0.7$ kpc) as shown in Figure 3 (left panel), consistent with the picture of star formation concentration in relation to the hosts’ stellar-mass concentration.
Figure 1. (Continued.)
formation extending over a similar area as the stellar mass distribution (see also Tacconi et al. 2013). Our median SFG radius is similar to those of the star formation-dominated subsample of Biggs & Ivison (2008) SMGs, which has a median $r_{\text{SF}}$ of 2.3 kpc (their star-forming subsample is reproduced in Figure 3). On the other hand, the SFG radii from our sample are ≈2–4 times larger than those of submillimeter (submm) galaxies studied with ALMA at high resolution ($r_{\text{SF}} = 0.7–1.2$ kpc, Ikarashi et al. 2015; Simpson et al. 2015). These submm galaxies have median SFRs 3–8 times those of our sample, hence an explanation is that the more compact star formation could be a characteristic of this higher SFR regime, but not for main-sequence SFGs (Figure 3).

The increasing compactness of SFGs at SFR $> 300 M_\odot$ yr$^{-1}$ is independently indicated by the decreasing ratio of mid-infrared aromatic luminosity to the total infrared luminosity ($L_{\alpha,2 \mu m}/L_{\text{IR}}$ trend shown in Figure 3; Nordon et al. 2012; Pope et al. 2013; Shipley et al. 2016). Consistent with this trend, the two most intense SFGs in our sample (UDF1 and 2) are also the most compact. While the origin of such a trend is at present unclear, if the SFR in the interstellar medium (ISM) is regulated by some internal process (such as a form of stellar feedback), then this result may be expected. At the centers of massive starbursts, the surface densities rise and the vertical weight on a cloud can hinder the effectiveness of internal feedback processes at dispersing gas (e.g., Ostriker & Shetty 2011; Hopkins et al. 2013). In these environments, the star formation efficiency can rise dramatically, decreasing the effective radius for star formation within the galaxy.

A major result from our observations is that main-sequence SFGs at $z \sim 2$ exhibit vigorous star formation in a distributed manner over size scales of a few kiloparsec, which is to be contrasted with the typical sub-kpc size scales of comparably luminous galaxies in the local universe (Condon et al. 1991; Rujopakarn et al. 2011; Lutz et al. 2016). Such a result is expected from numerical simulations of galaxy growth via in situ star formation that is fed by the accretion of gas from the intergalactic medium (e.g., Dekel et al. 2009; Davé et al. 2010). High-resolution cosmological simulations have found that star formation may remain distributed even in the most extreme systems at this epoch in the submillimeter-selected population (Narayanan et al. 2015).

The similarity of the effective radii of star formation and stellar-mass in our field-galaxy sample is to be contrasted with recent ALMA high-resolution studies that target high-SFR galaxies. Barro et al. (2016) found that the 870 $\mu$m sizes of main-sequence SFGs (median SFR $\sim 300 M_\odot$ yr$^{-1}$) selected by their compact optical sizes to represent the progenitors of compact galaxies at $z \sim 2$ have $r_{\text{SF}} \approx 1$ kpc, and do not show the star formation size trend with stellar-mass distribution sizes. Similarly, the Tadaki et al. (2016) 870 $\mu$m observations of H$\alpha$-selected SFGs at $z = 2.2$–2.5 (median SFR $\sim 230 M_\odot$ yr$^{-1}$, also in the main sequence) found their star formation radii to be $< 1.5$ kpc, a factor of two smaller than their average rest-frame optical size (3.2 kpc). These two samples exhibit a similar characteristic of star formation distribution to our sample in that their star formation concentrations also peak at the same locations as the hosts’ stellar-mass concentrations, but they appear to lack the extended component of galaxy-wide star formation. A possible explanation is that the extended star formation at $z \sim 3$ (on-sky diameter $\sim 2.5$ kpc) is not visible in the rest-frame UV image.
Table 2
Size Measurement of VLA-ALMA-selected Star-forming Galaxies

<table>
<thead>
<tr>
<th>ID</th>
<th>VLA r1/2 (kpc)</th>
<th>ALMA r1/2 (kpc)</th>
<th>Adopted r$_{SF}$ (kpc)</th>
</tr>
</thead>
<tbody>
<tr>
<td>UDF1</td>
<td>1.5 ± 0.2</td>
<td>1.4 ± 0.7</td>
<td>1.8 ± 0.6</td>
</tr>
<tr>
<td>UDF2</td>
<td>1.6 ± 0.2</td>
<td>2.5 ± 0.9</td>
<td>1.1 ± 0.5</td>
</tr>
<tr>
<td>UDF3</td>
<td>1.8 ± 0.3</td>
<td>2.0 ± 0.5</td>
<td>1.7 ± 0.5</td>
</tr>
<tr>
<td>UDF4</td>
<td>1.9 ± 0.5</td>
<td>2.4 ± 1.0</td>
<td>0.8 ± 0.6</td>
</tr>
<tr>
<td>UDF5</td>
<td>2.1 ± 0.6</td>
<td>2.0 ± 0.5</td>
<td>0.8 ± 0.6</td>
</tr>
<tr>
<td>UDF6</td>
<td>2.2 ± 0.7</td>
<td>&lt;0.6</td>
<td>1.9 ± 1.2</td>
</tr>
<tr>
<td>UDF7</td>
<td>&lt;3.0</td>
<td>2.7 ± 1.2</td>
<td>5.8 ± 0.6</td>
</tr>
<tr>
<td>UDF8</td>
<td>3.2 ± 1.0</td>
<td>1.8 ± 0.4</td>
<td>0.8 ± 0.6</td>
</tr>
<tr>
<td>UDF9</td>
<td>4.2 ± 1.7</td>
<td>4.0 ± 0.7</td>
<td>1.1 ± 0.7</td>
</tr>
<tr>
<td>UDF10</td>
<td>2.6 ± 1.1</td>
<td>1.9 ± 0.8</td>
<td>0.9 ± 0.7</td>
</tr>
<tr>
<td>UDF11</td>
<td>3.4 ± 0.9</td>
<td>2.5 ± 0.4</td>
<td>&lt;0.7</td>
</tr>
</tbody>
</table>

Note. Circularized SFG radii, r$_{SF}$, are measured from the deconvolved FWHMs of VLA and/or ALMA images with the following criteria, as noted for each object in the r$_{SF}$ column: (a) non-AGN, both VLA and ALMA are resolved, take average of the two bands; (b) non-AGN, only VLA is resolved, use VLA; (c) AGN, ALMA is resolved, use ALMA; (d) AGN, ALMA is unresolved, adopt 2r deconvolved size limit as upper limit; and (e) AGN present, but has negligible contribution to radio emission, both VLA and ALMA are resolved, same as (a). Uncertainties in the deconvolved FWHM columns are 1 – σ. UDF11 has two VLA components; its r$_{SF}$ reflects the total size.

Table 3
Sizes and Locations of Star Formation

<table>
<thead>
<tr>
<th>ID</th>
<th>r$_{SF}$ (kpc)</th>
<th>r$_{SM}$ (kpc)</th>
<th>$\Delta_\Sigma(M_\odot, SF)$ (kpc)</th>
<th>Average $\Sigma_{\text{AVG}}$ ($M_\odot$ yr$^{-1}$ kpc$^{-2}$)</th>
</tr>
</thead>
<tbody>
<tr>
<td>UDF1</td>
<td>1.4 ± 0.2</td>
<td>1.4 ± 0.7</td>
<td>0.6 ± 0.5</td>
<td>18.0 ± 0.3</td>
</tr>
<tr>
<td>UDF2</td>
<td>1.6 ± 0.2</td>
<td>2.5 ± 0.9</td>
<td>0.7 ± 0.5</td>
<td>11.3 ± 0.7</td>
</tr>
<tr>
<td>UDF3</td>
<td>1.8 ± 0.3</td>
<td>2.5 ± 0.7</td>
<td>1.7 ± 0.5</td>
<td>6.4 ± 3.0</td>
</tr>
<tr>
<td>UDF4</td>
<td>1.9 ± 0.5</td>
<td>2.5 ± 1.0</td>
<td>0.8 ± 0.6</td>
<td>3.3 ± 1.8</td>
</tr>
<tr>
<td>UDF5</td>
<td>2.1 ± 0.6</td>
<td>2.5 ± 0.5</td>
<td>0.9 ± 0.6</td>
<td>2.5 ± 1.4</td>
</tr>
<tr>
<td>UDF6</td>
<td>2.2 ± 0.7</td>
<td>3.5 ± 0.7</td>
<td>&lt;0.6</td>
<td>1.9 ± 1.2</td>
</tr>
<tr>
<td>UDF7</td>
<td>&lt;3.0</td>
<td>2.7 ± 1.2</td>
<td>5.8 ± 0.6</td>
<td>&gt;2.9</td>
</tr>
<tr>
<td>UDF8</td>
<td>3.2 ± 1.0</td>
<td>1.8 ± 0.4</td>
<td>0.8 ± 0.6</td>
<td>1.2 ± 0.9</td>
</tr>
<tr>
<td>UDF9</td>
<td>4.2 ± 1.7</td>
<td>4.0 ± 0.7</td>
<td>1.1 ± 0.7</td>
<td>1.0 ± 1.0</td>
</tr>
<tr>
<td>UDF10</td>
<td>2.6 ± 1.1</td>
<td>1.9 ± 0.8</td>
<td>0.9 ± 0.7</td>
<td>1.1 ± 0.9</td>
</tr>
<tr>
<td>UDF11</td>
<td>3.4 ± 0.9</td>
<td>2.5 ± 0.4</td>
<td>&lt;0.7</td>
<td>0.4 ± 0.2</td>
</tr>
</tbody>
</table>

Note. $\Delta_\Sigma(M_\odot, SF)$ is the physical separation between the barycenters of stellar mass and star formation concentrations (details of sizes and separation measurements are in Section 3.2); upper limits indicate separations smaller than the positional uncertainties. The $\Sigma_{\text{AVG}}$ is averaged over the entire star formation surface areas indicated by r$_{SF}$.

formation found in field galaxies could be a distinct evolutionary state from (and possibly preceding) the compact star formation in the Barro et al. (2016) and Tacconi et al. (2016) samples, with the latter representing a state closer to the conclusion of bulge formation and the subsequent cessation of the bulk of star formation (e.g., Barro et al. 2015).

3.3. Dust and Gas Properties

We estimate the dust masses from the 1.3 mm fluxes using the Li & Draine (2001) dust mass absorption coefficients. The dominant source of uncertainty in dust mass estimates is the dust temperature $T_{\text{dust}}$, which is not known precisely for our sample, but Herschel studies have shown that $T_{\text{dust}}$ of main-sequence SFGs at $z \sim 2$ is in the range of 20–30 K (Dunne et al. 2011; Auld et al. 2013), prompting us to adopt $T_{\text{dust}} = 25$ K (Magnelli et al. 2014; Kirkpatrick et al. 2015; Scoville et al. 2016) and estimate the uncertainties using $T_{\text{dust}} = 20$ and 30 K for upper and lower limits, respectively. Since the ALMA measurements are in the Rayleigh–Jeans regime, the dust masses have only a moderate dependence on the assumed temperature (Hildebrand 1983; Casey 2012; Elia & Pezzuto 2016). Dust masses are in the range of log($M_{\text{dust}}/M_\odot$) = 8.4–9.1 with a typical uncertainty of 0.2 dex (tabulated in Table 1). The dust mass estimates are in agreement with the Magdis et al. (2012) results of stacking analysis for $z \sim 2$ main-sequence SFGs, suggesting that the dust masses of our HUDF sample are representative of main-sequence populations at $z \sim 2$.

If we assume a gas-to-dust ratio of 100 (Leroy et al. 2011; Magdis et al. 2012), the gas fraction, $f_{\text{gas}} = M_{\text{gas}}/(M_{\text{star}} + M_{\text{gas}})$, of our sample is on average 0.5 ± 0.2, indicating the gas-rich nature of these systems similar to the normal SFGs at $z \sim 2$ studied by Tacconi et al. (2013), although significantly more gas rich than typical submm galaxies (e.g., Narayanan et al. 2012). The average gas depletion time, $\tau_{\text{dep}} = M_{\text{gas}}/f_{\text{gas}}$, of our sample is 0.4 ± 0.2 Gyr, comparable to their average mass-depleting time of 0.4 Gyr, assuming that star formation-driven outflows do not result in irreversible mass losses. These $\tau_{\text{dep}}$ are similar to the main-sequence samples of Sargent et al. (2014) and Silverman et al. (2015), further ensuring the main-sequence nature of our sample independent of the presence of AGN. We note that if the $T_{\text{dust}}$ is warmer than the assumed 25 K, the gas fraction could be lower (e.g., $T_{\text{dust}} = 30$ K implies average $f_{\text{gas}}$ and $\tau_{\text{dep}}$ of 0.4 ± 0.2 and 0.3 ± 0.1 Gyr) and hence requires a significant amount of cold gas accretion to sustain the current level of star formation for a significant period of time. However, the metallicities of our SFGs are expected to be lower than for local analogs (e.g., Zahid et al. 2013). The rapid reduction in dust mass with decreasing metallicity (Draine et al. 2007) would then argue for larger gas fractions in these galaxies.

The cospatiality and the similar effective radii of radio and millimeter emissions imply that in most cases, the cold gas associated with star formation resides in the star-forming regions, which we have shown occur near the stellar-mass concentrations. Hence, if this level of SFR is sustained through a mass-doubling time, we expect that the newly formed stellar masses will occur near the existing stellar-mass concentrations, i.e., enhancing the stellar mass surface density near the central
region, consistent with a picture of in situ bulge formation (see also Barro et al. 2016; Tadaki et al. 2016).

3.4. $\Sigma_{\text{SFR}}$ Implies Galaxy-wide Star Formation-driven Outflows?

The spatially resolved $\Sigma_{\text{SFR}}$ for 9 of 11 galaxies peaks within $\lesssim \! 1 \, \text{kpc}$ of the stellar-mass distribution (with the exception of UDF3 and 7; Table 3), and all have $\Sigma_{\text{SFR}}$ in the range of $0.4$–$18 \, M_\odot \, \text{yr}^{-1} \, \text{kpc}^{-2}$ (median $\Sigma_{\text{SFR}} = 2.5 \, M_\odot \, \text{yr}^{-1} \, \text{kpc}^{-2}$), the higher end of which is comparable to local starbursts (Kennicutt 1998). The level of $\Sigma_{\text{SFR}}$ found in our main-sequence sample is significantly below those found in luminous submm galaxies of 90–100 $M_\odot \, \text{yr}^{-1} \, \text{kpc}^{-2}$ (Ikarashi et al. 2015; Simpson et al. 2015), suggesting very different conditions for star formation between them.

Recent studies show that $\Sigma_{\text{SFR}}$ is a strong predictor of star formation-driven outflows at $z = 1$–2. An increased broad fraction of H$_\alpha$ line flux at $z \sim 2$ associated with star formation-driven outflows is found to strongly depend on $\Sigma_{\text{SFR}}$, with a $\Sigma_{\text{SFR}}$ of $1 \, M_\odot \, \text{yr}^{-1} \, \text{kpc}^{-2}$ being the threshold above which the large broad-line fractions are observed (Newman et al. 2012). At a similar $\Sigma_{\text{SFR}}$ threshold of $\approx 0.3 \, M_\odot \, \text{yr}^{-1} \, \text{kpc}^{-2}$, Bordoloi et al. (2014) also found a significant increase in the equivalent width of the Mg II absorption line at $z \sim 1$, which indicates cool outflowing gas. If $\Sigma_{\text{SFR}}$ is indeed a predictor of star formation-driven outflows at these redshifts, we expect that main-sequence SFGs like those in our sample, whose galaxy-average $\Sigma_{\text{SFR}}$ are higher than these threshold levels (Table 3), will harbor outflows across a significant extent comparable to the areas of stellar-mass buildup (i.e., galaxy-wide). While the presence of outflows appears to depend on $\Sigma_{\text{SFR}}$, how the outflow rate depends on $\Sigma_{\text{SFR}}$ is not well established, especially for molecular gas in a turbulent galactic environment at $z \sim 2$.

If a strong relationship also exists between them, we expect that the areas exhibiting high $\Sigma_{\text{SFR}}$ near galaxy centers will also be the areas with the largest outflows. Future spatially resolved observations of cold gas dynamics could provide definitive evidence whether this intense galaxy-wide star formation also drives galaxy-wide outflows.

3.5. Remarks on Star Formation in AGN Hosts at $z \sim 2$

In Section 3.1 we have shown that the 4 Ms Chandra observations in the HUDF allows us to detect X-ray AGN candidates residing in main-sequence SFGs with varying degrees of radio emission enhancement over the level corresponding to their SFRs. Six of 11 galaxies in our SFG sample are detected in X-rays, with the most luminous reaching $L_{0.5-8 \text{keV}} \simeq 5 \times 10^{43}$ ergs s$^{-1}$. Two of the X-ray detected SFGs (UDF3 and 11) are only detected in the soft X-ray band (0.5–2 keV), suggesting that their X-ray emission could originate solely from star formation. However, UDF11 also has radio emission enhanced by more than twice the level associated with star formation (Figure 2 Left; this level of radio enhancement is $15 \times$ the scatter around the emission level implied by the far-IR/radio correlation, Section 3.1), which suggests the presence of radio AGN.

Omitting the only ambiguous case of UDF3 from the AGN subsample, the AGN fraction (5 of 11 galaxies) is still higher than typically reported (e.g., Xue et al. 2010) at these redshifts, but not unexpected since AGN activity is known to depend on both stellar mass and gas mass, as evident by its clear dependence on star formation of its host even on a galaxy-wide scale (e.g., Silverman et al. 2009; Brusa et al. 2009). With the ALMA continuum detections at 1.3 mm preferential to the most massive galaxies with high SFRs at $z \sim 1$–3 (D16), such a selection is also effective at probing rapidly growing
supermassive black holes at $z \sim 2$ including those with mild obscuration. Further support of such a high detection rate of AGN using ALMA comes from Umehata et al. (2015), who report a high X-ray AGN fraction in the SSA22 field.

Despite the high sensitivity of the Chandra observations in the HUDF, the X-ray counts in the full 0.5–8 keV band for four of six AGN candidates are still low (16–26 counts), highlighting the necessity of sensitive X-ray observations to characterize AGN populations in main-sequence SFGs at $z \sim 2$. In the remaining two sources with sufficient count statistics, UDF1 and UDF8, which have 434 and 878 full-band counts, respectively, the column densities are $1.5 \times 10^{21}$ and $7 \times 10^{22}$ cm$^{-2}$ respectively, based on Web-PIMMS, provided by HEASARC, assuming an intrinsic photon index of 2 and the redshift of the source, indicating mild obscuration.

Regardless of the AGN selection method, ALMA imaging shows that all AGN candidates in the sample contain large amounts of cold dust and gas (Table 1) associated with star formation. The trend of the star formation size with SFR shown in Figure 3 is consistent with the ALMA 870 μm sizes of star formation in the hosts of six X-ray detected AGN that have SFRs of $\approx 130-400 M_\odot$ yr$^{-1}$ (Harrison et al. 2016).

For UDF7, 11, and 13 where the radio flux enhancement allows pinpointing the AGN location using the VLA images, we find that the AGN are within 0.5 kpc of the barycenter of star formation of their host galaxies as measured from the ALMA images (an upper limit given by our ability to pinpoint AGN and star formation given their S/N). For UDF11 and 13, the star formation concentrations are $\approx 1$ kpc from the stellar-mass barycenters, whereas UDF7 is notable for these activities being $\approx 6$ kpc away from the dominant stellar-mass concentration, a possible consequence of galaxy interactions.

4. CONCLUSION
We have made the first comparison between ultra-deep VLA and ALMA imaging of $z \sim 2$ main-sequence SFGs. The far-infrared/radio correlation appears to hold for individual main-sequence SFGs with SFR $\sim 100 M_\odot$ yr$^{-1}$ out to $z \sim 3$, and the extinction-independent distributions of star formation are consistent between the data sets.

The intense star formation in our blank-field-selected sample extends over a large galactic area regardless of their stellar-mass morphology (isolated or morphologically disturbed) or the presence of AGN, with the SFR and dust-mass surface densities both peaking near the existing stellar-mass concentration. These findings provide direct-imaging evidence of a gas-rich galactic environment with widespread occurrence of intense star formation. The spatially resolved SFR surface densities are sufficiently high across the areas of dominant stellar-mass buildsups that they may drive galaxy-wide outflows. Where radio excess permits pinpointing of the AGN, it is found to be cospatial with the dust mass concentrations.

The median star formation diameter in our main-sequence SFGs sample with SFR $\sim 100 M_\odot$ yr$^{-1}$ is $4.2 \pm 1.8$ kpc, twice larger than those of submm galaxies forming stars at 3–8 times higher than the main-sequence SFGs, indicating that at $z \sim 2$, the SFR threshold above which a significant population of more compact SFGs appears to emerge is $\sim 300 M_\odot$ yr$^{-1}$.

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http://heasarc.gsfc.nasa.gov/Tools