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Ab-Initio Surface Hopping and Multiphoton Ionisation Study of the Photodissociation Dynamics of CS$_2$

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Abstract

New ab-initio surface hopping simulations of the excited state dynamics of CS$_2$ including spin-orbit coupling are compared to new experimental measurements using a multiphoton ionisation probe in a photoelectron spectroscopy experiment. The calculations highlight the importance of the triplet states even in the very early time dynamics of the dissociation process and allow us to unravel the signatures in the experimental spectrum, linking the observed changes to both electronic and nuclear degrees of freedom within the molecule.

Keywords: Photodissociation, Photoelectron spectroscopy, Theoretical Chemistry, Non-adiabatic dynamics

1. Introduction

The dissociation dynamics of CS$_2$ following UV excitation have been a benchmark in chemical dynamics for many years, with numerous experimental studies in both the time and frequency domain, see for example [1, 2, 3, 4, 5, 6, 7, 8, 9, 10]. This lasting fascination with CS$_2$ can be traced to the efficient dissociation, dictated by complex dynamics on multiple coupled electronic states. Despite intense experimental study, the fast dynamics

and the high ionisation limits of intermediates and final products have limited the experimental view to specific points along the full dissociation path, such that open questions remain even for this structurally simple molecule. The origin of the complexity derives from the near degeneracy of the optically bright $^1\text{B}_2(^1\Sigma^+_u)$ state with multiple other electronic states at linear geometry, which leads to highly efficient population transfer and strongly coupled multistate dynamics. The mixing of the electronic states leads to dissociation and the formation of a ground state CS ($^1\Sigma^+$) molecule in conjunction with atomic sulphur in either the spin forbidden ground state, $^3\text{P}$, or a spin allowed excited state, $^1\text{D}$. While the exact branching ratio has proven difficult to define accurately, the spin forbidden product is seen to dominate in most experimental studies [11, 2, 12], highlighting the importance of spin-orbit coupling for an accurate description. Considering how well studied this molecule has been experimentally, calculations of the dynamics have been limited with, as far as we are aware, no simulations accounting for the spin-orbit coupling that drives the dominant dissociation process. In this work we combine ab-initio surface hopping simulations of the dissociation dynamics of CS$_2$ with new photoelectron spectroscopy measurements using a multiphoton probe to study the effect of spin-orbit coupling on the early time dynamics of the molecule.

Previous dynamics calculations have focused on the singlet state dynamics and the effect of non-adiabatic coupling on measured photoelectron angular distributions obtained following excitation at 201 nm and ionisation with 268 nm[13]. These calculations provide a very good measure of the photoangular distribution, which they claim suggests that the lack of spin-orbit coupling in the model does not affect the calculated early time dynamics. It should be noted that the rather low energy probe used in the experiment means that only the singlet states could be ionised and that any effect of the triplet in the angular distributions would not be observable. The measurements and theory therefore do not take into account the population transfer between the initially excited singlet manifold and the accessible triplet states. This point is highlighted by recent time-resolved VUV photoelectron spectroscopy experiments by Spesyvtsev et al. [4] using a 20 fs 159 nm (7.8 eV) probe pulse. These experiments provide the most detailed maps of the excited state dynamics to date, and show large changes in electron kinetic energy as the molecule undergoes bending vibrations with an almost 3 eV shift in the measured electron kinetic energy in 40 fs. The probed dynamics occur on the singlet surfaces and, as the molecule continues to vibrate, the population is
transferred into lower lying electronic states which are outside the observation window provided even by their VUV probe.

2. Theory

2.1. Computational methods

Ab-initio electronic structure calculations were performed using the Molo pro 2015.1 [14] suite of programs at the CAS(16,12)/aug-cc-pvTZ and aug-
cc-pvQZ level of theory using the full valence (16,12) active space with bonding and antibonding $\sigma$ and $\pi$ molecular orbitals and sulfur atom lone pairs. Ground state geometry optimisation using CAS(16,12)/aug-cc-pvQZ resulted in $R_{CS}=1.569$ Å and $\theta_{SCS}=0^\circ$. Angular and radial cuts through the singlet and triplet potential energy surfaces are shown in Fig. 1, and vertical excitation energies and oscillator strengths for the first four excited singlet states are given in Table 1. The potential energy curves in Fig. 1 are broadly in keeping with previous ab-initio calculations [15, 16, 17, 18, 19].

Table 1: Vertical excitation energies ($\Delta E = E(S_i) - E(S_0)$) and oscillator strengths from the ground state to the first four excited singlet states of CS$_2$, calculated using SA5-CAS(16,12)-SCF/aug-cc-pvQZ with CASPT2 corrections to the energies. The excitation energies are calculated at the equilibrium geometry ($\theta_{SCS}=180^\circ$ and $R_{CS}=1.569$ Å), while oscillator strengths are calculated at $\theta_{SCS}=160^\circ$ since transitions are very weak in the linear geometry.

<table>
<thead>
<tr>
<th>State</th>
<th>Energy (eV)</th>
<th>Oscillator strength</th>
</tr>
</thead>
<tbody>
<tr>
<td>$S_1$</td>
<td>3.821</td>
<td>0</td>
</tr>
<tr>
<td>$S_2$</td>
<td>3.836</td>
<td>0.004282</td>
</tr>
<tr>
<td>$S_3$</td>
<td>3.836</td>
<td>0</td>
</tr>
<tr>
<td>$S_4$</td>
<td>6.430</td>
<td>0.000834</td>
</tr>
</tbody>
</table>

We simulate the dynamics of photoexcited CS$_2$ using the code SHARC [20, 21] interfaced with MOLPRO [14]. SHARC treats nuclear motion classically, but nonadiabatic effects and spin-orbit coupling [22] are included using the fewest-switches surface-hopping approach [23]. In contrast to previous singlet-only simulations [13], we propagate the dynamics on the four lowest singlet and triplet electronic states. To keep the simulations computationally feasible, we perform the electronic structure calculations at the SA8-CAS(8,6)-SCF/6-31+G* level, which reproduces the potential energy curves shown in Fig. 1 qualitatively. Initial positions are generated from a Wigner distribution based on the CAS(8,6)/6-31G* ground state vibrational frequencies and the oscillator strength of each geometry, and kinetic energy is assigned based on the required excitation energy and the experimental pump pulse energy. Following this protocol, 85\% of trajectories begin in the $B^1\Sigma_2$ state. A total of 369 trajectories are launched, of which 197 reach 500 fs and 114 reach 1000 fs, using a time step of 0.5 fs.
2.2. **Computational results**

Excitation of CS$_2$ triggers bending and vibrational motion in the molecule, as can be seen in Fig. 2, which shows the probability density evolution of the
molecular geometry as a function of the C-S bond-lengths, $R_{CS}$, and the bending angle, $\theta_{SCS}$. During the first 100 fs the vibrations are dominated by the symmetric stretch, but at later times energy flows into the asymmetric stretch. The frequencies of vibrations are somewhat over-estimated compared to the experimental values, presumably due to slight differences in the \textit{ab-initio} potential energy surfaces at the CAS(8,6) level. The total fraction of dissociated molecules in the full set of 369 trajectories is 22\%, which constitutes a lower bound since only a quarter of the trajectories reach 1000 fs. Dissociation occurs predominantly in the triplet states, with 89\% of the trajectories that dissociate occurring on the triplet surfaces. The lower degree of dissociation compared to the experiments can be traced to the topology of the potential energy surfaces at the level of \textit{ab-initio} theory employed in the simulations, as discussed below.

The electronic state populations as a function of time are shown in Fig. 3. Initial excitation onto the $S_2 \ 1B_2(1\Sigma^+_u)$ state is followed by rapid decay onto the singlet $S_3$ and $S_1$ potentials, as well as a redistribution of population onto the manifold of triplet states via spin-orbit coupling. The nonadiabatic transfer of population between the singlet states correlates strongly with the bending motion of the molecule, with efficient transfer predominantly occurring close to the linear geometry where states are (near)-degenerate. This gives rise to a rather periodic beating in both the individual singlet state populations, and in the total singlet population. Over time there is a build-up of population in $T_2$ at $t > 400$ fs, and a subsequent rise of population in $T_1$ at around $t > 800$ fs, due to population transfer from $T_2$ to $T_1$. The build-up in $T_1$ appears to be an artifact due to the SA8-CAS(8,6)-SCF/6-31+G* \textit{ab-initio} calculations, which increases the relative barrier height for dissociation on the $T_1$ and $T_2$ potentials by $\approx 1.5$ eV, hindering dissociation and leading to the observed accumulation of population in $T_2$. Consequently, it is reasonable to assume that the population trapped in $T_2$ in actual fact dissociates as observed in the experiment. Nevertheless, despite that the simulations underestimate the amount of $t < 1$ ps dissociation via the triplet states, the short-time $t < 400$ fs dynamics appears quite reliable.

3. Experiment

3.1. Experimental methods

The experiment has been described in detail previously [24]. Briefly, an amplified femtosecond laser system (Red Dragon, KM Labs) generates 30 fs
pulses of 800 nm light, with a pulse energy of up to 10 mJ at a repetition rate of 1 kHz. The pump pulse is produced via fourth harmonic generation of the fundamental (800 nm) beam, generating photons at around 200 nm. The 200 nm beam is produced using standard non-linear optics with sequential second, third and fourth harmonic generation in BBO giving a pulse energy of ~1 µJ. The 400 nm probe is generated by second harmonic generation of the fundamental laser output, producing approximately 5 µJ per pulse. The pump and probe beams are reflection focused in a near collinear geometry and cross at the centre of the interaction region of a velocity-map imaging (VMI) spectrometer [25], where they intersect the CS\(_2\) molecular beam. The pump and probe beams are both linearly polarised in the plane of the VMI detector, perpendicular to the time-of-flight axis. The molecular beam is generated through the expansion of 5% CS\(_2\) in He at 1 bar through a 1 kHz pulsed nozzle (Amsterdam cantilever [26]) with a 100 µm aperture. The resulting expansion passes through a 1 mm skimmer and enters the interaction region of the spectrometer through a hole in the centre of the repeller plate of the VMI spectrometer. The photoelectron spectra are obtained through polar onion-peeling of the background subtracted images [27]. Although the photoelectron angular distributions are obtained, they show no time dependence and as such are not discussed in the results section.

3.2. Experimental results

The 200 nm pump excites a vibrational wavepacket, predominantly in the \(S_2 \ ^1B_2\) excited electronic state. The motion is then probed by non-resonant two-photon absorption at 400 nm. This provides a total energy of 12.5 eV, with the ionisation potential of CS\(_2\) at 10.07 eV. The photoelectron signal obtained when the pump and probe pulse are overlapped in time is plotted in Fig. 4(a) with three main features around 2.1 eV, 1.4 eV and 0.9 eV electron kinetic energy. The spacing between the features is similar to that seen in previous single-photon ionisation measurements [1, 5]. The use of a multiphoton probe maintains a clean experimental measurement, without any probe-pump contributions at early times, while maximising the available energy for ionisation, such that we can observe much of the initial excited state dynamics.

The time-dependence of the photoelectron spectrum is shown in Fig. 4(b). The three peaks in the spectrum have different appearance times, with those at lower electron kinetic energy appearing after those at higher electron kinetic energy. The low-energy feature rises approximately 35 fs after the high-
Figure 4: (a) Photoelectron spectrum obtained when the pump and probe pulses are temporally overlapped. The highlighted regions mark those used in the integrated plots shown in figure 5. (b) Photoelectron spectra as a function of pump-probe delay.
Figure 5: Total integrated photoelectron intensity (a) and intensity within the shaded areas shown in Fig. 4; 0.75 - 1.00 eV (b), 1.20 - 1.60 eV (C) and 1.90 - 2.20 eV (d). Solid lines represent fits to the data. (a,c,d) were fit with a single exponential decay, modulated by a single damped oscillator and convoluted with the instrument response function. However, a second oscillation with a period $\sim 200$ fs is clearly visible in the data. The data in (b) was fit with two damped oscillating components, which reproduce both observed oscillatory features well.
highest energy feature at 2.1 eV. This maps the initial bending motion as seen in the calculated dynamics, Fig. 2b, and in previous measurements [4]. At longer delay times it is also clear that the centre of mass of the photoelectron spectrum shifts to lower electron kinetic energies, such that the lifetime of the measured photoelectron features is longer at lower electron kinetic energies. To obtain a clearer view of the changes observed at the various electron kinetic energies measured, we plot the integrated intensity over the features highlighted in Fig. 4(a) in Fig. 5. The difference in lifetime is apparent in the plots, as well as the appearance of clear oscillations in intensity that peak at times after time zero. The effect of the oscillations is most prominent in the feature centered around 0.9 eV, Fig. 5(b), which has a maximum intensity \( \sim 200 \) fs after excitation. None of the transients can therefore be fit to a simple exponential decay, but are modulated by at least one oscillating component. We therefore fit the transients to an exponential decay modulated by either one or two damped oscillations, convoluted with the instrument response function, corresponding to the laser pulse cross-correlation [28],

\[
g \otimes \left( A_0 \exp\left( -\frac{t - t_0}{\tau} \right) \times \prod_1^n A_n \cos(\omega_n(t - t_0) + \delta_n) \right). \tag{1}
\]

Here \( A_n \) represent intensity scaling parameters, \( t_0 \) the arrival time of the laser pulse, \( \tau \) the exponential lifetime and \( \omega \) and \( \delta \) the angular frequency and phase of the oscillatory component. Fits are plotted as solid lines in Fig. 5. The highest energy feature, Fig. 5(d), provides the clearest data set and contains a single oscillation of period \( \sim 0.9 \) ps, 38 cm\(^{-1}\), as has previously been experimentally observed [9, 10]. This corresponds to the beat between the \( \nu_1 \) and \( \nu_2 \) vibrational modes [8]. This is present in each of the other features in the spectrum, along with a second beat with a period around 200 fs. The effect of this oscillation is clearest in the trace presented in Fig. 5(b), however the mixing with the other oscillation and relatively low contrast makes assigning the absolute value of this oscillation difficult, leading to significant error margins. Nonetheless we extract an oscillation period of \( 220 \) fs, corresponding to \( 149 \) cm\(^{-1}\), from this data. While this oscillation period does not fit with any of the known vibrational periods of the molecule, similar frequencies were also observed in a previous study [4] but were not discussed or assigned. The fits to the experimental data furthermore yield an increase in lifetime towards the lower electron kinetic energy regions. The \( 1/e \) lifetimes extracted are 401 fs, 457 fs and 451 fs for the peaks at 2.1 eV,
1.4 eV and 0.9 eV respectively.

4. Discussion

We now provide a comparison of the experimental measurements and the theoretical calculations. For both the calculations and experiment it is clear that the triplet states play a large role in the dynamics from very early times. Significant population is transferred into the triplet states very rapidly with over 50% of the total population in the triplet states within 250 fs. The overall transfer of population approximately matches the decay rate measured in the experiment such that we are only sensitive to the singlet state population. As mentioned above the initial shift in the measured electron kinetic energy maps the initial bending motion of the molecule. As the pump-probe delay increases, the electron kinetic energy shifts towards lower values, such that we observe a longer lifetime for the lower electron kinetic energy regions in the spectrum in Fig. 5. To compare the measured signal to theory, in Fig. 6 we plot the singlet state component of Fig. 2. Initial excitation leads to a wavepacket that oscillates between linear geometries and an angle of \( \sim 110^\circ \). With increasing pump-probe delay, the range of angles explored narrows and moves away from the linear geometries associated with the spectral feature at the highest electron kinetic energy. Within the current experiment we do not have the time resolution to fully resolve the bending motion, but we do observe the effect of the narrowing and shifting of the angles explored by the molecule as a corresponding narrowing and shifting to lower electron kinetic energies in the photoelectron spectrum.

As mentioned above, the short time period oscillation seen both in the total photoelectron count rate and, in individual regions of the photoelectron spectrum (Fig. 5), do not match any vibrational periods within the molecule. The simulations also show no obvious changes in the molecular structure that appear to provide an explanation for the oscillations. The calculations do show periodic changes in total singlet excited state population that correlate with the observed changes the photoelectron yield. We therefore tentatively assign the oscillations in the experimental spectrum to changes in the total singlet state population.

5. Summary

We have performed a combined theory and experiment study of the excited state dynamics of CS\(_2\). The ab-initio surface-hopping simulations high-
light the importance of the triplet states in the early time dynamics with significant population transfer predicted, and observed in the complementary time-resolved photoelectron spectroscopy measurements. The combined work demonstrates that one can now do on-the-fly dynamics including spin-orbit coupling. The accuracy of the calculation is such that we are able to directly compare the results of the calculation with experiment and explain the shifting and narrowing of the photoelectron spectrum in terms of the bending motion and angles explored by the vibrational wavepacket, while oscillation in the measured photoelectron count rate are explained by the complex coupling of the electronic states that leads to rapid population transfer between manifolds of multiple singlet and triplet excited states.

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