Measurement of CP asymmetry in B-s(0) -> D-s(-/+ ) K--/+ decays
Measurement of $CP$ asymmetry in $B_s^0 \rightarrow D_s^\mp K^{\pm}$ decays

The LHCb collaboration

E-mail: Manuel.Schiller@cern.ch

Abstract: We report on measurements of the time-dependent $CP$ violating observables in $B_s^0 \rightarrow D_s^\mp K^{\pm}$ decays using a dataset corresponding to 1.0 fb$^{-1}$ of $pp$ collisions recorded with the LHCb detector. We find the $CP$ violating observables $C_f = 0.53\pm0.25\pm0.04$, $A_f^{\Delta\Gamma} = 0.37\pm0.42\pm0.20$, $A_f^{\Delta\Gamma} = 0.20\pm0.41\pm0.20$, $S_f = -1.09\pm0.33\pm0.08$, $S_f = -0.36 \pm 0.34 \pm 0.08$, where the uncertainties are statistical and systematic, respectively. Using these observables together with a recent measurement of the $B_s^0$ mixing phase $-2\beta_s$ leads to the first extraction of the CKM angle $\gamma$ from $B_s^0 \rightarrow D_s^\mp K^{\pm}$ decays, finding $\gamma = (115^{+28}_{-43})^\circ$ modulo 180$^\circ$ at 68% CL, where the error contains both statistical and systematic uncertainties.

Keywords: CP violation, CKM angle gamma, B physics, Flavor physics, Hadron-Hadron Scattering

ArXiv ePrint: 1407.6127
1 Introduction

Time-dependent analyses of tree-level $B_{(s)}^0 \rightarrow D_{(s)}^{\mp} \pi^\pm, K^\pm$ decays\footnote{Inclusion of charge conjugate modes is implied except where explicitly stated.} are sensitive to the angle $\gamma \equiv \arg \left( -V_{ud} V_{ub}^* / V_{cd} V_{cb}^* \right)$ of the unitarity triangle of the Cabibbo-Kobayashi-Maskawa (CKM) matrix \cite{1,2} through $CP$ violation in the interference of mixing and decay amplitudes \cite{3–5}. The determination of $\gamma$ from such tree-level decays is important because it is not sensitive to potential effects from most models of physics beyond the Standard

---

The LHCb collaboration

---

The LHCb collaboration

---

1 Introduction

Time-dependent analyses of tree-level $B_{(s)}^0 \rightarrow D_{(s)}^{\mp} \pi^\pm, K^\pm$ decays\footnote{Inclusion of charge conjugate modes is implied except where explicitly stated.} are sensitive to the angle $\gamma \equiv \arg \left( -V_{ud} V_{ub}^* / V_{cd} V_{cb}^* \right)$ of the unitarity triangle of the Cabibbo-Kobayashi-Maskawa (CKM) matrix \cite{1,2} through $CP$ violation in the interference of mixing and decay amplitudes \cite{3–5}. The determination of $\gamma$ from such tree-level decays is important because it is not sensitive to potential effects from most models of physics beyond the Standard

---

The LHCb collaboration

---

1 Introduction

Time-dependent analyses of tree-level $B_{(s)}^0 \rightarrow D_{(s)}^{\mp} \pi^\pm, K^\pm$ decays\footnote{Inclusion of charge conjugate modes is implied except where explicitly stated.} are sensitive to the angle $\gamma \equiv \arg \left( -V_{ud} V_{ub}^* / V_{cd} V_{cb}^* \right)$ of the unitarity triangle of the Cabibbo-Kobayashi-Maskawa (CKM) matrix \cite{1,2} through $CP$ violation in the interference of mixing and decay amplitudes \cite{3–5}. The determination of $\gamma$ from such tree-level decays is important because it is not sensitive to potential effects from most models of physics beyond the Standard

---

The LHCb collaboration

---

1 Introduction

Time-dependent analyses of tree-level $B_{(s)}^0 \rightarrow D_{(s)}^{\mp} \pi^\pm, K^\pm$ decays\footnote{Inclusion of charge conjugate modes is implied except where explicitly stated.} are sensitive to the angle $\gamma \equiv \arg \left( -V_{ud} V_{ub}^* / V_{cd} V_{cb}^* \right)$ of the unitarity triangle of the Cabibbo-Kobayashi-Maskawa (CKM) matrix \cite{1,2} through $CP$ violation in the interference of mixing and decay amplitudes \cite{3–5}. The determination of $\gamma$ from such tree-level decays is important because it is not sensitive to potential effects from most models of physics beyond the Standard

---

The LHCb collaboration
In contrast to the CP violating observables in these decays are functions of a combination of \( \gamma \) and the relevant mixing phase, namely \( \gamma + 2\beta \) (\( \beta \equiv \text{arg}(\mathcal{V}_{cb}\mathcal{V}_{us}^{\ast})/\text{arg}(\mathcal{V}_{ub}\mathcal{V}_{cs}^{\ast})) \) in the \( B^0 \) and \( \gamma - 2\beta_s \) (\( \beta_s \equiv \text{arg}(-\mathcal{V}_b\mathcal{V}_{sb}^{\ast})/\text{arg}(\mathcal{V}_c\mathcal{V}_{tb}^{\ast})) \) in the \( B_s^0 \) system. A measurement of these physical observables can therefore be interpreted in terms of \( \gamma \) or \( \beta_s \) by using an independent measurement of the other parameter as input.

Such measurements have been performed by both the BaBar [6, 7] and the Belle [8, 9] collaborations using \( B^0 \to D_s^{(\pm)}\pi^\mp \) decays. In these decays, however, the ratios \( r_{D_s^{(\pm)}\pi} = |A(B^0 \to D_s^{(\pm)}\pi^\mp)/A(B^0 \to D_s^{(\pm)}\pi^-)| \) between the interfering \( b \to u \) and \( b \to c \) amplitudes are small, \( r_{D_s^{(\pm)}\pi} \approx 0.02 \), limiting the sensitivity on \( \gamma \) [10].

The leading order Feynman diagrams contributing to the interference of decay and mixing in \( B_s^0 \to D_s^+K^- \) are shown in figure 1. In contrast to \( B^0 \to D^{(\pm)}\pi^\mp \) decays, here both the \( B_s^0 \to D_s^-K^+ \) (\( b \to cs\bar{u} \)) and \( B_s^0 \to D_s^+K^- \) (\( b \to u\bar{c}s \)) amplitudes are of the same order in the sine of the Cabibbo angle \( \lambda = 0.2252 \pm 0.0007 \) [11, 12], \( O(\lambda^3) \), and the amplitude ratio of the interfering diagrams is approximately \( |V_{ub}V_{cs}/V_{cb}V_{us}| \approx 0.4 \). Moreover, the decay width difference in the \( B_s^0 \) system, \( \Delta \Gamma_s \), is nonzero [13], which allows a determination of \( \gamma - 2\beta_s \) from the sinusoidal and hyperbolic terms in the decay time evolution, up to a two-fold ambiguity.

This paper presents the first measurements of the \( CP \) violating observables in \( B_s^0 \to D_s^+K^- \) decays using a dataset corresponding to 1.0 fb\(^{-1}\) of \( pp \) collisions recorded with the LHCb detector at \( \sqrt{s} = 7 \) TeV, and the first determination of \( \gamma - 2\beta_s \) in these decays.

1.1 Decay rate equations and \( CP \) violation observables

The time-dependent decay rates of the initially produced flavour eigenstates \( |B_s^0(t = 0)\rangle \) and \( |B_s^0(t = 0\rangle \) are given by

\[
\frac{d\Gamma_{B_s^0 \to f}(t)}{dt} = \frac{1}{2} |A_f|^2 (1 + |\lambda_f|^2)e^{-\Gamma_s t}\left[ \cosh\left( \frac{\Delta \Gamma_s t}{2} \right) + A_f^\Delta \Gamma \sinh\left( \frac{\Delta \Gamma_s t}{2} \right) \right] + C_f \cos(\Delta m_s t) - S_f \sin(\Delta m_s t),
\]

(1.1)
where \( \lambda_f \equiv (q/p)(\bar{A}_f/A_f) \) and \( A_f (\bar{A}_f) \) is the decay amplitude of a \( B^0_s \) to decay to a final state \( f (\bar{f}) \). \( \Gamma_s \) is the average \( B^0_s \) decay width, and \( \Delta \Gamma_s \) is the positive [14] decay-width difference between the heavy and light mass eigenstates in the \( B^0_s \) system. The complex coefficients \( p \) and \( q \) relate the \( B^0_s \) meson mass eigenstates, \( |B^0_s^{L,H}\rangle \), to the flavour eigenstates, \( \left| B^0_s \right\rangle \) and \( \left| \bar{B}^0_s \right\rangle \)

\[
|B_{L,H}\rangle = p|B^0_s^{L,H}\rangle + q|\bar{B}^0_s^{L,H}\rangle,
\]

with \( |p|^2 + |q|^2 = 1 \). Similar equations can be written for the \( CP \)-conjugate decays replacing \( C_f \) by \( C_{\bar{f}} \), \( S_f \) by \( S_{\bar{f}} \), and \( A_f^{\Delta \Gamma} \) by \( A_{\bar{f}}^{\Delta \Gamma} \). In our convention \( f \) is the \( D^+_s K^- \) final state and \( \bar{f} \) is \( D^+_s K^- \). The \( CP \) asymmetry observables \( C_f, S_f, A_f^{\Delta \Gamma}, C_{\bar{f}}, S_{\bar{f}} \) and \( A_{\bar{f}}^{\Delta \Gamma} \) are given by

\[
\begin{align*}
C_f &= \frac{1 - |\lambda_f|^2}{1 + |\lambda_f|^2} = -C_\bar{f} = \frac{1 - |\lambda_\bar{f}|^2}{1 + |\lambda_\bar{f}|^2}, \\
S_f &= \frac{2Im(\lambda_f)}{1 + |\lambda_f|^2}, \quad A_f^{\Delta \Gamma} = \frac{-2Re(\lambda_f)}{1 + |\lambda_f|^2}, \\
S_{\bar{f}} &= \frac{2Im(\lambda_\bar{f})}{1 + |\lambda_\bar{f}|^2}, \quad A_{\bar{f}}^{\Delta \Gamma} = \frac{-2Re(\lambda_\bar{f})}{1 + |\lambda_\bar{f}|^2}.
\end{align*}
\]

The equality \( C_f = -C_\bar{f} \) results from \( |q/p| = 1 \) and \( |\lambda_f| = |\lambda_\bar{f}| \), i.e. the assumption of no \( CP \) violation in either the decay or mixing amplitudes. The \( CP \) observables are related to the magnitude of the amplitude ratio \( r_{D_sK} \equiv |\lambda_{D_sK}| = |A(\bar{B}^0_s \to D_s^- K^+)/A(B^0_s \to D_s^- K^+)| \), the strong phase difference \( \delta \), and the weak phase difference \( \gamma - 2\beta_s \) by the following equations:

\[
\begin{align*}
C_f &= \frac{1 - r_{D_sK}^2}{1 + r_{D_sK}^2}, \\
A_f^{\Delta \Gamma} &= \frac{-2r_{D_sK} \cos(\delta - (\gamma - 2\beta_s))}{1 + r_{D_sK}^2}, \quad A_{\bar{f}}^{\Delta \Gamma} = \frac{-2r_{D_sK} \cos(\delta + (\gamma - 2\beta_s))}{1 + r_{D_sK}^2}, \\
S_f &= \frac{2r_{D_sK} \sin(\delta - (\gamma - 2\beta_s))}{1 + r_{D_sK}^2}, \quad S_{\bar{f}} = \frac{-2r_{D_sK} \sin(\delta + (\gamma - 2\beta_s))}{1 + r_{D_sK}^2}.
\end{align*}
\]

### 1.2 Analysis strategy

To measure the \( CP \) violating observables defined in section 1.1, it is necessary to perform a fit to the decay-time distribution of the selected \( B^0_s \to D_s^+ K^\pm \) candidates. The kinematically similar mode \( B^0_s \to D^- s \pi^+ \) is used as control channel which helps in the determination of the time-dependent efficiency and flavour tagging performance. Before a fit to the decay time can be performed, it is necessary to distinguish the signal and background candidates in the selected sample. This analysis uses three variables to maximise sensitivity.
when discriminating between signal and background: the $B_s^0$ mass; the $D_{s}^{-}$ mass; and the log-likelihood difference $L(K/\pi)$ between the pion and kaon hypotheses for the companion particle.

In section 4, the signal and background shapes needed for the analysis are obtained in each of the variables. Section 5 describes how a simultaneous extended maximum likelihood fit (in the following referred to as multivariate fit) to these three variables is used to determine the yields of signal and background components in the samples of $B_s^0 \rightarrow D_{s}^{-} \pi^+$ and $B_s^0 \rightarrow D_{s}^{+} K^\pm$ candidates. Section 6 describes how to obtain the flavour at production of the $B_s^0 \rightarrow D_{s}^{\mp} K^\pm$ candidates using a combination of flavour-tagging algorithms, whose performance is calibrated with data using flavour-specific control modes. The decay-time resolution and acceptance are determined using a mixture of data control modes and simulated signal events, described in section 7.

Finally, section 8 describes two approaches to fit the decay-time distribution of the $B_s^0 \rightarrow D_{s}^{\mp} K^\pm$ candidates which extract the $CP$ violating observables. The first fit, henceforth referred to as the $sFit$, uses the results of the multivariate fit to obtain the so-called $sWeights$ [15] which allow the background components to be statistically subtracted [16]. The $sFit$ to the decay-time distribution is therefore performed using only the probability density function (PDF) of the signal component. The second fit, henceforth referred to as the $cFit$, uses the various shapes and yields of the multivariate fit result for the different signal and background components. The $cFit$ subsequently performs a six-dimensional maximum likelihood fit to these variables, the decay-time distribution and uncertainty, and the probability that the initial $B_s^0$ flavour is correctly determined, in which all contributing signal and background components are described with their appropriate PDFs. In section 10, we extract the CKM angle $\gamma$ using the result of one of the two approaches.

2 Detector and software

The LHCb detector [17] is a single-arm forward spectrometer covering the pseudorapidity range $2 < \eta < 5$, designed for the study of particles containing $b$ or $c$ quarks. The detector includes a high-precision tracking system consisting of a silicon-strip vertex detector surrounding the $pp$ interaction region [18], a large-area silicon-strip detector located upstream of a dipole magnet with a bending power of about 4 Tm, and three stations of silicon-strip detectors and straw drift tubes [19] placed downstream of the magnet. The tracking system provides a measurement of momentum, $p$, with a relative uncertainty that varies from 0.4% at low momentum to 0.6% at 100 GeV/c. The minimum distance of a track to a primary $pp$ collision vertex, the impact parameter, is measured with a resolution of $(15 + 29/p_T) \mu m$, where $p_T$ is the component of $p$ transverse to the beam, in GeV/c. Different types of charged hadrons are distinguished using information from two ring-imaging Cherenkov detectors [20]. The magnet polarity is reversed regularly to control systematic effects.

The trigger [21] consists of a hardware stage, based on information from the calorimeter and muon systems, followed by a software stage, which applies a full event reconstruction. The software trigger requires a two-, three- or four-track secondary vertex with a large sum of the transverse momentum of the charged particles and a significant displacement from
the primary \( pp \) interaction vertices (PVs). A multivariate algorithm [22] is used for the identification of secondary vertices consistent with the decay of a \( b \) hadron.

In the simulation, \( pp \) collisions are generated using PYTHIA [23] with a specific LHCb configuration [24]. Decays of hadrons are described by EvtGen [25], in which final state radiation is generated using PHOTOS [26]. The interaction of the generated particles with the detector and its response are implemented using the GEANT4 toolkit [27, 28] as described in ref. [29].

### 3 Event selection

The event selection begins by building \( D_s^- \to K^- K^+ \pi^- \), \( D_s^- \to K^- \pi^+ \pi^- \), and \( D_s^- \to \pi^- \pi^+ \pi^- \) candidates from reconstructed charged particles. These \( D_s^- \) candidates are subsequently combined with a fourth particle, referred to as the “companion”, to form \( B_s^0 \rightarrow D_s^\mp K^\pm \) and \( B_s^0 \rightarrow D_s^- \pi^+ \) candidates. The flavour-specific Cabibbo-favoured decay mode \( B_s^0 \rightarrow D_s^- \pi^+ \) is used as a control channel in the analysis, and is selected identically to \( B_s^0 \rightarrow D_s^0 K^\pm \) except for the PID criteria on the companion particle. The decay-time and \( B_s^0 \) mass resolutions are improved by performing a kinematic fit [30] in which the \( B_s^0 \) candidate is constrained to originate from its associated proton-proton interaction, i.e. the one with the smallest IP with respect to the \( B_s^0 \) candidate, and the \( B_s^0 \) mass is computed with a constraint on the \( D_s^- \) mass.

The \( B_s^0 \rightarrow D_s^- \pi^+ \) mode is used for the optimisation of the selection and for studying and constraining physics backgrounds to the \( B_s^0 \rightarrow D_s^0 K^\pm \) decay. The \( B_s^0 \rightarrow D_s^0 K^\pm \) and \( B_s^0 \rightarrow D_s^- \pi^+ \) candidates are required to be matched to the secondary vertex candidates found in the software trigger. Subsequently, a preselection is applied to the \( B_s^0 \rightarrow D_s^0 K^\pm \) and \( B_s^0 \rightarrow D_s^- \pi^+ \) candidates using a similar multivariate displaced vertex algorithm to the trigger selection, but with offline-quality reconstruction.

A selection using the gradient boosted decision tree (BDTG) [31] implementation in the TMVA software package [32] further suppresses combinatorial backgrounds. The BDTG is trained on data using the \( B_s^0 \rightarrow D_s^- \pi^+ \), \( D_s^- \rightarrow K^- K^+ \pi^- \) decay sample, which is purified with respect to the previous preselection exploiting PID information from the Cherenkov detectors. Since all channels in this analysis are kinematically similar, and since no PID information is used as input to the BDTG, the resulting BDTG performs equally well on the other \( D_s^- \) decay modes. The optimal working point is chosen to maximise the expected sensitivity to the CP violating observables in \( B_s^0 \rightarrow D_s^0 K^\pm \) decays. In addition, the \( B_s^0 \) and \( D_s^- \) candidates are required to be within \( m(B_s^0) \in [5300, 5800] \text{MeV}/c^2 \) and \( m(D_s^-) \in [1930, 2015] \text{MeV}/c^2 \), respectively.

Finally, the different final states are distinguished by using PID information. This selection also strongly suppresses cross-feed and peaking backgrounds from other misidentified decays of \( b\)-hadrons to \( c\)-hadrons. We will refer to such backgrounds as “fully reconstructed” if no particles are missed in the reconstruction, and “partially reconstructed” otherwise. The decay modes \( B^0 \rightarrow D^- \pi^+ \), \( B^0 \rightarrow D_s^- \pi^+ \), \( \bar{B}_s^0 \rightarrow \bar{D}_s^- \pi^+ \), \( B_s^0 \rightarrow D_s^0 K^\pm \), and \( B_s^0 \rightarrow D_s^- \pi^+ \) are backgrounds to \( B_s^0 \rightarrow D_s^- \pi^+ \), while \( B_s^0 \rightarrow D_s^- \pi^+ \), \( B_s^0 \rightarrow D_s^- \pi^+ \), \( B_s^0 \rightarrow D_s^- \rho^+ \), \( B_s^0 \rightarrow D_s^- K^+ \), \( B_s^0 \rightarrow D^- K^+ \), \( B_s^0 \rightarrow D^- \pi^+ \), \( \bar{B}_s^0 \rightarrow \bar{D}_s^- K^+ \), \( \bar{B}_s^0 \rightarrow \bar{D}_s^- \pi^+ \), and
$A_b^0 \rightarrow D_s^{(*)-}p$ are backgrounds to $B_s^0 \rightarrow D_s^+K^\mp$. This part of the selection is necessarily different for each $D_s^-$ decay mode, as described below.

- For $D_s^- \rightarrow \pi^-\pi^+\pi^-$ none of the possible misidentified backgrounds fall inside the $D_s^-$ mass window. Loose PID requirements are nevertheless used to identify the $D_s^-$ decay products as pions in order to suppress combinatorial background.

- For $D_s^- \rightarrow K^-\pi^+\pi^-$, the relevant peaking backgrounds are $\Lambda_c^- \rightarrow p\pi^+\pi^-$ in which the antiproton is misidentified, and $D^- \rightarrow K^+\pi^-\pi^-$ in which both the kaon and a pion are misidentified. As this is the smallest branching fraction $D_s^-$ decay mode used, and hence that most affected by background, all $D_s^-$ decay products are required to pass tight PID requirements.

- The $D_s^- \rightarrow K^- K^+ \pi^-$ mode is split into three submodes. We distinguish between the resonant $D_s^- \rightarrow \phi\pi^-$ and $D_s^- \rightarrow K^{*0}K^-$ decays, and the remaining decays. Candidates in which the $K^+K^-$ pair falls within 20 MeV$/c^2$ of the $\phi$ mass are identified as a $D_s^- \rightarrow \phi\pi^-$ decay. This requirement suppresses most of the cross-feed and combinatorial background, and only loose PID requirements are needed. Candidates within a 50 MeV$/c^2$ window around the $K^{*0}$ mass are identified as a $D_s^- \rightarrow K^{*0}K^-$ decay; it is kinematically impossible for a candidate to satisfy both this and the $\phi$ requirement. In this case there is non-negligible background from misidentified $D^- \rightarrow K^+\pi^-\pi^-$ and $\Lambda_c^- \rightarrow \bar{p}\pi^-K^+$ decays which are suppressed through tight PID requirements on the $D_s^-$ kaon with the same charge as the $D_s^-$ pion. The remaining candidates, referred to as nonresonant decays, are subject to tight PID requirements on all decay products to suppress cross-feed backgrounds.

Figure 2 shows the relevant mass distributions for candidates passing and failing this PID selection. Finally a loose PID requirement is made on the companion track. After all selection requirements, fewer than 2% of retained events contain more than one signal candidate. All candidates are used in the subsequent analysis.

4 Signal and background shapes

The signal and background shapes are obtained using a mixture of data-driven approaches and simulation. The simulated events need to be corrected for kinematic differences between simulation and data, as well as for the kinematics-dependent efficiency of the PID selection requirements. In order to obtain kinematic distributions in data for this weighting, we use the decay mode $B^0 \rightarrow D^-\pi^+$, which can be selected with very high purity without the use of any PID requirements and is kinematically very similar to the $B_s^0$ signals. The PID efficiencies are measured as a function of particle momentum and event occupancy using prompt $D^{*+} \rightarrow D^0(K^-\pi^+)\pi^+$ decays which provide pure samples of pions and kaons [33], henceforth called $D^{*+}$ calibration sample.
Figure 2. Mass distributions for $D_s^-$ candidates passing (black, open circles) and failing (red, crosses) the PID selection criteria. In reading order: $D_s^- \to K^- K^+ \pi^-$, $D_s^- \to \pi^- \pi^+ \pi^-$, and $D_s^- \to K^- \pi^+ \pi^-$. 

4.1 $B_s^0$ candidate mass shapes

In order to model radiative and reconstruction effects, the signal shape in the $B_s^0$ mass is the sum of two Crystal Ball [34] functions with common mean and oppositely oriented tails. The signal shapes are determined separately for $B_s^0 \to D_s^\mp K^\pm$ and $B_s^0 \to D_s^- \pi^+$ from simulated candidates. The shapes are subsequently fixed in the multivariate fit except for the common mean of the Crystal Ball functions which floats for both the $B_s^0 \to D_s^- \pi^+$ and $B_s^0 \to D_s^\mp K^\pm$ channel.

The functional form of the combinatorial background is taken from the upper $B_s^0$ sideband, with its parameters left free to vary in the subsequent multivariate fit. Each $D_s^-$ mode is considered independently and parameterised by either an exponential function or by a combination of an exponential and a constant function.

The shapes of the fully or partially reconstructed backgrounds are fixed from simulated events using a non-parametric kernel estimation method (KEYS, [35]). Exceptions to this are the $B^0 \to D^- \pi^+$ background in the $B_s^0 \to D^- \pi^+$ fit and the $B_s^0 \to D^- \pi^+$ background in the $B_s^0 \to D_s^\mp K^\pm$ fit, which are obtained from data. The latter two backgrounds are
reconstructed with the “wrong” mass hypothesis but without PID requirements, which
would suppress them. The resulting shapes are then weighted to account for the effect
of the momentum-dependent efficiency of the PID requirements from the $D^{*+}$ calibration
samples, and KEYS templates are extracted for use in the multivariate fit.

4.2 $D_s^-$ candidate mass shapes

The signal shape in the $D_s^-$ mass is again a sum of two Crystal Ball functions with common
mean and oppositely oriented tails. The signal shapes are extracted separately for each
$D_s^-$ decay mode from simulated events that have the full selection chain applied to them.
The shapes are subsequently fixed in the multivariate fit except for the common mean of
the Crystal Ball functions, which floats independently for each $D_s^-$ decay mode.

The combinatorial background consists of both random combinations of tracks which
do not peak in the $D_s^-$ mass, and, in some $D_s^-$ decay modes, backgrounds that contain a
true $D_s^-$, and a random companion track. It is parameterised separately for each $D_s^-$ decay
mode either by an exponential function or by a combination of an exponential function and
the signal $D_s^-$ shape.

The fully and partially reconstructed backgrounds which contain a correctly recon-
structed $D_s^-$ candidate ($B_s^0 \to D_s^{\mp} K^\pm$ and $B^0 \to D_s^- \pi^+$ as backgrounds in the $B_s^0 \to D_s^- \pi^+$
fit; $B^0 \to D_s^- K^+$ and $B^0 \to D_s^- \pi^+$ as backgrounds in the $B_s^0 \to D_s^{\mp} K^\pm$ fit) are assumed
to have the same mass distribution as the signal. For other backgrounds, the shapes are
KEYS templates taken from simulated events, as in the $B_s^0$ mass.

4.3 Companion $L(K/\pi)$ shapes

We obtain the PDFs describing the $L(K/\pi)$ distributions of pions and kaons from dedicated
$D^{*+}$ calibration samples. We obtain the PDF describing the protons using a calibration
sample of $\Lambda_c^+ \to pK^- \pi^+$ decays. These samples are weighted to match the signal kinematic
and event occupancy distributions in the same way as the simulated events. The weighting
is done separately for each signal and background component, as well as for each magnet
polarity. The shapes for each magnet polarity are subsequently combined according to the
integrated luminosity in each sample.

The signal companion $L(K/\pi)$ shape is obtained separately for each $D_s^-$ decay mode
to account for small kinematic differences between them. The combinatorial background
companion $L(K/\pi)$ shape is taken to consist of a mixture of pions, protons, and kaons, and
its normalisation is left floating in the multivariate fit. The companion $L(K/\pi)$ shape for
fully or partially reconstructed backgrounds is obtained by weighting the PID calibration
samples to match the event distributions of simulated events, for each background type.

5 Multivariate fit to $B_s^0 \to D_s^{\mp} K^\pm$ and $B_s^0 \to D_s^- \pi^+$

The total PDF for the multivariate fit is built from the product of the signal and back-
ground PDFs, since correlations between the fitting variables are measured to be small in
simulation. These product PDFs are then added for each $D_s^-$ decay mode, and almost all
background yields are left free to float. The only exceptions are those backgrounds whose
yield is below 2% of the signal yield. These are $B^0 \to D^- K^+$, $B^0 \to D^- \pi^+$, $\Lambda_b^0 \to \Lambda_c^- K^+$, and $\Lambda_b^0 \to \Lambda_c^- \pi^+$ for the $B^0_s \to D_s^+ K^±$ fit, and $B^0 \to D^- \pi^+$, $\Lambda_b^0 \to \Lambda_c^- \pi^+$, and $B^0_s \to D_s^± K^±$ for the $B^0_s \to D_s^- \pi^+$ fit. These background yields are fixed from known branching fractions and relative efficiencies measured using simulated events. The multivariate fit results in a signal yield of 28 260 $\pm$ 180 $B^0_s \to D_s^- \pi^+$ and 1770 $\pm$ 50 $B^0_s \to D_s^± K^±$ decays, with an effective purity of 85% for $B^0_s \to D_s^- \pi^+$ and 74% for $B^0_s \to D_s^± K^±$. The multivariate fit is checked for biases using large samples of data-like pseudoexperiments, and none is found. The results of the multivariate fit are shown in figure 3 for both the $B^0_s \to D_s^- \pi^+$ and $B^0_s \to D_s^± K^±$, summed over all $D_s^-$ decay modes.

6 Flavour tagging

At the LHC, $b$ quarks are produced in pairs $b\bar{b}$; one of the two hadronises to form the signal $B^0_s$, the other $b$ quark hadronises and decays independently. The identification of the $B^0_s$ initial flavour is performed by means of two flavour-tagging algorithms which exploit this pair-wise production of $b$ quarks. The opposite side (OS) tagger determines the flavour of the non-signal $b$-hadron produced in the proton-proton collision using the charge of the lepton ($\mu$, $e$) produced in semileptonic $B$ decays, or that of the kaon from the $b \to c \to s$ decay chain, or the charge of the inclusive secondary vertex reconstructed from $b$-decay products. The same side kaon (SSK) tagger searches for an additional charged kaon accompanying the fragmentation of the signal $B^0_s$ or $\bar{B}^0_s$.

Each of these algorithms has an intrinsic mistag rate $\omega = \frac{\text{wrong tags}}{\text{all tags}}$ and a tagging efficiency $\varepsilon_{\text{tag}} = \frac{\text{tagged candidates}}{\text{all candidates}}$. Candidates can be tagged incorrectly due to tracks from the underlying event, particle misidentifications, or flavour oscillations of neutral $B$ mesons. The intrinsic mistag $\omega$ can only be measured in flavour-specific, self-tagging final states.

The tagging algorithms predict for each $B^0_s$ candidate an estimate $\eta$ of the mistag probability, which should closely follow the intrinsic mistag $\omega$. This estimate $\eta$ is obtained by using a neural network trained on simulated events whose inputs are the kinematic, geometric, and PID properties of the tagging particle(s).

The estimated mistag $\eta$ is treated as a per-candidate variable, thus adding an observable to the fit. Due to variations in the properties of tagging tracks for different channels, the predicted mistag probability $\eta$ is usually not exactly the (true) mistag rate $\omega$, which requires $\eta$ to be calibrated using flavour specific, and therefore self-tagging, decays. The statistical uncertainty on $C_f$, $S_f$, and $S_f$ scales with $\frac{1}{\sqrt{\varepsilon_{\text{eff}}}}$, defined as $\varepsilon_{\text{eff}} = \varepsilon_{\text{tag}}(1 - 2\omega)^2$. Therefore, the tagging algorithms are tuned for maximum effective tagging power $\varepsilon_{\text{eff}}$.

6.1 Tagging calibration

The calibration for the OS tagger is performed using several control channels: $B^+ \to J/\psi K^+$, $B^+ \to D^0 \pi^+$, $B^0 \to D^{*+} \mu^+ \nu_\mu$, $B^0 \to J/\psi K^{*0}$ and $B^0 \to D_s^- \pi^+$. This calibration of $\eta$ is done for each control channel using the linear function

$$\omega = p_0 + p_1 \cdot (\eta - \langle \eta \rangle), \quad (6.1)$$
Figure 3. The multivariate fit to the (left) $B_s^0 \rightarrow D_s^- \pi^+$ and (right) $B_s^0 \rightarrow D_s^{\ast} K^\pm$ candidates for all $D_s^-$ decay modes combined. From top to bottom: distributions of candidates in $B_s^0$ mass, $D_s^-$ mass, companion PID log-likelihood difference. The solid, blue, line represents the sum of the fit components.
where the values of $p_0$ and $p_1$ are called calibration parameters, and $\langle \eta \rangle$ is the mean of the $\eta$ distribution predicted by a tagger in a specific control channel. Systematic uncertainties are assigned to account for possible dependences of the calibration parameters on the final state considered, on the kinematics of the $B_s^0$ candidate and on the event properties. The corresponding values of the calibration parameters are summarised in table 1. For each control channel the relevant calibration parameters are reported with their statistical and systematic uncertainties. These are averaged to give the reference values including a systematic uncertainty accounting for kinematic differences between different channels. The resulting calibration parameters for the $B_s^0 \rightarrow D_s^\mp \eta \pi^+$ fit are: $p_0 = 0.3834 \pm 0.0014 \pm 0.0040$ and $p_1 = 0.972\pm0.012\pm0.035$, where the $p_0$ for each control channel needs to be translated to the $\langle \eta \rangle$ of $B_s^0 \rightarrow D_s^- \pi^+$, the channel which is most similar to the signal channel $B_s^0 \rightarrow D_s^+ \eta \pi^+$. This is achieved by the transformation $p_0 \rightarrow p_0 + p_1(\langle \eta \rangle - 0.3813)$ in each control channel.

<table>
<thead>
<tr>
<th>Control channel</th>
<th>$\langle \eta \rangle$</th>
<th>$p_0 - \langle \eta \rangle$</th>
<th>$p_1$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$B^+ \rightarrow J/\psi K^+$</td>
<td>0.3919</td>
<td>0.0008 $\pm$ 0.0014 $\pm$ 0.0015</td>
<td>0.982 $\pm$ 0.017 $\pm$ 0.005</td>
</tr>
<tr>
<td>$B^+ \rightarrow D^0 \pi^+$</td>
<td>0.3836</td>
<td>0.0018 $\pm$ 0.0016 $\pm$ 0.0015</td>
<td>0.972 $\pm$ 0.017 $\pm$ 0.005</td>
</tr>
<tr>
<td>$B^0 \rightarrow J/\psi K^{*0}$</td>
<td>0.390</td>
<td>0.0090 $\pm$ 0.0030 $\pm$ 0.0060</td>
<td>0.882 $\pm$ 0.043 $\pm$ 0.039</td>
</tr>
<tr>
<td>$B^0 \rightarrow D^{*-} \mu^+ \nu_\mu$</td>
<td>0.3872</td>
<td>0.0081 $\pm$ 0.0019 $\pm$ 0.0069</td>
<td>0.946 $\pm$ 0.019 $\pm$ 0.061</td>
</tr>
<tr>
<td>$B_s^0 \rightarrow D_s^- \pi^+$</td>
<td>0.3813</td>
<td>0.0159 $\pm$ 0.0097 $\pm$ 0.0071</td>
<td>1.000 $\pm$ 0.116 $\pm$ 0.047</td>
</tr>
<tr>
<td>Average</td>
<td>0.3813</td>
<td>0.0021 $\pm$ 0.0014 $\pm$ 0.0040</td>
<td>0.972 $\pm$ 0.012 $\pm$ 0.035</td>
</tr>
</tbody>
</table>

Table 1. Calibration parameters of the combined OS tagger extracted from different control channels. In each entry the first uncertainty is statistical and the second systematic.

The SSK algorithm uses a neural network to select fragmentation particles, giving improved flavour tagging power [36] with respect to earlier cut-based [37] algorithms. It is calibrated using the $B_s^0 \rightarrow D_s^- \pi^+$ channel, resulting in $\langle \eta \rangle = 0.4097$, $p_0 = 0.4244 \pm 0.0086 \pm 0.0071$ and $p_1 = 1.255 \pm 0.140 \pm 0.104$, where the first uncertainty is statistical and second systematic. The systematic uncertainties include the uncertainty on the decay-time resolution, the $B_s^0 \rightarrow D_s^- \pi^+$ fit model, and the backgrounds in the $B_s^0 \rightarrow D_s^- \pi^+$ fit.

Figure 4 shows the measured mistag probability as a function of the mean predicted mistag probability in $B_s^0 \rightarrow D_s^- \pi^+$ decays for the OS and SSK taggers. The data points show a linear correlation corresponding to the functional form in eq. (6.1). We additionally validate that the obtained tagging calibration parameters can be used in $B_s^0 \rightarrow D_s^- \pi^+$ decays by comparing them for $B_s^0 \rightarrow D_s^+ \eta \pi^+$ and $B_s^0 \rightarrow D_s^- \pi^+$ in simulated events; we find excellent agreement between the two. We also evaluate possible tagging asymmetries between $B$ and $\bar{B}$ mesons for the OS and SSK taggers by performing the calibrations split by $B$ meson flavour. The OS tagging asymmetries are measured using $B^+ \rightarrow J/\psi K^+$ decays, while the SSK tagging asymmetries are measured using prompt $D^\pm$ mesons whose $p_T$ distribution has been weighted to match the $B_s^0 \rightarrow D_s^- \pi^+$ signal. The resulting initial flavour asymmetries for $p_0$, $p_1$ and $\varepsilon_{\text{tag}}$ are taken into account in the decay-time fit.
Figure 4. Measured mistag rate against the average predicted mistag rate for the (left) OS and (right) SSK taggers in $B^0_s \to D_s^- \pi^+$ decays. The error bars represent only the statistical uncertainties. The solid curve is the linear fit to the data points, the shaded area defines the 68% confidence level region of the calibration function (statistical only).

<table>
<thead>
<tr>
<th>Event type</th>
<th>$\varepsilon_{\text{tag}}$ [%]</th>
<th>$\varepsilon_{\text{eff}}$ [%]</th>
</tr>
</thead>
<tbody>
<tr>
<td>OS-only</td>
<td>$19.80 \pm 0.23$</td>
<td>$1.61 \pm 0.03 \pm 0.08$</td>
</tr>
<tr>
<td>SSK-only</td>
<td>$28.85 \pm 0.27$</td>
<td>$1.31 \pm 0.22 \pm 0.17$</td>
</tr>
<tr>
<td>OS-SSK</td>
<td>$18.88 \pm 0.23$</td>
<td>$2.15 \pm 0.05 \pm 0.09$</td>
</tr>
<tr>
<td>Total</td>
<td>$67.53$</td>
<td>$5.07$</td>
</tr>
</tbody>
</table>

Table 2. Flavour tagging performance for the three different tagging categories for $B^0_s \to D_s^- \pi^+$ candidates.

6.2 Combination of OS and SSK taggers

Since the SSK and OS taggers rely on different physical processes they are largely independent, with a correlation measured as negligible. The tagged candidates are therefore split into three different samples depending on the tagging decision: events only tagged by the OS tagger (OS-only), those only tagged by the SSK tagger (SSK-only), and those tagged by both the OS and SSK taggers (OS-SSK). For the candidates that have decisions from both taggers a combination is performed using the calibrated mistag probabilities. The combined tagging decision and calibrated mistag rate are used in the final time-dependent fit, where the calibration parameters are constrained using the combination of their associated statistical and systematic uncertainties. The tagging performances, as well as the effective tagging power, for the three sub-samples and their combination as measured using $B^0_s \to D_s^- \pi^+$ events are reported in table 2.

6.3 Mistag distributions

Because the fit uses the per-candidate mistag prediction, it is necessary to model the distribution of this observable for each event category (SS-only, OS-only, OS-SSK for the signal and each background category). The mistag probability distributions for all $B^0_s$ decay modes, whether signal or background, are obtained using $s$Weighted $B^0_s \to D_s^- \pi^+$ events.
The mistag probability distributions for combinatorial background events are obtained from the upper $B_s^0$ mass sideband in $B_s^0 \to D_s^- \pi^+$ decays. For $B^0$ and $\Lambda_b^0$ backgrounds the mistag distributions are obtained from $s$-Weighted $B^0 \to D^- \pi^+$ events. For the SSK tagger this is justified by the fact that these backgrounds differ by only one spectator quark and should therefore have similar properties with respect to the fragmentation of the $s\bar{s}$ pair. For the OS tagger, the predicted mistag distributions mainly depend on the kinematic properties of the $B$ candidate, which are similar for $B^0$ and $\Lambda_b^0$ backgrounds.

7 Decay-time resolution and acceptance

The decay-time resolution of the detector must be accounted for because of the fast $B_s^0$-$B_s^0$ oscillations. Any mismodelling of the resolution function also potentially biases and affects the precision of the time-dependent $CP$ violation observables. The signal decay-time PDF is convolved with a resolution function that has a different width for each candidate, making use of the per-candidate decay-time uncertainty estimated by the decay-time kinematic fit. This approach requires the per-candidate decay-time uncertainty to be calibrated. The calibration is performed using prompt $D_s^-$ mesons combined with a random track and kinematically weighted to give a sample of “fake $B_s^{0\text{pr}}$” candidates, which have a true lifetime of zero. From the spread of the observed decay times, a scale factor to the estimated decay time resolution is found to be $1.37 \pm 0.10$ [38]. Here the uncertainty is dominated by the systematic uncertainty on the similarity between the kinematically weighted “fake $B_s^{0\text{pr}}$” candidates and the signal. As with the per-candidate mistag, the distribution of per-candidate decay-time uncertainties is modelled for the signal and each type of background. For the signal these distributions are taken from $s$-Weighted data, while for the combinatorial background they are taken from the $B_s^0$ mass sidebands. For other backgrounds, the decay-time error distributions are obtained from simulated events, which are weighted for the data-simulation differences found in $B_s^0 \to D_s^- \pi^+$ signal events.

In the case of background candidates which are either partially reconstructed or in which a particle is misidentified, the decay-time is incorrectly estimated because either the measured mass of the background candidate, the measured momentum, or both, are systematically misreconstructed. For example, in the case of $B_s^0 \to D_s^- \pi^+$ as a background to $B_s^0 \to D_s^0 K^\pm$, the momentum measurement is unbiased, while the reconstructed mass is systematically above the true mass, leading to a systematic increase in the reconstructed decay-time. This effect causes an additional non-Gaussian smearing of the decay-time distribution, which is accounted for in the decay time resolution by nonparametric PDFs obtained from simulated events, referred to as $k$-factor templates.

The decay-time acceptance of $B_s^0 \to D_s^\pm K^\mp$ candidates cannot be floated because its shape is heavily correlated with the $CP$ observables. In particular the upper decay-time acceptance is correlated with $A_f^{\Delta \Gamma}$ and $A_f^{\Delta \Gamma}$. However, in the case of $B_s^0 \to D_s^- \pi^+$, the acceptance can be measured by fixing $\Gamma_s$ and floating the acceptance parameters. The decay-time acceptance in the $B_s^0 \to D_s^\pm K^\mp$ fit is fixed to that found in the $B_s^0 \to D_s^- \pi^+$ data fit, corrected by the acceptance ratio in the two channels in simulated signal events. These simulated events have been weighted in the manner described in section 4. In all cases, the acceptance is described using segments of smooth polynomial functions (“splines”), which
Figure 5. Result of the $s$Fit to the decay-time distribution of $B_s^0 \rightarrow D_s^- \pi^+$ candidates, which is used to measure the decay-time acceptance in $B_s^0 \rightarrow D_s^\mp K^\pm$ decays. The solid curve is measured decay-time acceptance.

can be implemented in an analytic way in the decay-time fit [39]. The spline boundaries ("knots") were chosen in an ad hoc fashion to model reliably the features of the acceptance shape, and placed at 0.5, 1.0, 1.5, 2.0, 3.0, 12.0 ps. Doubling the number of knots results in negligible changes to the nominal fit result. The decay-time fit to the $B_s^0 \rightarrow D_s^- \pi^+$ data is an $s$Fit using the signal PDF from section 1.1, with $S_f$, $S_f$, $A^\Delta\Gamma_f$, and $A^\Delta\Gamma_f$ all fixed to zero, and the knot magnitudes and $\Delta m_s$ floating. The measured value of $\Delta m_s = 17.772 \pm 0.022 \text{ps}^{-1}$ (the uncertainty is statistical only) is in excellent agreement with the published LHCb measurement of $\Delta m_s = 17.768 \pm 0.023 \pm 0.006 \text{ps}^{-1}$ [38]. The time fit to the $B_s^0 \rightarrow D_s^- \pi^+$ data together with the measured decay-time acceptance is shown in figure 5.

8 Decay-time fit to $B_s^0 \rightarrow D_s^\mp K^\pm$

As described previously, two decay-time fitters are used: in one all signal and background time distributions are described ($c$Fit), and in a second the background is statistically subtracted using the $s$Plot technique [15] where only the signal time distributions are described ($s$Fit). In both cases an unbinned maximum likelihood fit is performed to the $CP$ observables defined in eq. (1.5), and the signal decay-time PDF is identical in the two fitters. Both the signal and background PDFs are described in the remainder of this section, but it is important to bear in mind that none of the information about the background PDFs or fixed background parameters is relevant for the $s$Fit. When performing the fits to the decay-time distribution, the following parameters are fixed from independent measurements [12, 13, 40]:

$$\Gamma_s = 0.661 \pm 0.007 \text{ps}^{-1}, \quad \Delta\Gamma_s = 0.106 \pm 0.013 \text{ps}^{-1}, \quad \rho(\Gamma_s, \Delta\Gamma_s) = -0.39,$$
$$\Gamma_{A^0_b} = 0.676 \pm 0.006 \text{ps}^{-1}, \quad \Gamma_d = 0.658 \pm 0.003 \text{ps}^{-1}, \quad \Delta m_s = 17.768 \pm 0.024 \text{ps}^{-1}.$$

Here $\rho(\Gamma_s, \Delta\Gamma_s)$ is the correlation between these two measurements, $\Gamma_{A^0_b}$ is the decay-width of the $A^0_b$ baryon, $\Gamma_d$ is the $B^0$ decay width, and $\Delta m_s$ is the $B_s^0$ oscillation frequency.
The signal production asymmetry is fixed to zero because the fast $B_s^0$ oscillations wash out any initial asymmetry and make its effect on the CP observables negligible. The signal detection asymmetry is fixed to (1.0 ± 0.5)%, with the sign convention in which positive detection asymmetries correspond to a higher efficiency to reconstruct positive kaons [41, 42]. The background production and detection asymmetries are floated within constraints of ±1% for $B_0^0$ and $B^0$ decays, and ±3% for $A_0^0$ decays.

The signal and background mistag and decay-time uncertainty distributions, including k-factors, are modelled by kernel templates as described in section 6 and 7. The tagging calibration parameters are constrained to the values obtained from the control channels for all $B_s^0$ decay modes, except for $B^0$ and $A_0^0$ decays where the calibration parameters of the SSK tagger are fixed to $p_0 = 0.5$, $p_1 = 0$. All modes use the same spline-based decay-time acceptance function described in section 7.

The backgrounds from $B_s^0$ decay modes are all flavour-specific, and are modelled by the decay-time PDF used for $B_s^0 \rightarrow D_\mp K^\pm$ decays convolved with the appropriate decay-time resolution and k-factors model for the given background. The backgrounds from $A_0^0$ decay modes are all described by a single exponential convolved with the appropriate decay-time resolution and k-factor models. The $B^0 \rightarrow D^- K^+$ background is flavour specific and is described with the same PDF as $B_s^0 \rightarrow D_\mp K^\pm$, except with $\Delta m_d$ instead of $\Delta m_s$ in the oscillating terms, $\Gamma_d$ instead of $\Gamma_s$ and the appropriate decay-time resolution and k-factor KEYS templates. The $B^0 \rightarrow D^- \pi^+$ background, on the other hand, is not a flavour specific decay, and is itself sensitive to CP violation as discussed in section 1. Its decay-time PDF therefore includes nonzero $S_f$ and $S_{\bar{f}}$ terms which are constrained to their world-average values [12]. The decay-time PDF of the combinatorial background used in the cFit is a double exponential function split by the tagging category of the event, whose parameters are measured using events in the $B_s^0$ mass sidebands.

All decay-time PDFs include the effects of flavour tagging, are convolved with a single Gaussian representing the per-candidate decay-time resolution, and are multiplied by the decay-time acceptance described in section 7. Once the decay-time PDFs are constructed, the sFit proceeds by fitting the signal PDF to the sWeighted $B_s^0 \rightarrow D_\mp K^\pm$ candidates. The cFit, on the other hand, performs a six-dimensional fit to the decay time, decay-time error, predicted mistag, and the three variables used in the multivariate fit. The $B_s^0$ mass range is restricted to $m(B_s^0) \in [5320, 5420]$ MeV/c$^2$, and the yields of the different signal and background components are fixed to those found in this fit range in the multivariate fit. The decay-time range of the fit is $\tau(B_s^0) \in [0.4, 15.0]$ ps in both cases.

The results of the cFit and sFit for the CP violating observables are given in table 3, and their correlations in table 4. The fits to the decay-time distribution are shown in figure 6 together with the folded asymmetry plots for $D_s^+ K^-$ and $D_s^- K^+$ final states. The folded asymmetry plots show the difference in the rates of $B_s^0$ and $B_s^0$ tagged $D_s^+ K^-$ and $D_s^- K^+$ candidates, plotted in slices of $2\pi/\Delta m_s$, where the sWeights obtained with the multivariate fit have been used to subtract background events. The plotted asymmetry function is drawn using the sFit central values of the CP observables, and is normalised using the expected dilution due to mistag and time resolution.
Table 3. Fitted values of the CP observables to the $B^0_s \to D_s^\mp K^\pm$ time distribution for (left) $sFit$ and (right) $cFit$, where the first uncertainty is statistical, the second is systematic. All parameters other than the CP observables are constrained in the fit.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>$sFit$ fitted value</th>
<th>$cFit$ fitted value</th>
</tr>
</thead>
<tbody>
<tr>
<td>$C_f$</td>
<td>$0.52 \pm 0.25 \pm 0.04$</td>
<td>$0.53 \pm 0.25 \pm 0.04$</td>
</tr>
<tr>
<td>$A^{\Delta \Gamma}_f$</td>
<td>$0.29 \pm 0.42 \pm 0.17$</td>
<td>$0.37 \pm 0.42 \pm 0.20$</td>
</tr>
<tr>
<td>$A^{\Delta \Gamma}_{\bar{f}}$</td>
<td>$0.14 \pm 0.41 \pm 0.18$</td>
<td>$0.20 \pm 0.41 \pm 0.20$</td>
</tr>
<tr>
<td>$S_f$</td>
<td>$-0.90 \pm 0.31 \pm 0.06$</td>
<td>$-1.09 \pm 0.33 \pm 0.08$</td>
</tr>
<tr>
<td>$S_{\bar{f}}$</td>
<td>$-0.36 \pm 0.34 \pm 0.06$</td>
<td>$-0.36 \pm 0.34 \pm 0.08$</td>
</tr>
</tbody>
</table>

Figure 6. Result of the decay-time (top left) $sFit$ and (top right) $cFit$ to the $B^0_s \to D_s^\mp K^\pm$ candidates; the $cFit$ plot groups $B^0_s \to D_s^- \pi^+$ and $B^0_s \to D_s^+ \rho^-$, and also groups $B^0 \to D^- K^+$, $B^0 \to D^\mp \pi^+$, $A^0 \to \omega^- K^+$, $\bar{A}^0 \to \omega^- K^+$, $A^0 \to \rho^- p$, $A^0 \to \bar{D}_s^- p$, and $B^0 \to D_s^- K^+$ together for the sake of clarity. The folded asymmetry plots for (bottom left) $D_s^\mp K^\pm$, and (bottom right) $D_s^\pm K$ are also shown.
Table 4. Statistical correlation matrix of the $B^0_s \to D^{\mp}_s K^\pm$ (top) sFit and (bottom) cFit CP parameters. Other fit parameters have negligible correlations with the CP parameters and are omitted for brevity.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>$C_f$</th>
<th>$A_f^{\Delta \Gamma}$</th>
<th>$A_f^{\Delta \Gamma}$</th>
<th>$S_f$</th>
<th>$S_f$</th>
</tr>
</thead>
<tbody>
<tr>
<td>sFit</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$C_f$</td>
<td>1.000</td>
<td>0.071</td>
<td>0.097</td>
<td>0.117</td>
<td>−0.042</td>
</tr>
<tr>
<td>$A_f^{\Delta \Gamma}$</td>
<td>1.000</td>
<td>0.500</td>
<td>−0.044</td>
<td>−0.003</td>
<td></td>
</tr>
<tr>
<td>$A_f^{\Delta \Gamma}$</td>
<td>1.000</td>
<td>−0.013</td>
<td>−0.005</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S_f$</td>
<td>1.000</td>
<td>0.007</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S_{\bar{f}}$</td>
<td>1.000</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

<table>
<thead>
<tr>
<th>Parameter</th>
<th>$C_f$</th>
<th>$A_f^{\Delta \Gamma}$</th>
<th>$A_f^{\Delta \Gamma}$</th>
<th>$S_f$</th>
<th>$S_f$</th>
</tr>
</thead>
<tbody>
<tr>
<td>cFit</td>
<td></td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$C_f$</td>
<td>1.000</td>
<td>0.084</td>
<td>0.103</td>
<td>−0.008</td>
<td>−0.045</td>
</tr>
<tr>
<td>$A_f^{\Delta \Gamma}$</td>
<td>1.000</td>
<td>0.544</td>
<td>−0.117</td>
<td>−0.022</td>
<td></td>
</tr>
<tr>
<td>$A_f^{\Delta \Gamma}$</td>
<td>1.000</td>
<td>−0.067</td>
<td>−0.032</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S_f$</td>
<td>1.000</td>
<td>0.002</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$S_{\bar{f}}$</td>
<td>1.000</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

9 Systematic uncertainties

Systematic uncertainties arise from the fixed parameters $\Delta m_s$, $\Gamma_s$, and $\Delta \Gamma_s$, and from the limited knowledge of the decay time resolution and acceptance. These uncertainties are estimated using large sets of simulated pseudoexperiments, in which the relevant parameters are varied. The pseudoexperiments are generated with the average of the cFit and sFit central values reported in section 8. They are subsequently processed by the full data fitting procedure: first the multivariate fit to obtain the sWeights, and then the decay time fits. The fitted values of the observables are compared between the nominal fit, where all fixed parameters are kept at their nominal values, and the systematic fit, where each parameter is varied according to its systematic uncertainty. A distribution is formed by normalising the resulting differences to the uncertainties measured in the nominal fit, and the mean and width of this distribution are added in quadrature and conservatively assigned as the systematic uncertainty. The systematic uncertainty on the acceptance is strongly anti-correlated with that due to the fixed value of $\Gamma_s$. This is because the acceptance parameters are determined from the fit to $B^0_s \to D^-_s \pi^+$ data, where $\Gamma_s$ determines the expected exponential slope, so that the acceptance parameterises any difference between the observed and the expected slope. The systematic pseudoexperiments are also used to compute the systematic covariance matrix due to each source of uncertainty.

The total systematic covariance matrix is obtained by adding the individual covariance matrices. The resulting systematic uncertainties are shown in tables 5 and 6 relative to the corresponding statistical uncertainties. The contributions from $\Gamma_s$ and $\Delta \Gamma_s$ are listed independently for comparison to convey a feeling for their relative importance. For this comparison, $\Gamma_s$ and $\Delta \Gamma_s$ are treated as uncorrelated systematic effects. When computing the total, however, the correlations between these two, as well as between them and the
Table 5. Systematic errors, relative to the statistical error, for (top) sFit and (bottom) cFit. The daggered contributions ($\Gamma_s$, $\Delta\Gamma_s$) are given separately for comparison (see text) with the other uncertainties and are not added in quadrature to produce the total.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>$C_f$</th>
<th>$A_f^{\Delta\Gamma}$</th>
<th>$A_f^{\Delta\Gamma}$</th>
<th>$S_f$</th>
<th>$S_f$</th>
</tr>
</thead>
<tbody>
<tr>
<td>sFit $\Delta m_s$</td>
<td>0.062</td>
<td>0.013</td>
<td>0.013</td>
<td>0.104</td>
<td>0.100</td>
</tr>
<tr>
<td>scale factor</td>
<td>0.104</td>
<td>0.004</td>
<td>0.004</td>
<td>0.092</td>
<td>0.096</td>
</tr>
<tr>
<td>$\Delta\Gamma_s^{\dagger}$</td>
<td>0.007</td>
<td>0.261</td>
<td>0.286</td>
<td>0.007</td>
<td>0.007</td>
</tr>
<tr>
<td>$\Gamma_s^{\dagger}$</td>
<td>0.043</td>
<td>0.384</td>
<td>0.385</td>
<td>0.039</td>
<td>0.038</td>
</tr>
<tr>
<td>acceptance, $\Gamma_s$, $\Delta\Gamma_s$</td>
<td>0.043</td>
<td>0.427</td>
<td>0.437</td>
<td>0.039</td>
<td>0.038</td>
</tr>
<tr>
<td>sample splits</td>
<td>0.124</td>
<td>0.000</td>
<td>0.000</td>
<td>0.072</td>
<td>0.071</td>
</tr>
<tr>
<td>total</td>
<td>0.179</td>
<td>0.427</td>
<td>0.437</td>
<td>0.161</td>
<td>0.160</td>
</tr>
</tbody>
</table>

<table>
<thead>
<tr>
<th>Parameter</th>
<th>$C_f$</th>
<th>$A_f^{\Delta\Gamma}$</th>
<th>$A_f^{\Delta\Gamma}$</th>
<th>$S_f$</th>
<th>$S_f$</th>
</tr>
</thead>
<tbody>
<tr>
<td>cFit $\Delta m_s$</td>
<td>0.068</td>
<td>0.014</td>
<td>0.011</td>
<td>0.131</td>
<td>0.126</td>
</tr>
<tr>
<td>scale factor</td>
<td>0.131</td>
<td>0.004</td>
<td>0.004</td>
<td>0.101</td>
<td>0.103</td>
</tr>
<tr>
<td>$\Delta\Gamma_s^{\dagger}$</td>
<td>0.008</td>
<td>0.265</td>
<td>0.274</td>
<td>0.009</td>
<td>0.008</td>
</tr>
<tr>
<td>$\Gamma_s^{\dagger}$</td>
<td>0.049</td>
<td>0.395</td>
<td>0.394</td>
<td>0.048</td>
<td>0.042</td>
</tr>
<tr>
<td>acceptance, $\Gamma_s$, $\Delta\Gamma_s$</td>
<td>0.050</td>
<td>0.461</td>
<td>0.464</td>
<td>0.050</td>
<td>0.043</td>
</tr>
<tr>
<td>comb. bkg. lifetime</td>
<td>0.016</td>
<td>0.069</td>
<td>0.072</td>
<td>0.015</td>
<td>0.005</td>
</tr>
<tr>
<td>sample splits</td>
<td>0.102</td>
<td>0.000</td>
<td>0.000</td>
<td>0.156</td>
<td>0.151</td>
</tr>
<tr>
<td>total</td>
<td>0.187</td>
<td>0.466</td>
<td>0.470</td>
<td>0.234</td>
<td>0.226</td>
</tr>
</tbody>
</table>

acceptance parameters, are accounted for, and the full systematic uncertainty which enters into the total is listed as “acceptance, $\Gamma_s$, $\Delta\Gamma_s$”. The cFit contains fixed parameters describing the decay time of the combinatorial background. These parameters are found to be correlated to the $C_P$ parameters, and a systematic uncertainty is assigned.

The result is cross-checked by splitting the sample into two subsets according to the two magnet polarities, the hardware trigger decision, and the BDTG response. There is good agreement between the cFit and the sFit in each subsample. However, when the sample is split by BDTG response, the weighted averages of the subsamples show a small discrepancy with the nominal fit for $C_f$, $S_f$, and $S_f$, and a corresponding systematic uncertainty is assigned. In addition, fully simulated signal and background events are fitted in order to check for systematic effects due to neglecting correlations between the different variables in the signal and background PDFs. No bias is found.

A potential source of systematic uncertainty is the imperfect knowledge on the tagging parameters $p_0$ and $p_1$. Their uncertainties are propagated into the nominal fits by means of Gaussian constraints, and are therefore included in the statistical error. A number of other possible systematic effects were studied, but found to be negligible. These include possible production and detection asymmetries, and missing or imperfectly modelled backgrounds. Potential systematic effects due to fixed background yields are evaluated by generating pseudoexperiments with the nominal value for these yields, and fitting back with the yields.
Table 6. Systematic uncertainty correlations for (top) sFit and (bottom) cFit.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>$C_f$</th>
<th>$A_{f}^{\Delta \Gamma}$</th>
<th>$A_{\bar{f}}^{\Delta \Gamma}$</th>
<th>$S_f$</th>
<th>$S_{\bar{f}}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>sFit</td>
<td></td>
<td>0.18</td>
<td>0.18</td>
<td>-0.04</td>
<td>-0.04</td>
</tr>
<tr>
<td></td>
<td></td>
<td>1.00</td>
<td>0.95</td>
<td>-0.17</td>
<td>-0.16</td>
</tr>
<tr>
<td></td>
<td></td>
<td>1.00</td>
<td>-0.17</td>
<td>-0.16</td>
<td></td>
</tr>
<tr>
<td></td>
<td></td>
<td>1.00</td>
<td>0.05</td>
<td></td>
<td></td>
</tr>
<tr>
<td></td>
<td></td>
<td>1.00</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>cFit</td>
<td></td>
<td>0.22</td>
<td>0.22</td>
<td>-0.04</td>
<td>-0.03</td>
</tr>
<tr>
<td></td>
<td></td>
<td>1.00</td>
<td>0.96</td>
<td>-0.17</td>
<td>-0.14</td>
</tr>
<tr>
<td></td>
<td></td>
<td>1.00</td>
<td>-0.17</td>
<td>-0.14</td>
<td></td>
</tr>
<tr>
<td></td>
<td></td>
<td>1.00</td>
<td>0.09</td>
<td></td>
<td></td>
</tr>
<tr>
<td></td>
<td></td>
<td>1.00</td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

fixed to twice their nominal value. No significant bias is observed and no systematic uncertainty assigned. No systematic uncertainty is attributed to the imperfect knowledge of the momentum and longitudinal scale of the detector since both effects are taken into account by the systematic uncertainty in $\Delta m_s$.

Both the cFit and sFit are found to be unbiased through studies of large ensembles of pseudoexperiments generated at the best-fit point in data. In addition, differences between the cFit and sFit are evaluated from the distributions of the per-pseudoexperiment differences of the fitted values. Both fitters return compatible results. Indeed, an important result of this analysis is that the sFit technique has been successfully used in an environment with such a large number of variables, parameters and categories. The sFit technique was able to perform an accurate subtraction of a variety of time-dependent backgrounds in a multidimensional fit, including different oscillation frequencies, different tagging behaviours, and backgrounds with modified decay-time distributions due to misreconstructed particles.

10 Interpretation

The measurement of the CP-sensitive parameters is interpreted in terms of $\gamma - 2\beta_s$ and subsequently $\gamma$. For this purpose we have arbitrarily chosen the cFit as the nominal fit result. The strategy is to maximise the following likelihood

$$L(\vec{\alpha}) = \exp\left(-\frac{1}{2} \left(\vec{A}(\vec{\alpha}) - \vec{A}_{\text{obs}}\right)^T V^{-1} \left(\vec{A}(\vec{\alpha}) - \vec{A}_{\text{obs}}\right)\right),$$

(10.1)

where $\vec{\alpha} = (\gamma, \phi_s, r_{D_sK}, \delta)$ is the vector of the physics parameters, $\vec{A}$ is the vector of observables expressed through eqs. (1.6), $\vec{A}_{\text{obs}}$ is the vector of the measured CP violating observables and $V$ is the experimental (statistical and systematic) covariance matrix. Confidence intervals are computed by evaluating the test statistic $\Delta \chi^2 \equiv \chi^2(\vec{\alpha}_{\text{min}}') - \chi^2(\vec{\alpha}_{\text{min}})$,
Figure 7. Graph showing $1 - \text{CL}$ for $\gamma$, together with the central value and the 68.3% CL interval as obtained from the frequentist method described in the text (top). Profile likelihood contours of $r_{D,K}$ vs. $\gamma$ (bottom left), and $\delta$ vs. $\gamma$ (bottom right). The contours are the $1\sigma$ ($2\sigma$) profile likelihood contours, where $\Delta \chi^2 = 1$ ($\Delta \chi^2 = 4$), corresponding to 39% CL (86% CL) in Gaussian approximation. The markers denote the best-fit values.

where $\chi^2(\vec{\alpha}) = -2 \ln \mathcal{L}(\vec{\alpha})$, in a frequentist way following ref. [43]. Here, $\vec{\alpha}_{\text{min}}$ denotes the global maximum of eq. (10.1), and $\vec{\alpha}'_{\text{min}}$ is the conditional maximum when the parameter of interest is fixed to the tested value. The value of $\beta_s$ is constrained to the LHCb measurement from $B^0_s \to J/\psi K^+ K^-$ and $B^0_s \to J/\psi \pi^+ \pi^-$ decays, $\phi_s = 0.01 \pm 0.07$ (stat) $\pm 0.01$ (syst) rad [13]. Neglecting penguin pollution and assuming no BSM contribution in these decays, $\phi_s = -2\beta_s$. The resulting confidence intervals are, at 68% CL,

$$
\gamma = (115^{+28}_{-43})^\circ, \\
\delta = (3^{+19}_{-20})^\circ, \\
r_{D,K} = 0.53^{+0.17}_{-0.16},
$$

where the intervals for the angles are expressed modulo 180°. Figure 7 shows the $1-\text{CL}$ curve for $\gamma$, and the two-dimensional contours of the profile likelihood $\mathcal{L}(\vec{\alpha}'_{\text{min}})$. The systematic contributions to the uncertainty are quoted separately as $\gamma =$
\((115^{+26}_{-35}\,\text{stat})^{+8}_{-25}\,\text{syst}) \pm 4\,(\phi_s)\)°, assuming the central value to be independent from systematic uncertainties and taking the difference in squares of the total and statistical uncertainties.

11 Conclusion

The \(CP\) violation sensitive parameters which describe the \(B^0_s \rightarrow D^{\mp} K^\pm\) decay rates have been measured using a dataset of 1.0 fb\(^{-1}\) of \(pp\) collision data. Their values are found to be

\[
\begin{align*}
C_f &= 0.53 \pm 0.25 \pm 0.04, \\
A_f^{\Delta\Gamma} &= 0.37 \pm 0.42 \pm 0.20, \\
A_f^{\Delta\bar{\Gamma}} &= 0.20 \pm 0.41 \pm 0.20, \\
S_f &= -1.09 \pm 0.33 \pm 0.08, \\
S_{\bar{f}} &= -0.36 \pm 0.34 \pm 0.08,
\end{align*}
\]

where the first uncertainties are statistical and the second are systematic. The results are interpreted in terms of the CKM angle \(\gamma\), which yields \(\gamma = (115^{+28}_{-43})°\), \(\delta = (3^{+19}_{-20})°\) and \(r_{D_s K} = 0.53^{+0.17}_{-0.16}\) (all angles are given modulo 180°) at the 68% confidence level. This is the first measurement of \(\gamma\) performed in this channel.

Acknowledgments

We express our gratitude to our colleagues in the CERN accelerator departments for the excellent performance of the LHC. We thank the technical and administrative staff at the LHCb institutes. We acknowledge support from CERN and from the national agencies: CAPES, CNPq, FAPERJ and FINEP (Brazil); NSFC (China); CNRS/IN2P3 (France); BMBF, DFG, HGF and MPG (Germany); SFI (Ireland); INFN (Italy); FOM and NWO (The Netherlands); MNiSW and NCN (Poland); MEN/IFA (Romania); MinES and FANO (Russia); MinECo (Spain); SNSF and SER (Switzerland); NASU (Ukraine); STFC (United Kingdom); NSF (U.S.A.). The Tier1 computing centres are supported by IN2P3 (France), KIT and BMBF (Germany), INFN (Italy), NWO and SURF (The Netherlands), PIC (Spain), GridPP (United Kingdom). We are indebted to the communities behind the multiple open source software packages on which we depend. We also are thankful for the computing resources and the access to software R&D tools provided by Yandex LLC (Russia). Individual groups or members have received support from EPLANET, Marie Skłodowska-Curie Actions and ERC (European Union), Conseil général de Haute-Savoie, Labex ENIGMASS and OCEVU, Région Auvergne (France), RFBR (Russia), XuntaGal and GENCAT (Spain), Royal Society and Royal Commission for the Exhibition of 1851 (United Kingdom).

Open Access. This article is distributed under the terms of the Creative Commons Attribution License (CC-BY 4.0), which permits any use, distribution and reproduction in any medium, provided the original author(s) and source are credited.
References


[5] R. Fleischer, New strategies to obtain insights into CP violation through $B(s) \to D^{\pm}\pi^\mp, D^*(s)\pi^\mp$ and $B(d) \to D^{\pm}\pi^\mp, D^*(s)\pi^\mp$ decays, Nucl. Phys. B 671 (2003) 459 [hep-ph/0304027] [inSPIRE].

[6] BABAR collaboration, B. Aubert et al., Measurement of time-dependent CP-violating asymmetries and constraints on $\sin(2\beta + \gamma)$ with partial reconstruction of $B \to D^{+\mp}\pi^\pm$ decays, Phys. Rev. D 71 (2005) 112003 [hep-ex/0504035] [inSPIRE].

[7] BABAR collaboration, B. Aubert et al., Measurement of time-dependent CP asymmetries in $B^0 \to D^{(*)+}\pi^-$ and $B^0 \to D^-\rho^+$ decays, Phys. Rev. D 73 (2006) 111101 [hep-ex/0602049] [inSPIRE].


[13] LHCb collaboration, Measurement of CP-violation and the $B_s^0$ meson decay width difference with $B_s^0 \to J/\psi K^+K^-$ and $B_s^0 \to J/\psi\pi^+\pi^-$ decays, Phys. Rev. D 87 (2013) 112010 [arXiv:1304.2600] [inSPIRE].


[38] LHCb collaboration, *Precision measurement of the $B_s^0 - \overline{B}_s^0$ oscillation frequency with the decay $B_s^0 \rightarrow D_s^- \pi^+$*, *New J. Phys.* **15** (2013) 053021 [arXiv:1304.4741] [INSPIRE].


[40] LHCb collaboration, *Precision measurement of the ratio of the $\Lambda_b^0$ to $\overline{B}^0$ lifetimes*, *Phys. Lett. B* **734** (2014) 122 [arXiv:1402.6242] [INSPIRE].


The LHCb collaboration

G. Simi\textsuperscript{22}, M. Sirendi\textsuperscript{47}, N. Skidmore\textsuperscript{46}, T. Skwarnicki\textsuperscript{59}, N.A. Smith\textsuperscript{52}, E. Smith\textsuperscript{55,49}, E. Smith\textsuperscript{53}, J. Smith\textsuperscript{47}, M. Smith\textsuperscript{54}, H. Snoek\textsuperscript{41}, M.D. Sokoloff\textsuperscript{57}, F.J.P. Soles\textsuperscript{51}, F. Soomro\textsuperscript{39}, D. Souza\textsuperscript{46}, B. Souza De Paula\textsuperscript{2}, B. Spaan\textsuperscript{2}, A. Sparkes\textsuperscript{50}, P. Spradlin\textsuperscript{51}, S. Sridharan\textsuperscript{38}, F. Stagni\textsuperscript{38}, M. Stahl\textsuperscript{11}, S. Stahl\textsuperscript{11}, O. Steinkamp\textsuperscript{40}, O. Stenyakin\textsuperscript{35}, S. Stevenson\textsuperscript{55}, S. Stoica\textsuperscript{29}, S. Stone\textsuperscript{59}, B. Storaci\textsuperscript{40}, S. Stracke\textsuperscript{23,38}, M. Stratil\textsuperscript{29}, U. Straumann\textsuperscript{40}, R. Stroili\textsuperscript{22}, V.K. Subbiah\textsuperscript{38}, L. Sun\textsuperscript{57}, W. Sutcliffe\textsuperscript{7}, S. Swientek\textsuperscript{27}, S. Szczekowski\textsuperscript{28}, P. Szczypa\textsuperscript{39,38}, D. Szilard\textsuperscript{2}, T. Szumlak\textsuperscript{27}, S. T'Jampens\textsuperscript{4}, M. Teklishyn\textsuperscript{7}, G. Tellarini\textsuperscript{16,f}, F. Teubert\textsuperscript{38}, C. Thomas\textsuperscript{55}, E. Thomas\textsuperscript{38}, J. van Tilburg\textsuperscript{41}, V. Tisserand\textsuperscript{14}, M. Tobin\textsuperscript{39}, S. Tolk\textsuperscript{42}, L. Tomassetti\textsuperscript{16,f}, D. Tonelli\textsuperscript{38}, S. Topp-Joergensen\textsuperscript{55}, N. Torr\textsuperscript{55}, E. Tournefier\textsuperscript{4}, S. Tourneur\textsuperscript{39}, M.T. Tran\textsuperscript{39}, M. Tresch\textsuperscript{40}, A. Tsaregorodtsev\textsuperscript{6}, P. Tsopelas\textsuperscript{41}, N. Tuning\textsuperscript{41}, M. Ubeda Garcia\textsuperscript{38}, A. Ukraje\textsuperscript{28}, A. Ustyuzhanin\textsuperscript{63}, U. Uwer\textsuperscript{11}, V. Vagnoni\textsuperscript{14}, G. Valenti\textsuperscript{14}, A. Vallier\textsuperscript{7}, R. Vazquez Gomez\textsuperscript{18}, P. Vazquez Regueiro\textsuperscript{37}, C. Vázquez Sierra\textsuperscript{47}, S. Vecchi\textsuperscript{16}, J.J. Veltinis\textsuperscript{46}, M. Velti\textsuperscript{17,b}, G. Veneziano\textsuperscript{10}, M. Vesterinen\textsuperscript{11}, B. Viaud\textsuperscript{7}, D. Vieira\textsuperscript{2}, M. Vieites Diaz\textsuperscript{37}, X. Vilasis-Cardona\textsuperscript{36,p}, A. Vollhardt\textsuperscript{40}, D. Volyanskyy\textsuperscript{10}, D. Voong\textsuperscript{46}, A. Vorobyev\textsuperscript{30}, V. Vorobyev\textsuperscript{34}, C. Voß\textsuperscript{92}, H. Voss\textsuperscript{10}, J.A. de Vries\textsuperscript{41}, R. Waldi\textsuperscript{92}, C. Wallace\textsuperscript{48}, R. Wallace\textsuperscript{12}, J. Walsh\textsuperscript{23}, S. Wandernoth\textsuperscript{11}, J. Wang\textsuperscript{59}, D.R. Ward\textsuperscript{47}, N.K. Watson\textsuperscript{45}, D. Websdale\textsuperscript{33}, M. Whitehead\textsuperscript{48}, J. Wicht\textsuperscript{38}, D. Wiedner\textsuperscript{11}, G. Wilkinson\textsuperscript{55}, M.P. Williams\textsuperscript{45}, M. Williams\textsuperscript{56}, F.F. Wilson\textsuperscript{49}, J. Wimerley\textsuperscript{58}, J. Wisniewski\textsuperscript{9}, W. Wislicki\textsuperscript{26}, M. Wormser\textsuperscript{7}, S.A. Wotton\textsuperscript{47}, S. Wright\textsuperscript{47}, S. Wu\textsuperscript{3}, K. Wyllie\textsuperscript{38}, Y. Xie\textsuperscript{61}, Z. Xing\textsuperscript{59}, Z. Xu\textsuperscript{39}, Z. Yang\textsuperscript{3}, X. Yuan\textsuperscript{3}, O. Yushchenko\textsuperscript{35}, M. Zangoli\textsuperscript{14}, M. Zavertyaev\textsuperscript{10,b}, L. Zhang\textsuperscript{59}, W.C. Zhang\textsuperscript{12}, Y. Zhang\textsuperscript{3}, A. Zhelezov\textsuperscript{11}, A. Zhokhov\textsuperscript{11}, L. Zhong\textsuperscript{3}, A. Zvyagin

---

1 Centro Brasileiro de Pesquisas Físicas (CBPF), Rio de Janeiro, Brazil
2 Universidade Federal do Rio de Janeiro (UFRJ), Rio de Janeiro, Brazil
3 Center for High Energy Physics, Tsinghua University, Beijing, China
4 LAPP, Université de Savoie, CNRS/IN2P3, Annecy-Le-Vieux, France
5 Clermont Université, Université Blaise Pascal, CNRS/IN2P3, LPC, Clermont-Ferrand, France
6 CPPM, Aix-Marseille Université, CNRS/IN2P3, Marseille, France
7 LAL, Université Paris-Sud, CNRS/IN2P3, Orsay, France
8 LPNHE, Université Pierre et Marie Curie, Université Paris Diderot, CNRS/IN2P3, Paris, France
9 Fakultät Physik, Technische Universität Dortmund, Dortmund, Germany
10 Max-Planck-Institut für Kernphysik (MPIK), Heidelberg, Germany
11 Physikalisches Institut, Ruprecht-Karls-Universität Heidelberg, Heidelberg, Germany
12 School of Physics, University College Dublin, Dublin, Ireland
13 Sezione INFN di Bari, Bari, Italy
14 Sezione INFN di Bologna, Bologna, Italy
15 Sezione INFN di Cagliari, Cagliari, Italy
16 Sezione INFN di Ferrara, Ferrara, Italy
17 Sezione INFN di Firenze, Firenze, Italy
18 Laboratori Nazionali dell’INFN di Frascati, Frascati, Italy
19 Sezione INFN di Genova, Genova, Italy
20 Sezione INFN di Milano Bicocca, Milano, Italy
21 Sezione INFN di Milano, Milano, Italy
22 Sezione INFN di Padova, Padova, Italy
23 Sezione INFN di Pisa, Pisa, Italy
24 Sezione INFN di Roma Tor Vergata, Roma, Italy
25 Sezione INFN di Roma La Sapienza, Roma, Italy
26 Henryk Niewodniczanski Institute of Nuclear Physics Polish Academy of Sciences, Kraków, Poland
27 AGH - University of Science and Technology, Faculty of Physics and Applied Computer Science, Kraków, Poland
28 National Center for Nuclear Research (NCBJ), Warsaw, Poland
Horia Hulubei National Institute of Physics and Nuclear Engineering, Bucharest-Magurele, Romania

Petersburg Nuclear Physics Institute (PNPI), Gatchina, Russia

Institute of Theoretical and Experimental Physics (ITEP), Moscow, Russia

Institute of Nuclear Physics, Moscow State University (SINP MSU), Moscow, Russia

Institute for Nuclear Research of the Russian Academy of Sciences (INR RAN), Moscow, Russia

Budker Institute of Nuclear Physics (SB RAS) and Novosibirsk State University, Novosibirsk, Russia

Institute for High Energy Physics (IHEP), Protvino, Russia

Universitat de Barcelona, Barcelona, Spain

Universidad de Santiago de Compostela, Santiago de Compostela, Spain

European Organization for Nuclear Research (CERN), Geneva, Switzerland

Ecole Polytechnique Fédérale de Lausanne (EPFL), Lausanne, Switzerland

Physik-Institut, Universität Zürich, Zürich, Switzerland

Nikhef National Institute for Subatomic Physics, Amsterdam, The Netherlands

Nikhef National Institute for Subatomic Physics and VU University Amsterdam, Amsterdam, The Netherlands

NSC Kharkiv Institute of Physics and Technology (NSC KIPT), Kharkiv, Ukraine

Institute for Nuclear Research of the National Academy of Sciences (KINR), Kyiv, Ukraine

University of Birmingham, Birmingham, United Kingdom

H.H. Wills Physics Laboratory, University of Bristol, Bristol, United Kingdom

Cavendish Laboratory, University of Cambridge, Cambridge, United Kingdom

Department of Physics, University of Warwick, Coventry, United Kingdom

STFC Rutherford Appleton Laboratory, Didcot, United Kingdom

School of Physics and Astronomy, University of Edinburgh, Edinburgh, United Kingdom

School of Physics and Astronomy, University of Glasgow, Glasgow, United Kingdom

Oliver Lodge Laboratory, University of Liverpool, Liverpool, United Kingdom

Imperial College London, London, United Kingdom

School of Physics and Astronomy, University of Manchester, Manchester, United Kingdom

Department of Physics, University of Oxford, Oxford, United Kingdom

Massachusetts Institute of Technology, Cambridge, MA, United States

University of Cincinnati, Cincinnati, OH, United States

University of Maryland, College Park, MD, United States

Syracuse University, Syracuse, NY, United States

Pontifícia Universidade Católica do Rio de Janeiro (PUC-Rio), Rio de Janeiro, Brazil, associated to

Institute of Particle Physics, Central China Normal University, Wuhan, Hubei, China, associated to

Institut für Physik, Universität Rostock, Rostock, Germany, associated to

National Research Centre Kurchatov Institute, Moscow, Russia, associated to

Instituto de Fisica Corpuscular (IFIC), Universitat de Valencia-CSIC, Valencia, Spain, associated to

KVI - University of Groningen, Groningen, The Netherlands, associated to

Celal Bayar University, Manisa, Turkey, associated to

Universidade Federal do Triângulo Mineiro (UFTM), Uberaba-MG, Brazil

P.N. Lebedev Physical Institute, Russian Academy of Science (LPI RAS), Moscow, Russia

Università di Bari, Bari, Italy

Università di Bologna, Bologna, Italy

Università di Cagliari, Cagliari, Italy

Università di Ferrara, Ferrara, Italy

Università di Firenze, Firenze, Italy

Università di Urbino, Urbino, Italy
Università di Modena e Reggio Emilia, Modena, Italy

Università di Genova, Genova, Italy

Università di Milano Bicocca, Milano, Italy

Università di Roma Tor Vergata, Roma, Italy

Università di Roma La Sapienza, Roma, Italy

Università della Basilicata, Potenza, Italy

AGH - University of Science and Technology, Faculty of Computer Science, Electronics and Telecommunications, Kraków, Poland

LIFAELS, La Salle, Universitat Ramon Llull, Barcelona, Spain

Hanoi University of Science, Hanoi, Viet Nam

Università di Padova, Padova, Italy

Università di Pisa, Pisa, Italy

Scuola Normale Superiore, Pisa, Italy

Università degli Studi di Milano, Milano, Italy