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Extraplanar H\textsc{ii} regions in the edge-on spiral galaxies NGC 3628 and NGC 4522

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ABSTRACT

Context. Gas infall and outflow are critical for determining the star formation rate and chemical evolution of galaxies but direct measurements of gas flows are difficult to make. Young massive stars and H\textsc{ii} regions in the halos of galaxies are potential tracers for accretion and/or outflows of gas.

Aims. Gas phase abundances of three H\textsc{ii} regions in the lower halos of the edge-on galaxies NGC 3628 and NGC 4522 are determined by analyzing optical long-slit spectra. The observed regions have projected distances to the midplane of their host from 1.4 to 3 kpc.

Methods. With the measured flux densities of the optical nebular emission lines, we derived the oxygen abundance 12 + log(O/H) for the three extraplanar H\textsc{ii} regions. The analysis was based on one theoretical and two empirical strong-line calibration methods.

Results. The resulting oxygen abundances of the extraplanar H\textsc{ii} regions are comparable to the disk H\textsc{ii} regions in one case and are a little lower in the other case. Since our results depend on the accuracy of the metallicity determinations, we critically discuss the difference of the calibration methods we applied and confirm previously noted offsets. From our measurements, we argue that these three extraplanar H\textsc{ii} regions were formed in the disk or at least from disk material. We discuss the processes that could transport disk material into the lower halo of these systems and conclude that gravitational interaction with a companion galaxy is most likely for NGC 3628 while ram pressure is favored in the case of NGC 4522.

Key words. galaxies: abundances – ISM: abundances – galaxies: individual: NGC 3628 – galaxies: individual: NGC 4522 – H\textsc{ii} regions – galaxy: evolution

1. Introduction

Star formation in spiral galaxies usually takes place in dense regions of the interstellar medium (ISM) close to the midplane of the disk. However, there is evidence for low-level star formation high above the midplane in the Milky Way and in some external galaxies. OB stars exist at large distances from the plane in the Milky Way (e.g., Greenstein & Sargent 1974; Kilkenny et al. 1995) and in the halos of other galaxies (e.g., Comerón et al. 2001). Furthermore, extraplanar H\textsc{ii} regions have been detected in galaxies with a strong indication for an exchange of matter between the disk and the halo; see, for example, NGC 55 (Ferguson et al. 1996), NGC 891, and NGC 4013 (Howk & Savage 1999), as well as the Virgo galaxy NGC 4402 (Cortese et al. 2004). In some objects, such extraplanar H\textsc{ii} regions are found at even larger galactocentric distances and hence are attributed to the intracluster or intergalactic medium of interacting galaxies, where the H\textsc{ii} regions probably formed in tidal debris (Mendes de Oliveira et al. 2004; Ryan-Weber et al. 2004), or by ram pressure stripping, such as in the Virgo cluster (Gerhard et al. 2002; Oosterloo & van Gorkom 2005).

There are different scenarios to explain the presence of young stars and star formation in these unusual environments.

(I) For the young stars, ejection from the disk by either asymmetric supernovae of a close binary system (e.g., Stone 1982) or dynamical three- or four-body encounters in dense stellar systems (e.g., Poveda et al. 1967) could be responsible for their existence far from the galactic midplane. Supersonic space motions of massive stars are possible with this mechanism, as is apparent from bow shocks induced by these stars (e.g., Gvaramadze & Bomans 2008). (II) Star formation in the halo itself would not only explain the presence of young stars in the halos (e.g., Keenan et al. 1986) but would also explain the existence of H\textsc{ii} regions in the halo, formed with halo material. (III) Tidal forces could induce clumps in material drawn out of galactic disks or in gas falling back onto a disk (e.g., Barnes & Hernquist 1996), which then would form stars and thus explain extraplanar H\textsc{ii} regions.

One way to constrain the origin of the above-mentioned star formation in galactic halos is the determination of metallicities of the extraplanar and disk H\textsc{ii} regions. Low metallicities of extraplanar regions in comparison to high metallicity disk regions would indicate that the region is formed by low metallicity halo gas (second scenario). Comparable metallicities between disk and extraplanar regions would lead to the conclusion that the extraplanar H\textsc{ii} region is formed by disk material (first and third scenario).

Tüllmann et al. (2003) studied extraplanar H\textsc{ii} regions in the galaxy NGC 55. They found a lower metallicity of the extraplanar regions in comparison to H\textsc{ii} regions in the midplane...
and concluded that the second scenario is most likely responsible and that the HII regions formed from halo material. This result was verified by Kudritzki et al. (2016). The metallicities of the extragalactic HII regions in Stephan’s Quintet (Mendes de Oliveira et al. 2004), in the interacting galaxies studied by Ryan-Weber et al. (2004), and in the extraplanar HII regions in the Virgo galaxy NGC 4402 are all enhanced in metallicity. In the case of the Virgo galaxy NGC 4388, subsolar metallicities are observed. This suggests there may be no single mechanism for producing extraplanar HII regions. However, there have only been a limited number of studies so far and the nature and origin of such objects requires further study.

NGC 3628 in the Leo Triplet and NGC 4522 in the Virgo Cluster are two edge-on galaxies. These galaxies are known to have extended diffuse Hα structure, which we refer to as extraplanar diffuse ionized gas (eDIG; Rossa & Dettmar 2000). Three extraplanar HII regions are known in these systems. We here present optical long-slit spectroscopy of these regions, which we analyze to determine metallicities and thus place further constraints on the origins of the young stars and gas at large scale heights.

The extraplanar HII region (1) of NGC 4522 is offset from the main star-forming region defined by Hα and lies at a projected distance of 1.4 kpc. This is also visible in the HI map of NGC 4522 (Kenney et al. 2004). According to Kenney et al. (2004, Fig. 4) region (1) resides in extraplanar gas that was stripped out and removed from the disk via ram pressure. The extraplanar HII regions (2) and (3) lie in NGC 3628 and have projected distances from the midplane of their host galaxy of 2.8 kpc and 3.0 kpc, respectively. We use the oxygen abundance 12 + log(O/H) to derive metallicities. Oxygen is widely used as a reference element because it is relatively abundant, emits strong lines in the optical (as an important cooling element), and is observed in many ionization states (Kewley & Dopita 2002). With the aid of excitation models, oxygen abundances can also be used to derive electron temperatures. There are different ways to derive oxygen abundances and temperatures. Direct methods use ratios of auroral and nebular lines as they are temperature sensitive; (e.g., the auroral line [OII] λ3726 in combination with [OII] λ4959, 5007). The line intensities of the auroral lines in spectra are weak, especially at solar metallicities, or higher. We did not detect any of these weak lines. Our spectra contain the strong emission lines such as the Balmer lines, [OII] λ3726, 3729, [OIII] λ4959, 5007, and [NII] λλ6548, 6583. Owing to the spectral resolution, the lines of [OII] λ3726.0 and [O II] λ3728.8 were blended in the spectra and hence we use the sum of the doublet lines. To derive oxygen abundances, we use strong-line calibrations. There are different strong-line calibration types, i.e., empirical (e.g., Pilyugin 2001; Pilyugin & Thuan 2005; Pettini & Pagel 2004; Pilyugin et al. 2014) and theoretical (e.g., Kewley & Dopita 2002; Kobulnicky & Kewley 2004), which show well-known offsets between 0.1 and 0.7 dex (Stasińska 2002; Modjaz et al. 2008; Moustakas et al. 2010). The empirical methods use a fitted relationship between direct metallicities and strong-line ratios in the optical of a sample of HII regions to derive metallicity relations while the theoretical methods use photoionization models such as Cloudy (Ferland et al. 1998) to predict how theoretical emission line ratios depend on the input metallicity. We use three different calibrations, investigate their offsets, and discuss possible reasons. Using our measurements of the oxygen abundances, we explore the origin of the extraplanar HII regions.

The data reduction and observational strategy is given in Sect. 2. We then describe the measurement of the emission lines and the different calibration types for the analysis in Sect. 3. Section 4 contains the results. In Sect. 5 a short summary and the conclusions are given.

2. Data

2.1. Spectroscopic data

The spectroscopic data were obtained at the European Southern Observatory (ESO) La Silla with ESO Faint Object Spectrograph and Camera (EFOSC2 v.2) attached to the 3.6 m telescope during the nights of March 7–9, 2000. A 2048 × 2048 CCD detector with a pixel size of 15 μm was used. The field of view is 5.4′ × 5.4′ with a pixel scale of 0.316″/pixel. With the two different grisms 7 (3270–5240 Å) and 9 (4700–6700 Å), we were able to detect the important lines in the optical. The measured spectral resolution is 6.7 Å in both grisms. The slit had a length of 5′ and a width of 1″. The seeing was around 1.1″ (see Table 1). The long slit was oriented differently in the galaxies. The slit positions for the galaxies NGC 3628 and NGC 4522 are shown in Figs. 1 and 2 and we report the position angles in Table 1. The slit position was chosen to go through both extraplanar regions and the disk. The total integration time for NGC 3628 was one hour in each grism. For NGC 4522, two slit positions were observed with the first covering the extraplanar region (s1) and the second the galaxy (s2). The total integration time was 20 min in.

Table 1. Galaxy parameter and observation log.

<table>
<thead>
<tr>
<th>Galaxy</th>
<th>RA 2000</th>
<th>Dec 2000</th>
<th>Hubble type</th>
<th>v [km s⁻¹]</th>
<th>D [Mpc]</th>
<th>i [°]</th>
<th>Grism#</th>
<th>t_int [s]</th>
<th>PA</th>
<th>Seeing</th>
<th>AM</th>
</tr>
</thead>
<tbody>
<tr>
<td>NGC 3628</td>
<td>13°20'17.0''</td>
<td>13°35'23''</td>
<td>Sb pec</td>
<td>843 ± 1</td>
<td>12.2 ± 2.4</td>
<td>87°</td>
<td>7</td>
<td>2 × 1800</td>
<td>9</td>
<td>2 × 1800</td>
<td>1.11</td>
</tr>
<tr>
<td>NGC 4522</td>
<td>12°33'39.7''</td>
<td>09°10'30''</td>
<td>SB(s)cD</td>
<td>2329 ± 1</td>
<td>16.8 ± 1.3</td>
<td>75 ± 5°</td>
<td>7</td>
<td>s1</td>
<td>1 × 1200</td>
<td>7</td>
<td>1 × 1200</td>
</tr>
</tbody>
</table>


slit position s1 with grism 7 and slit position s2 with both grisms. Slit position s1 with grism 9 was observed with two exposures of 20 min each (see Table 1 for details of the observations).

The data reduction was carried out using the image reduction and analysis facility (IRAF, Tody 1986, 1993). This includes the corrections for bias and overscan, flat-field response, and cosmic rays with L.A. Cosmic (van Dokkum 2001). Additionally, the wavelength calibration was achieved with helium-argon spectra and the flux calibration was carried out with standard stars (EG274, LTT2415). In order to obtain the pure emission from the HII regions, night sky background emission above and below these regions was averaged and subtracted with the IRAF task background. There are some residuals at 6300 Å in the spectrum of NGC 4522 grism 9 left, which do not influence the data analysis. The galaxy continuum was subtracted by the IRAF task continuum, which was used with a linear spline 12th order and finally the spectra were coadded. Figures 3, 4 show the resulting spectra. To correct for stellar absorption, the non-continuum-subtracted spectra were coadded. All listed flux spectral densities are taken from these non-continuum-subtracted spectra.

2.2. Narrowband imaging

The Hα imaging for the galaxies discussed in this paper was generated from two different sources: NGC 3628 was observed on May 8, 1991 using the ESO New Technology Telescope (NTT, red arm of the ESO Multi-Mode Instrument (EMMI)) for 900 s in Hα and 300 s in R as part of ESO program 047.01-003. The detector was a Ford 20482 single CCD (ESO CCD #24). Both images were taken with subarcsecond seeing. The data and all relevant calibration files taken during the run and a few days before and after the observing run were retrieved from the ESO archive and re-reduced by us using IRAF in the usual manner. L.A. Cosmic (van Dokkum 2001) was used to clear the image...
of cosmic rays to produce the final Hα and final R image. The continuum subtracted image was then produced by aligning Hα and R band, homogenizing the PSF, and subtracting the appropriately scaled R image from the Hα image \((\text{e.g., Skillman et al. } 1997)\). To determine the scaling factor, we measured the apparent fluxes for several stars in the field on both the Hα and R images. Scaling, subtraction, and correction to reach an optimal correction of the continuum light of the galaxy was performed with our own IRAF scripts. The NGC 4522 Hα data were taken directly from the work of Koopmann et al. (2001) and the Galaxy On Line Database Milano Network (GOLD-Mine, Gavazzi et al. 2003). The images were astrometrically calibrated to the Digitized Sky Survey (DSS) system. Flux calibration of the continuum subtracted Hα images was performed by transferring the total Sloan Digital Sky Survey (SDSS) r′ band flux of each galaxy to its continuum image. With the derived scaling factor \(1\) (see above), the flux scale can then be directly transferred to the continuum corrected Hα image. A conversion from magnitudes to flux was performed using the fact that SDSS is calibrated in AB magnitudes, so that the zero-point flux density of each filter is 3631 Jy \((\text{with } 1 \text{ Jy} = 1 \text{ Jansky} = 10^{-26} \text{ W Hz}^{-1} \text{ m}^{-2} = 10^{-23} \text{ erg s}^{-1} \text{ Hz}^{-1} \text{ cm}^{-2})\)\].

### 3. Analysis

#### 3.1. Emission line measurements

The spectra were analyzed using the splot task in IRAF. We chose the size of the extraction region of the extraplanar and disk HII regions depending on the extent of their Hα emission along the spatial axis in the spectra. The emission line flux densities were measured with the deblending tool of IRAF/splot, with which we fit Gaussians to obtain fluxes and central peak positions. We used the overlapping wavelength range of the two grisms for putting the spectra on a common flux scale. Unfortunately, the straight calibration showed different absolute fluxes of the emission lines due to differences in sky transparency conditions. Thus, it was necessary to scale the fluxes of the two grisms

\[1\] http://classic.sdss.org/dr7/algorithms/fluxcal.html
3.2. Stellar continuum and reddening correction

Since absorption lines in the underlying stellar continuum can affect the measurement of the emission lines (e.g., the Balmer lines), we need to investigate this effect in our spectra. We identified a weak contribution of underlying stellar continuum in the disk H\textsc{ii} regions, which was even smaller in the extraplanar H\textsc{ii} region (1) of NGC 4522 in our spectra. We corrected the H\textalpha and H\beta lines for absorption with 2 Å in the equivalent width (EW; Moustakas et al. 2010). There were no metal line absorption features visible in the stellar continuum of the disk H\textsc{ii} regions and therefore it was not possible to perform a reliable population synthesis, for example, with STARLIGHT (Cid Fernandes et al. 2004). Nevertheless, we tried pPXF fitting (Cappellari & Emsellem 2004; Cappellari 2017) for one H\textsc{ii} region. This routine does full spectral fitting to extract stellar and gas kinematics and stellar populations information simultaneously (Cappellari 2017). This was performed for the disk H\textsc{ii} region (2) in grism 7 of NGC 4522 to be sure that our applied standard absorption correction is reliable. As Fig. 5 shows, the pPXF-predicted absorption features are very weak. The resulting corrected line ratio of H\alpha/H\beta from pPXF was 0.26 and, after dereddening, 0.29, which is at odds with the theoretical value of 0.46. Clearly, because of the missing information of the metal-absorption features, the pPXF code cannot find a good solution for our data. Therefore, a first order EW correction of the Balmer lines seems to be the best option. There was no continuum visible in the extraplanar H\textsc{ii} regions (2), (3) of NGC 3628 and therefore no EW correction was applied, which is similar to the approach taken for the analysis of faint high redshift galaxies (e.g., Mannucci et al. 2009). To account for extinction, the fluxes were dereddened with the IRAF task redcorr. For this we used the theoretical Balmer decrement of H\alpha/H\beta = 2.87 (Case B, 10 000 K), the Galactic extinction law f(\lambda) from Savage & Mathis (1979) and the extinction parameter c (see Table 2) with the following equation used in IRAF:

$$I_{\text{corr}}(\lambda) = I_{\text{obs}}(\lambda) \times 10^{-f(\lambda)c},$$

where $I_{\text{corr}}(\lambda)$ is the corrected flux and $I_{\text{obs}}(\lambda)$ the observed flux. To judge the effect of the chosen extinction law, we tested our reddening correction using the LMC extinction curve (Howarth 1983). For the H\alpha/H\beta line ratio, the change is 0.3%, whereas the biggest change is 3.5% with the [OII]/H\beta line ratio. We claim the reddening correction is only weakly dependent on the adopted reddening curve. If possible, the lines were measured in grism 9. The H\beta fluxes are listed in Table 2. The dereddened and scaled fluxes relative to H\beta are listed in Table 3.

3.3. Errors

3.3.1. Flux errors

Errors in the emission line fluxes can arise from the calibration and from the measurement of the emission lines with the task splot in IRAF. Statistical errors are calculated within splot. These errors are explained and discussed below.

- The error in the flux calibration was tested by applying the calibration to a standard star and comparing the resultant fluxes with the known fluxes. The error was found to be 12% in grism 7 and 6% in grism 9.

<table>
<thead>
<tr>
<th>Galaxy</th>
<th>Region</th>
<th>H\beta grism 7</th>
<th>H\beta grism 9</th>
<th>c</th>
<th>v_\text{helio} [\pm 23 \text{ km s}^{-1}]</th>
<th>Size H\textsc{ii} region [pc]</th>
</tr>
</thead>
<tbody>
<tr>
<td>NGC 3628</td>
<td>1</td>
<td>(4.63 ± 0.67) × 10^{−16}</td>
<td>(4.31 ± 0.40) × 10^{−16}</td>
<td>0.95</td>
<td>671</td>
<td></td>
</tr>
<tr>
<td></td>
<td>2 (e)</td>
<td>(3.89 ± 0.56) × 10^{−16}</td>
<td>(3.89 ± 0.36) × 10^{−16}</td>
<td>0.12(1)</td>
<td>648</td>
<td>150 ± 30</td>
</tr>
<tr>
<td></td>
<td>3 (e)</td>
<td>(2.53 ± 0.36) × 10^{−16}</td>
<td>(2.53 ± 0.23) × 10^{−16}</td>
<td>0.00(2)</td>
<td>650</td>
<td>150 ± 30</td>
</tr>
<tr>
<td>NGC 4522</td>
<td>1 (e)</td>
<td>(6.22 ± 0.90) × 10^{−16}</td>
<td>(6.12 ± 0.56) × 10^{−16}</td>
<td>0.30</td>
<td>2366</td>
<td>325 ± 40</td>
</tr>
<tr>
<td></td>
<td>2</td>
<td>(1.22 ± 0.18) × 10^{−15}</td>
<td>(1.35 ± 0.12) × 10^{−15}</td>
<td>0.33</td>
<td>2320</td>
<td></td>
</tr>
<tr>
<td></td>
<td>3</td>
<td>(5.67 ± 0.82) × 10^{−16}</td>
<td>(6.04 ± 0.56) × 10^{−16}</td>
<td>0.37</td>
<td>2253</td>
<td></td>
</tr>
</tbody>
</table>

Notes. H\beta fluxes in erg cm^{-2} s^{-1}; (e) = extraplanar. (1) No EW correction applied; (2) No EW correction, and no reddening correction applied.

to one another first. The common lines of H\beta and [OIII] \lambda 5007 were used for this in grism 7 and grism 9.
The total error in wavelength is 2.13 Å in grism 7, which mainly originates from the high splot noise ($\sigma$). For each emission line, this error was determined separately. We tested these errors with the [OII] λλ 3726, 3728 and [OIII] λ5007 lines in grism 7 and with Hz and [OIII] λ5007 lines in grism 9. By repeated measurements of the same emission line with variations in applying the deblending tool within the splot task, for example, for variations in the initial estimate of the peak position and continuum, the resulting errors were found to be similar to the splot errors. Therefore the splot errors were used for further error propagation.

The resulting flux error (sum of the squares) for Hβ is 14.4% in grism 7 and 9.2% in grism 9. The errors are propagated throughout the complete abundance analysis (Gaussian error propagation). To calculate the final error in the oxygen abundance determination, we added the errors of the different calibrations to the flux error (sum of squares).

### 3.3.2. Wavelength errors

The errors of the velocity determination result from the wavelength calibration and the measurement of the peak positions of the emission lines, which are given below in more detail.

- The applied wavelength calibration function had a dispersion of 0.5 Å, which was taken as the error of the calibration. This was further checked by measuring the position of three sky lines per grism and comparing these with values in the literature (Slanger et al. 2003). The error was confirmed to be 0.5 Å.
- The error of the peak positions was determined by measuring four times the same emission line with variation in the initial estimate of the peak position using the deblending tool within the splot task. The average of the errors given by splot was taken as the resulting error.

The total error in wavelength is 2.13 Å in grism 7, which mainly originates from the high splot errors due to the low signal-to-noise ratio (S/N) in this grism. The total wavelength error is 0.50 Å in grism 9.

### 3.4. Parameter determination

The projected distances from the extraplanar regions to the host galaxy were determined as the projected height perpendicular to the central line of the disk, which was defined by an isophotal fit to the 3.6 μm image for NGC 4522 and by the central emission in the Hz image for NGC 3628. The heliocentric velocities of the HII regions were determined from the redshifts of the Hz line, which is the brightest line. The heliocentric correction was calculated by the IRAF task rvcorrect. The metallicities were derived using the oxygen abundance $12 + \log(O/H)$ by the following three different methods.

#### 3.4.1. Method 1

Kewley & Dopita (2002; theoretical) use stellar population synthesis and photoionization models to produce a set of ionization parameter and abundance diagnostics based on the use of strong optical emission lines. These authors claim the ratio of [NII] λ6583 to [OIII] λλ 3726, 3729 is a reliable diagnostic because it is independent of the ionization parameter and has no local maximum like the $R_3$ parameter. Additionally, the [NII], [OII] lines are relatively strong, even in low S/N spectra. With $R = I([NII]λ6583)/I([OIII]λλ 3726, 3729)$ a first estimate of the oxygen abundance is calculated

$$12 + \log(O/H) = \log(1.54020 + 1.26602R + 0.167977R^2) + 8.93.$$  

Following that paper, the oxygen abundance is taken as final if the result is $12 + \log(O/H) \geq 8.6$, which was the case in all regions discussed in this paper.

#### Limitsations: Kewley & Dopita (2002) mention that the observed [NII]/[OII] abundance ratio shows large scatter from object to object due to the varying age distribution of the stellar population for a sample of HII regions. Furthermore, in this method, the oxygen abundance depends only on the [NII]/[OII] ratio, which is taken from an empirical fit to the observed behavior of the [NII]/[OII] ratio in HII regions. Because the lines are well separated in wavelength, this method uses one line from each grism and hence it is highly dependent on the scaling and reddening correction.

We also used the ionization parameter, which was determined empirically (Kewley & Dopita 2002), as follows:

$$\log(\frac{I([OIII])}{I([OII])}) = k_0 + k_1\log(q) + k_2\log(q)^2 + k_3\log(q)^3,$$

with given coefficients for $12 + \log(O/H) = 8.9k_0 = 52.6367$, $k_1 = 16.088$, $k_2 = -1.67443$ and $k_3 = 0.0608004$.

#### 3.4.2. Method 2

Pilyugin et al. (2014) is a revised version of Pilyugin (2001), Pilyugin & Thuan (2005). They calibrated empirically the $R_3$...
and $P$ (excitation parameter) against the oxygen abundances determined with temperature-sensitive lines (direct $T_e$ method) of more than 250 HII regions from 3904 different spectra of 130 late-type galaxies, providing central oxygen abundances and abundance gradients. This is carried out for low metallicities ($12 + \log(O/H) \leq 8.0$) (Pilyugin 2000) and high metallicities ($12 + \log(O/H) \geq 8.3$). In this paper the relation for the higher metallicities (upper branch: $12 + \log(O/H) \geq 8.3$) is used.

The excitation parameter is defined as $P = R_3/R_{23}$ whereas

\[
R_3 = \frac{I([OIII]λ4959, 5007)/I(H\beta)}{I([OII]λ3726, 3729) + I([OIII]λ4959, 5007)}
\]

and

\[
R_{23} = \frac{I([OII]λ3726, 3729)}{I(H\beta)}
\]

Using this definition of the excitation parameter $P$ and $R_3$, they computed

\[
12 + \log(O/H) = 8.334 + 0.533P - (0.338 + 0.415P)\log R_3 - (0.086 - 0.225P)\log R_3^2.
\]

The oxygen abundance was determined with Eq. (6).

Limitations: in the intermediate-metallicity regime, $12 + \log(O/H) \sim 8.0-8.4$, the derived oxygen abundance is highly affected by the $R_{23}$ ambiguity, therefore the average value of the higher and lower metallicity regime should be assumed (López-Sánchez et al. 2012).

### 3.4.3. Method 3

Pettini & Pagel (2004) empirically calibrated the parameter O3N2 (Eq. (7)) against the oxygen abundances of a sample of 137 extragalactic HII regions determined mostly from temperature-sensitive lines (direct $T_e$ method) or photoionization models. This results in the oxygen abundance determination via a linear fit between $-1$ and 1.9, which is used in this paper (Eq. (8)), as follows:

\[
O3N2 = \log \left( \frac{I([OIII]λ5007)/I(H\beta)}{I([NII]λ6583)/I(H\alpha)} \right)
\]

\[
12 + \log(O/H) = 8.73 - 0.32 \cdot O3N2.
\]

Limitations: this method is only valid for $12 + \log(O/H) \gtrsim 8.7$ (López-Sánchez et al. 2012).

## 4. Results and discussion

Tables 4 and 5 show the results of the analysis. In this section, we discuss the Balmer decrement, compare the calibrations, and describe the results of the individual galaxies and their HII regions.

### 4.1. Balmer decrement of the HII regions

The measured Balmer decrement of $H\alpha/H\beta$ from the extraplanar HII region (3) of NGC 3628 is slightly lower than the theoretical value. Possibly, the region has a higher temperature. Nevertheless, considering the error, the Balmer decrement fits the theoretical value for all regions. The extinction parameter of the extraplanar HII region (3) of NGC 3628 is zero. Therefore, no reddening correction was applied. As the two extraplanar regions (2) and (3) of NGC 3628 are away from the galaxy disk, they are expected to be only slightly influenced by reddening, which is consistent with our observation.

We did a sanity check on region (3) of NGC 3628 to test how sensitive the oxygen abundance is to the uncertainty in $[OIII]$. The $[OIII]$, $H\alpha$, and $H\beta$ lines in grism 7 are very important for method 1, however grism 7 does not provide high S/N. We determined the influence of this line on the abundance measurement. We calculated the oxygen abundance with method 1 using the measured value of $H\alpha/H\beta = 2.84$ and additionally with considering the error of $0.31$, which leads to the upper limit of the $H\alpha/H\beta$ ratio of $3.15$. The latter ratio would lead to an extinction parameter of $c = 0.13$. That increases the $[OIII]$, $H\alpha$, and $H\beta$ line ratio from $1.71$ to $1.86$ and decreases $[NII]$, $H\alpha$, and $H\beta$ from $0.54$ to $0.49$. The resulting oxygen abundance of method 1 would change from $8.92$ to $8.87$. Thus, the oxygen abundance determination is robust toward our errors and the analysis is therefore reliable.

### 4.2. Oxygen abundances: general trends and comparison of the calibrations

The results of the abundance determination are consistent within all regions.

In general, method 1 (Kewley & Dopita 2002) shows the highest values of $12 + \log(O/H)$ while method 2 (Pilyugin et al. 2014) and method 3 (Pettini & Pagel 2004) provide lower values. The mean discrepancy between method 1 and method 2 is $0.41$ dex. The mean offset between method 1 and method 3 is $0.36$ dex and between method 2 and method 3 is $0.05$ dex. This result is in agreement with other analyses (Modjaz et al. 2008; Kewley & Ellison 2008).

To check the determined oxygen abundances, we searched for reference metallicities of the galaxies. For NGC 3628, an oxygen abundance of $12 + \log(O/H) = 8.57$ was determined by Engelbracht et al. (2008) using an integrated spectrum of the galaxy and empirical strong-line methods (Pettini & Pagel 2004; Pilyugin & Thuan 2005). To compare to their value, we also calculated the mean of the methods. For the disk HII region (1) of NGC 3628, the reference oxygen abundance is similar to the mean values of the two calibrations of $8.56$ dex.

The comparison of theoretical and empirical strong-line calibrations show well-known offsets (Stasińska 2002; Modjaz et al. 2008; Moustakas et al. 2010). We confirm the offsets with a mean difference of $0.37$ dex.

Strong-line calibrations use the fact that the metallicity in large HII regions is connected to the mean effective temperature of the stellar radiation field and the ionization parameter (Stasińska 2010). The theoretical calibrations use photoionization models to derive metallicities. These models are limited and depend strongly on the temperature structure calculations of the nebula (Kewley & Ellison 2008). This leads to discrepancies of up to $0.2$ dex within the theoretical calibrations and up to $0.6$ dex in comparison to empirical methods. Empirical calibrations use the relationship between direct metallicities and...
strong-line ratios of a sample of HII regions. This sample shows special characteristics in densities, the strength of the radiation field, and the state of ionization (Stasyszyn 2010). Therefore, empirical methods should be used with HII regions that are comparable to the calibration sample. One origin of discrepancies within various empirical methods is the variety of characteristics of HII regions. A possible reason for the offset in the comparison to theoretical calibrations is temperature fluctuations in HII regions. In metal-rich HII regions effective cooling processes could lead to strong temperature gradients (Garnett 1992). Therefore the temperature determined with the forbidden lines is overestimated and the oxygen abundance is underestimated (Garnett 1992). With knowledge of the fluctuations, the oxygen abundance would shift 0.2–0.3 dex (Moustakas et al. 2010). Another possible reason has been mentioned by Nicholls et al. (2012). If the electrons in HII regions do not follow a Maxwell-Boltzmann equilibrium energy distribution but follow a non-equilibrium kappa distribution, it should be possible to estimate the temperature and metallicity in these regions more accurately.

In addition to the calibrations used in this paper, Kewley & Ellison (2008) also investigated seven other calibrations and got a mean difference of up to 0.7 dex between theoretical and empirical calibrations. These authors propose to use oxygen abundance calibration with caution and to always specify the adopted calibration. They refer to relative oxygen abundances as a reliable tool, especially with the calibrations of Kewley & Dopita (2002) and Pettini & Pagel (2004).

4.3. Individual galaxies

In this section the individual galaxies with their extraplanar HII regions are presented. In interpreting our results, we account for the metallicity gradients in spiral galaxies (e.g., Garnett 1998; Sánchez-Menguiano et al. 2016), which can often change the metallicity by up to 0.3 dex within $R_{25}$ and does not depend on the different metallicity calculations (e.g., Magrini et al. 2011).

4.3.1. NGC 3628

NGC 3628 is a member of the Leo triplet and interacts with NGC 3627. The two analyzed extraplanar HII regions are located in a filament of NGC 3628 visible in HI (Wilding et al. 1993), which was probably formed due to tidal interactions. The regions have projected distances of 2.8 kpc (2) and 3 kpc (3).

Radial velocity: the determined radial velocities of the extraplanar HII regions are 648 km s$^{-1}$ of region (2) and 650 km s$^{-1}$ of region (3). Comparing this result to an HI velocity map of Wilding et al. (1993), the velocities are consistent. The same is true for the disk HII region (1). It is also obvious that the analyzed extraplanar regions (2), (3), and the disk region (1) are located in the same section of the velocity field. The extraplanar regions of NGC 3628 are located in a filament, which was probably formed by an interaction with NGC 3627. The velocity gradient within the filament as seen in the HI map is a hint that it has its origin in the outer parts of the galaxy. Therefore, the extraplanar regions and the disk region have comparable velocities within the rotation curve. The measured heliocentric velocities from our spectra confirm this (Table 2). There is a difference of only 20 km s$^{-1}$ between the disk region and the extraplanar regions. Therefore, the extraplanar regions are comparable with the disk region in terms of metallicity gradients.

Metallicities: the derived oxygen abundances of the HII disk region (1) and extraplanar region (3) are the same within the errors in all methods. In method 1 the extraplanar region (3) has a slightly lower value than the disk region (1). Nevertheless, in general, region (3) seems to have a comparable abundance to the disk region (1). We conclude that this extraplanar region was formed in a tidal arm of disk material.

Region (2) of NGC 3628 shows the lowest value in method 3. This value is confirmed by method 2. There is an indication that this region is highly ionized based on the high [OIII]/H$\beta$ ratio and the high ionization parameter in comparison to the other HII regions of this galaxy (see Table 4). In the diagnostic diagram (Fig. 6) the region shows the highest value of log(OIII/H$\beta$) of the HII regions analyzed in this paper but is still within the region of HII regions located.

4.3.2. NGC 4522

NGC 4522 is a member of the Virgo cluster. The optical distribution is undisturbed (Kenney & Koopmann 1999) while the HI distribution is asymmetric (Kenney et al. 2004). There is indication that ram pressure of the Virgo cluster stripped the gas of the galaxy (Kenney et al. 2004). The extraplanar HII region (1) is located in the outer gas and has a projected distance of 1.4 kpc.

Radial velocity: the determined heliocentric velocities of the three HII regions are around 2310 km s$^{-1}$. The extraplanar region (1) differs by 46 km s$^{-1}$ from disk region (2) and by

### Table 5. Resulting oxygen abundances of the different methods.

<table>
<thead>
<tr>
<th>Galaxy</th>
<th>Region</th>
<th>Method 1</th>
<th>Method 2</th>
<th>Method 3</th>
</tr>
</thead>
<tbody>
<tr>
<td></td>
<td></td>
<td>$[\pm 0.04]$</td>
<td>$[\pm 0.10]$</td>
<td>$[\pm 0.14]$</td>
</tr>
<tr>
<td>NGC 3628</td>
<td>1</td>
<td>8.96 ± 0.06</td>
<td>8.58 ± 0.11</td>
<td>8.72 ± 0.15</td>
</tr>
<tr>
<td></td>
<td>2 (e)</td>
<td>8.94 ± 0.07</td>
<td>8.43 ± 0.14</td>
<td>8.42 ± 0.21</td>
</tr>
<tr>
<td></td>
<td>3 (e)</td>
<td>8.92 ± 0.07</td>
<td>8.60 ± 0.11</td>
<td>8.67 ± 0.15</td>
</tr>
<tr>
<td>NGC 4522</td>
<td>1 (e)</td>
<td>8.92 ± 0.07</td>
<td>8.54 ± 0.13</td>
<td>8.63 ± 0.16</td>
</tr>
<tr>
<td></td>
<td>2</td>
<td>9.13 ± 0.04</td>
<td>8.68 ± 0.12</td>
<td>8.69 ± 0.15</td>
</tr>
<tr>
<td></td>
<td>3</td>
<td>9.12 ± 0.04</td>
<td>8.69 ± 0.11</td>
<td>8.72 ± 0.15</td>
</tr>
</tbody>
</table>

Notes. (e) = extraplanar, the error in the heading is the intrinsic error of the calibration according to the publication.
113 km s$^{-1}$ from disk region (3) in velocity (Table 2). A similar trend is seen in the velocity map of Kenney et al. (2004) at the locations of the regions, but with values that are, on average, lower by 30 km s$^{-1}$. One reason is that their central heliocentric velocity of 2337 km s$^{-1}$ is offset from other measurements that are lower, for example, from HI of de Vaucouleurs et al. (1991b) at 2324 $\pm$ 5 km s$^{-1}$. As the H$\alpha$ line lies at the edge of grism 9, we tested our measured heliocentric velocity using the $\lambda$5007 line, which lies in the middle of grism 9. We find on average 15 km s$^{-1}$ higher heliocentric velocities, which is more in agreement with Kenney et al. (2004) and indeed is compatible within the errors.

**Metallicities:** the extraplanar HII region (1) shows lower oxygen abundances in comparison to the disk HII regions (2) and (3) in all methods; the mean difference is 0.17 dex. In method 3 the offset between disk and extraplanar HII regions is the smallest and not significant, whereas the offset of 0.20 dex in method 1 is the largest and most significant. Method 2 shows a slightly lower value. Even considering the combined errors, this is not significant either. However, considering all methods together the lower metallicity is slightly more significant than for method 1 alone.

In the diagnostic diagram (Fig. 6), all regions are located close to each other. The Virgo Cluster galaxy NGC 4522 is close to peak ram pressure, where the outer gas disk has already been removed and the remaining gas disk is strongly truncated (Vollmer 2009). The slightly smaller oxygen abundance of the extraplanar HII region could be caused by the metallicity gradients that decrease the metallicity in the outer parts, which are the regions most affected by ram pressure stripping. We conclude that the extraplanar HII region (1) is located in these parts of stripped gas from the galaxy and therefore shows slightly lower metallicities.

**4.3.3. Origin of extraplanar HII regions**

The three analyzed extraplanar HII regions show different origins. These are tidal interactions in the galaxy NGC 3628 (tidal stream) and ram pressure in NGC 4522. These tidal interactions are both influenced by the group or cluster environment. As the extraplanar HII regions studied here show comparable or slightly lower oxygen abundances in comparison to the HII regions in the host galaxy, these extraplanar HII regions likely result from disk material. This result is different from the analysis of NGC 55 (Tüllmann et al. 2003; Kudritzki et al. 2016) in which two analyzed extraplanar HII regions show low oxygen abundance. Therefore, ejection from the disk was ruled out and Tüllmann et al. (2003) claimed that the regions were formed in the halo. The intergalactic regions of the Stephan’s Quintet (Mendes de Oliveira et al. 2004) show a high oxygen abundance. These regions were probably formed in tidal tales of the four interacting galaxies. The extragalactic region of NGC 4402 (Cortese et al. 2004) seems also to have a high metallicity and thus is built with pre-enriched material. They ruled out a tidal stripping scenario due to the absence of a clear interaction signal. One possible scenario would be that the extragalactic HII region is formed by an interaction between NGC 4402 and the ICM of the Virgo. The mechanism of ram pressure stripping is probably present in the region of NGC 4388 (Oosterloo & van Gorkom 2005). The same is true for the extraplanar HII region (1) of NGC 4522 analyzed in this paper, whereas the metallicity is low in the extraplanar region of NGC 4388 (Gerhard et al. 2002) and just slightly lower in our case.

**5. Summary and conclusions**

We present optical long-slit spectroscopic data of extraplanar HII regions and their edge-on host galaxies NGC 3628 and NGC 4522. In order to analyze these extraplanar HII regions, we determined metallicities (oxygen abundances). As the spectra contain strong emission lines such as the Balmer series, [OIII] $\lambda\lambda$3726, 3729, [OIII] $\lambda$4959, 5007, and [NII] $\lambda\lambda$6548, 6583, we calculate the oxygen abundances of the extraplanar HII regions with different strong-line calibrations based on empirical data and theoretical photoionization grids. First, we probed the shift between the different types of strong-line calibrations (empirical and theoretical). We analyzed the extraplanar HII regions and at least one HII region located in the disk of the...
galaxies via three different calibrations (Kewley & Dopita 2002; Pilyugin 2001; Pilyugin et al. 2014; Pettini & Pagel 2004). Our analysis again confirms the known offsets with a mean value of 0.37 dex. Secondly, we investigated the origin of the extraplanar H II regions by comparing the oxygen abundance of the extraplanar and disk H II regions.

In NGC 3628 we derived comparable oxygen abundances, which leads to the conclusion that the extraplanar H II regions are formed from disk material. The two analyzed extraplanar H II regions (2) and (3) of NGC 3628 are located in a filament that emerges from the disk, which was probably formed from tidal interactions with NGC 3627.

The extraplanar H II region (1) of NGC 4522 shows slightly lower measured oxygen abundances in comparison to the disk H II regions (2) and (3). Therefore, ram pressure of the Virgo Cluster probably stripped out the gas of the galaxy. This outer gas, where the extraplanar H II region is located, is not as metal rich as the inner disk region as a result of metallicity gradients. Thus, this extraplanar H II region (1) shows a slightly lower oxygen abundance than the inner disk.

Taking our results together with our previous work on NGC 55 (Tüllmann et al. 2003) and the results of other groups (see Sect. 1), it is clear that the mechanisms for forming extraplanar regions are diverse. The H II regions analyzed in this paper, along with most of the extraplanar H II regions mentioned in the literature, do not result from star formation in low-metallicity halos. The extraplanar H II regions seem largely to result from material pulled out during an interaction or from ram pressure stripping. Therefore, the environment has a strong impact on the generation of H II regions far above the disk plane.

We conclude that infalling gas, which is probably composed of falling back disk material and a small amount of intergalactic medium (e.g., Sancisi et al. 2008), does not form stars in the halo very frequently.

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