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First Run II Measurement of the W Boson Mass


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We describe a measurement of the W boson mass $m_W$ using 200 pb$^{-1}$ of $\sqrt{s}=1.96$ TeV $pp$ collision data taken with the CDF II detector. With a sample of 63,964 $W\to e\nu$ candidates and 51,128 $W\to \mu\nu$ candidates, we measure $m_W = (80.413 \pm 0.034\text{(stat)} \pm 0.034\text{(sys)}) = 80.413 \pm 0.048$ GeV/$c^2$. This is the single most precise $m_W$ measurement to date. When combined with other
I. INTRODUCTION

The discovery of the W and Z bosons in 1983 confirmed a central prediction of the unified model of electromagnetic and weak interactions [2]. Initial W and Z boson mass measurements verified the tree-level predictions of the theory, with subsequent measurements probing the predicted $O(3 \text{ GeV}/c^2)$ radiative corrections to the masses. The current knowledge of these masses and other electroweak parameters constrains additional radiative corrections from unobserved particles such as the Higgs boson or supersymmetric particles. These constraints are however limited by the precision of the W boson mass $m_W$, making improved measurements of $m_W$ a high priority in probing the masses and electroweak couplings of new hypothetical particles. We describe in this article the single most precise $m_W$ measurement to date.

The W boson mass can be written in terms of other precisely measured parameters in the “on-shell” scheme as [3, 4]:

$$m_W^2 = \frac{\hbar^2 \pi \alpha_{EM}}{c^2} \frac{Q^{\alpha_{EM}}}{\sqrt{2} G_F (1 - m_W^2/m_Z^2)(1 - \Delta r)},$$

where $\alpha_{EM}$ is the electromagnetic coupling at the renormalization energy scale $Q = m_Zc^2$, $G_F$ is the Fermi weak coupling extracted from the muon lifetime, $m_Z$ is the Z boson mass, and $\Delta r$ includes all radiative corrections. Fermionic loop corrections increase the W boson mass by terms proportional to $\ln(m_Z/m_f)$ for $m_f \ll m_Z$ [4], while the loop containing top and bottom quarks (Fig. 1) increases $m_W$ according to [2]:

$$\Delta r_{tb} = \frac{c}{\hbar^2} \frac{-3G_F m_{tb}^2}{8\sqrt{2} \pi^2 (m_Z^2 - m_W^2)} \times \left[ \frac{m_t^2 + m_b^2}{m_t^2 - m_b^2} \ln\left(\frac{m_t^2}{m_b^2}\right) \right],$$

where the second and third terms can be neglected since $m_t \gg m_b$. Higgs loops (Fig. 2) decrease $m_W$ with a contribution proportional to the logarithm of the Higgs mass ($m_H$). Contributions from possible supersymmetric particles are dominated by squark loops (Fig. 3) and tend to increase $m_W$. Generally, the lighter the squark masses and the larger the squark weak doublet mass splitting, the larger the contribution to $m_W$. The total radiative correction from supersymmetric particles can be as large as several hundred MeV/$c^2$ [7].

Table I [8] shows the change in $m_W$ for +1σ changes in the measured standard model input parameters and the effect of doubling $m_H$ from 100 GeV/$c^2$ to 200 GeV/$c^2$. In addition to the listed parameters, a variation of ±1.7 MeV/$c^2$ on the predicted $m_W$ arises from two-loop sensitivity to $\alpha_s$, e.g. via gluon exchange in the quark loop in Fig. 1. Theoretical corrections beyond second order, which have yet to be calculated, are estimated to affect the $m_W$ prediction by ±4 MeV/$c^2$ [8].

\[ \text{FIG. 1: The one-loop contribution to the } \text{W boson mass from top and bottom quarks.} \]

The uncertainties on the $m_W$ prediction can be compared to the 29 MeV/$c^2$ uncertainty on the world average from direct $m_W$ measurements (Table I),
which include results from four experiments, ALEPH [12], DELPHI [15], L3 [14], and OPAL [13], studying \( \sqrt{s} = 161 - 209 \) GeV \( e^+e^- \) collisions at the Large Electron Positron collider (LEP), and from two experiments, CDF [16] and DØ [17, 18], studying \( \sqrt{s} = 1.8 \) TeV \( p\bar{p} \) collisions in Run I of the Fermilab Tevatron. The current experimental \( m_W \) uncertainty is a factor of two larger than the uncertainty from radiative corrections, excluding the Higgs contribution (Table I). The Higgs mass constraint extracted from the \( W \) boson mass is thus limited by the direct \( m_W \) measurement. The precise \( m_W \) measurement described in this article has a significant impact on the world average \( m_W \).

### II. OVERVIEW

A measurement of \( m_W \) at a \( p\bar{p} \) collider [21] is complementary to that at an \( e^+e^- \) collider. Individual \( u \) (\( d \)) quarks inside the proton can interact with \( \bar{d} \) (\( \bar{u} \)) quarks inside the anti-proton (or vice versa), allowing single \( W^+ (W^-) \) boson production, which is not possible at an \( e^+e^- \) collider. In addition, \( p\bar{p} \) colliders have higher center of mass energies and \( W \) boson production cross sections. This provides high statistics for the leptonic decays of the \( W \) boson, which are studied exclusively because of the overwhelming hadronic-jet background in the quark decay channels.

The leptonic decays of singly produced \( Z \) bosons provide important control samples, since both leptons from \( Z \) boson decay are well measured. The production and decay uncertainties on the measurement of \( m_W \) from \( p\bar{p} \) and \( e^+e^- \) collider data are almost completely independent [22].

We present in this Section an overview of \( W \) and \( Z \) boson production at the Tevatron, a description of the coordinate definitions and symbol conventions used for this measurement, and a broad discussion of our \( m_W \) measurement strategy.

#### A. \( W \) and \( Z \) Boson Production and Decay

\( W \) and \( Z \) bosons are produced in \( \sqrt{s} = 1.96 \) TeV \( p\bar{p} \) collisions primarily through \( s \)–channel annihilation of valence \( u \) and/or \( d \) quarks (Fig. 4), with a smaller \( O(20\%) \) contribution from sea quarks. The quark (antiquark) has a fraction \( x_p \) (\( x_{\bar{p}} \)) of the pro-
ton’s (antiproton’s) total momentum, producing a $W$ or $Z$ boson at center of mass energy $\sqrt{s} \equiv Q$ equal to its mass times $c^2$. The rate of production can be predicted from two components: (1) the momentum fraction distributions of the quarks, $f_q(x, Q^2)$, which are determined from fits to world data [23, 24]; and (2) a perturbative calculation of the $qq^\prime \to W$ or $Z$ boson process [25].

$W$ and $Z$ bosons can decay to lepton or quark pairs. Decays to quark pairs are not observable given the large direct $qq^\prime$ background, and decays to $\tau \to \nu_\tau$ +hadrons are not as precisely measured as lepton decays to electrons or muons. For these reasons we restrict ourselves to the direct electronic and muonic decays ($W \to e\nu$, $W \to \mu\nu$, $Z \to ee$, and $Z \to \mu\mu$), with the corresponding decays to $\tau \to$ leptons considered as backgrounds to these processes (Section VIII). The branching ratio for each leptonic decay $W \to l\nu$ ($Z \to ll$) is $\approx 11\%$ (3.3\%), and the measured cross section times branching ratio is $(2749 \pm 174)$ pb [$(254.9 \pm 16.2)$ pb] [26].

**B. Conventions**

We use both Cartesian and cylindrical coordinate systems, in which $+z$ points in the direction of the proton beam (east) and the origin is at the center of the detector. In the right-handed Cartesian coordinate system, $+x$ points north (outward from the ring) and $+y$ points upwards; in the cylindrical system, $\phi$ is the azimuthal angle and $r$ is the radius from the center of the detector in the $x - y$ plane. The rapidity $y = -\ln[(E - p_zc)/(E + p_zc)]$ is additive under Lorentz boosts along the $z$ axis. For massless particles, this quantity is equal to the pseudorapidity $\eta = -\ln[tan(\theta/2)]$, where $\theta$ is the polar angle with respect to the $z$ axis. All angles are quoted in radians unless otherwise indicated.

Because the interacting quarks’ longitudinal momenta $p_z$ are not known for each event, we generally work with momenta transverse to the beam line. The interacting protons and antiprotons have no net transverse momentum. Electron energy (muon momentum) measured using the calorimeter (tracker) is denoted as $E$ $(p)$, and the corresponding transversemomenta $p_T$ are derived using the measured track direction and neglecting particle masses. The event calorimetric $p_T$, excluding the lepton(s), is calculated assuming massless particles using calorimeter tower energies (Section [III A 2]) and the lepton production vertex, and provides a measurement of the recoil momentum vector $u$. The component of recoil projected along the lepton direction is denoted $u_{||}$ and the orthogonal component is $u_{\perp}$ (Fig. 5). The transverse momentum imbalance in a $W$ boson event is a measure of the neutrino transverse momentum $p_T^\nu$ and is given by $p_T^\nu = -(\vec{p}_T - \vec{u}_T)$, where $\vec{p}_T$ is the measured charged lepton transverse momentum.

When electromagnetic charge is not indicated, both charges are considered. We use units where $\hbar = c = 1$ for the remainder of this paper.

**C. Measurement Strategy**

The measurement of the final state from $W \to l\nu$ decays involves a measurement of $p_T^l$ and the total recoil $\vec{u}_T$. The neutrino escapes detection and the unknown initial partonic $p_z$ precludes the use of $p_z$ conservation in the measurement. The boson invariant mass is thus not reconstructable; rather, the 2-
dimensional “transverse mass” $m_T$ is used in the $m_W$ fit:

$$m_T = \sqrt{2p_T^1 p_T^2 (1 - \cos \Delta \phi)}$$  \hspace{1cm} (3)

where $\Delta \phi$ is the angle in the transverse plane between the leptons, whose masses are negligible. The fit to the $m_T$ distribution provides the statistically most precise measurement of $m_W$.

The charged lepton, which can be measured precisely, carries most of the observable mass information in the event. We calibrate the muon momentum precisely, which carries most of the observable mass information as well. This results in a precise track momentum calibration, which we transfer to the calorimeter with a fit to the ratio of calorimeter energy to track momentum $(E/p)$ of electrons from $W$ boson decays. The accuracy of these calibrations is demonstrated by applying them to measurements of the $Z$ boson mass in the muon and electron decay channels. We then incorporate the known $Z$ boson mass as an additional calibration constraint.

The other directly measurable quantity needed for the calculation of $m_T$ is the recoil transverse momentum $\vec{u}_T$. Since the $W$ and $Z$ bosons are produced at a similar $Q^2$, they have similar recoil distributions. We use the leptons from the $Z$ boson decay to measure the $p_T$ of the $Z$ boson. We then calibrate our model of $\vec{u}_T$ by measuring the balance between the recoil and $Z$ boson $p_T$. The $Z$ boson statistics are sufficient to perform a recoil calibration to 1% accuracy, which leads to a systematic uncertainty commensurate with other uncertainties on $m_W$.

To accurately model the shape of the $m_T$ distribution, we use a fast Monte Carlo simulation of the $p\bar{p} \rightarrow W \rightarrow l\nu$ process including the recoil and detector response. The custom fast simulation allows flexibility in parametrizing the detector response and in separating the effects of the detector model components. We use a binned likelihood to fit the measured $m_T$ distributions to templates (Section III D), generated from the fast simulation, with $m_W$ as the free parameter. All $m_W$ and lepton energy scale fits are performed with this procedure.

Though less statistically precise, the $p_T^1$ and $p_T^2$ distributions provide additional information on the $W$ boson mass and are used as important tests of consistency. We separately fit these distributions for $m_W$ and combine all fits in our final result.

During the measurement process, all $W$ boson mass fits were offset by a single unknown random number chosen from a flat distribution in the range $[-100,100]$ MeV. The fit result was thus blinded to the authors until the analysis was complete [27]. The final measured $m_W$ and its uncertainty have not changed since the random offset was removed from the fit results.

We give a brief overview of the template likelihood fitting procedure in Section III D. Section III describes the detector and the fast detector simulation used in the analysis. The $W$ boson measurement samples are defined in Section V. We describe the precision measurements of muons and electrons in Sections VI and VII respectively. These sections include event selection, calibration, and resolution studies from the dilepton and $W$ boson data samples. Measurement of the recoil response and resolution is presented in Section VII. The backgrounds to the $W$ boson sample are discussed in Section VIII. Theoretical aspects of $W$ and $Z$ boson production and decay, including constraints from the current data sample, are described in Section IX. We present the $W$ boson mass fits and cross-checks in Section X. Finally, in Section XI we show the result of combining our measurement with previous measurements, and the corresponding implications on the predicted standard model Higgs boson mass.

D. Template Likelihood Fits

All the fits involving mass measurements and the energy scale (Sections V, VI, and VII) are performed with a template binned likelihood fitting procedure. A given distribution to be fit is generated as a discrete function of the fit parameter, using the fast simulation. These simulated distributions are referred to as “templates.” For each value of the fit parameter, the simulated distribution is compared to the data distribution and the logarithm of a binned likelihood is calculated. The binned likelihood is the Poisson probability for each bin to contain the $n_i$ observed data events given $m_i$ expected events, multiplied over the $N$ bins in the fit range:

$$L = \prod_{i=1}^{N} \frac{e^{-m_i} m_i^{n_i}}{n_i!}.$$  \hspace{1cm} (4)

We calculate the logarithm of the likelihood using the approximation $\ln n! \approx (n + 1/2) \ln(n + 1) - n$:

$$\ln L \approx \sum_{i=1}^{N} [n_i \ln m_i - m_i - (n_i + 1/2) \ln(n_i + 1) + n_i].$$  \hspace{1cm} (5)

The best-fit value of the parameter maximizes the likelihood (or equivalently minimizes $-\ln L$), and the $\pm 1\sigma$ values are those that increase $-\ln L$ by $1/2$. The
approximation for $\ln n!$ only affects the shape of the likelihood about the minimum and not the position of the minimum. The procedure is validated by fitting simulated data (“pseudoexperiments”) and no bias is found. We symmetrize the uncertainty by taking half the difference between the $+1\sigma$ and $-1\sigma$ values. For the $E/p$ fits in the $W$ boson sample, we reduce the effect of finite template statistics by fitting $-\ln \mathcal{L}$ to a parabola, and extracting the best-fit value and the uncertainty from this parabola.

III. DETECTOR AND MODEL

The CDF II detector [26, 28] is well suited for the $m_W$ measurement. Its high-resolution tracker and calorimeter measure individual charged lepton momenta from $W$ and $Z$ boson decays with a resolution of $\approx 2\%$. It has similar acceptance and resolution for central electrons and muons, giving the two channels similar weight in a combined mass measurement.

A. Detector Components

The CDF II detector (Fig. 6) is a multi-purpose detector consisting of: an inner silicon tracker designed to measure the production vertex of charged particles with high precision; an outer tracking drift chamber to measure charged particle momenta; a solenoid to provide a uniform 1.4 T magnetic field inside the trackers; electromagnetic calorimeters to contain and measure electron and photon showers; hadronic calorimeters for hadron energy measurements; and a muon system to detect muons escaping the calorimeters. The detector information is read out on-line and saved for later analysis when event topologies consistent with a particular physics process (or class of processes) are selected. The read-out decision is made with a fast three-level trigger system that has high efficiency for selecting the $W$ and $Z$ bosons to be used in the offline analysis.

1. Tracking System

The silicon tracker (Fig. 7) consists of three separate detectors: Layer 00, SVX II, and ISL. Layer 00 is a single layer of 300 $\mu$m thick sensors attached to the beam pipe at a radius of 1.3 cm. Five additional layers of sensors at radii ranging from 2.5 cm to 10.6 cm comprise SVX II. Surrounding these sensors are port cards, which transport deposited charge information from the silicon wafers to the readout system. The intermediate silicon layers (ISL) are located at radii of 20.2 cm and 29.1 cm. The SVX II is segmented longitudinally into three barrels in the region $|z| < 45$ cm. This covers the $p\bar{p}$ interaction region, which is well approximated by a Gaussian distribution with $\sigma_z \approx 30$ cm. We do not use the silicon measurements in this analysis, though we model the tracker’s effects on leptons and photons (Section II.B).

An open-cell drift chamber, the central outer tracker (COT) [29], surrounds the silicon tracker and covers the region $|z| < 155$ cm ($|\eta| \lesssim 1$) and $40 \text{ cm} < r < 137$ cm. The COT consists of eight concentric ‘superlayers,’ separated azimuthally into cells. Each cell contains 12 sense wires to measure the ionization produced by a charged particle in the ambient argon-ethane gas mixture. The superlayers alternate between a purely axial configuration, with sense wires parallel to the beam line, and a small-angle stereo configuration, with sense wires at a 2° angle relative to the $z$ axis.

The sense wires are strung from end to end in $z$ and held under tension at each aluminum endplate (Fig. 8). The wires are azimuthally sandwiched by field sheets, which provide a 1.9 kV/cm electric field. All cells are rotated at a 35° angle relative to a radial line, such that the ionized electrons travel approximately azimuthally to the wire under the combined influence of the local electric field and the global magnetic field from the solenoid.

Within a given cell the sense wires are slightly off-center relative to the field sheets. In addition, the sense wires and field sheets sag under the influence of gravity, with the field sheets sagging more due to their larger masses. These effects cause a small electrostatic deflection of the sense wires toward a particular field sheet. To prevent the relative deflection of sense wires within a cell, a support rod connects the sense wires at the center of the detector. The support rod results in a small ($\approx 2$ mm) region at $z = 0$ cm where charged particles are not measured.

Between the solenoid and the COT is a time-of-flight system (TOF) consisting of scintillator bars that precisely measure the time of incidence of charged particles. From this measurement and the tracker information, a particle’s velocity and mass can be inferred. The TOF is not utilized in this analysis.

2. Calorimeter System

The CDF calorimeter is segmented radially into electromagnetic and hadronic sections. The central
FIG. 6: A cut-away view of a section of the CDF detector. The slice is along the y-axis at \( x = 0 \) cm.

calorimeter covers \(|\eta| < 1.1\) and is split at the center into two separate barrels covering +\( \eta \) and −\( \eta \). Each barrel consists of 24 azimuthal “wedges” of size 0.26 radians (15°) with ten projective towers of size \( \Delta \eta \approx 0.11\). To allow a pathway for the solenoid cryogenic tubes, a two-tower region is removed, corresponding to \( 0.77 < \eta < 1.0, 75^\circ < \phi < 90^\circ, \) and \( z > 193 \) cm. The forward calorimeter covers \( 1.1 < |\eta| < 3.6, \) filling the forward gaps with a plug shape (Fig. 6).

The central electromagnetic calorimeter (CEM) \[30\] has a thickness of 18 radiation lengths, consisting of 31 radial scintillator layers interleaved with 30 layers of lead-aluminum plates. At a radius of 184 cm electromagnetic showers have traversed about six radiation lengths (including the solenoidal coil) and have their maximum energy deposition. At this radius finely segmented strip and wire chambers (CES) measure the energy deposition with a position resolution of \( \approx 2 \) mm.

The local shower position in the azimuthal direction in the tower is denoted as CES \( x \), which ranges from -24.1 cm to 24.1 cm. The wire chambers extend only to \(|x| \leq 22.5\) cm, and for \(|x| > 23.1\) cm no energy measurements are made. In this region wavelength shifters read out the light from the scintillator, and steel and foam separate the towers. Light follows a waveguide to a phototube positioned at the back end of the hadronic calorimeter.

Parallel to the beam line, the position at shower maximum is denoted CES \( z \). The strip chambers extend from 6-239 cm in \(|z|\), and there is no scintillator for \(|z| < 4.2\) cm, where the two calorimeter barrels meet.

The central hadronic calorimeter \[31\] is separated into a central region (CHA, \(|\eta| < 0.6\)) with 32 longitudinal layers of scintillator sandwiched with steel and a forward “wall” calorimeter (WHA, \(0.6 < |\eta| < 1.1\)) with 15 such layers. These calorimeters have thicknesses of \( \approx 4.5\) interaction lengths.

The plug calorimeter \[32\] has a comparable design to the central calorimeter with scintillator-lead electromagnetic calorimeters and scintillator-steel hadronic calorimeter compartments. The \( \phi \) seg-
FIG. 7: End view of the silicon detector. The innermost layer (Layer 00) is attached to the beam pipe, and is surrounded by five concentric layers of silicon wafers (SVX II). The outermost layers are the intermediate silicon layers (ISL), which sit just inside the outer tracking chamber.

3. Muon Detectors

The muon systems relevant for the W mass measurement cover the region $|\eta| \leq 1$. The central muon detector (CMU) and the central muon upgrade (CMP) cover $|\eta| \leq 0.6$, while the central muon extension (CMX) covers $0.6 \leq |\eta| \leq 1$.

The CMU detector is located at the outer edge of the CHA, 347 cm from the z axis. The CMU is segmented into 15° azimuthal wedges containing four layers of proportional drift chambers that cover 12.6°. The maximum drift time within a chamber is 800 ns, about twice as long as the 396 ns spacing between $p\bar{p}$ crossings. CMU information must therefore be combined with reconstructed COT particle tracks to determine the appropriate $p\bar{p}$ crossing.

Because the total thickness of the central calorimeter is about five interaction lengths, approximately 0.5% of high-momentum pions reach the CMU. To reduce this background, the CMP detector is located behind an additional 60 cm of steel. The CMP has a similar construction to the CMU, with the exception that wider drift chambers are used to cover the same solid angle, resulting in a maximum drift time of 1.8 µs rather than 800 ns.

The CMX detector consists of eight drift chamber layers beyond both the calorimeter and the steel detector support structure (6 – 10 interaction lengths). The CMX $\phi$ regions used in this analysis are $-45^\circ < \phi < 75^\circ$ and $105^\circ < \phi < 225^\circ$. New detectors for Run II cover much of the remaining $\phi$ region, but were not fully commissioned for the data-taking period of this analysis. Scintillator detectors (CSX) at the inner and outer surfaces of the CMX provide timing information to the trigger to separate collision particles from other sources such as beam halo or cosmic rays.

4. Trigger System

The trigger consists of three stages with progressively greater sophistication of event reconstruction. The first stage is hardware-based, the second a mix of hardware and software, and the third a farm of processors performing full event reconstruction.

The first trigger stage, level 1, includes tracker, calorimeter, and muon reconstruction. The charged particle track reconstruction is performed with the extremely fast tracker (XFT) based on the four axial COT superlayers. A track segment is reconstructed in a given superlayer if at least 11 of the 12 sense wires in a wide road have charge deposition above a given threshold (“hits”). The list of segments from the full tracker is compared to predefined groups of segments expected from charged particles above a given momentum threshold. When matches are found, track candidates are created and passed to the track extrapolator (XTRP). The XTRP determines the expected $\phi$ positions of the tracks in the calorimeter and muon detectors, for the purpose of forming electron and muon candidates.

The calorimeter reconstruction at level 1 defines separate electromagnetic and hadronic “trigger towers” as tower pairs adjacent in $\eta$. The tower $p_T$ is calculated assuming a collision vertex $z = 0$ and an electron candidate is formed if the ratio of hadronic to electromagnetic energy (Had/EM) in a trigger tower is less than 1/8. The high-momentum electron trigger used in this analysis requires a level 1 trigger tower with electromagnetic $p_T > 8$ GeV matched to a track with $p_T > 8$ GeV, and drops the Had/EM requirement for electromagnetic $p_T > 14$ GeV.

Level 1 muon reconstruction includes a $p_T$ estimate within the CMU and CMX chambers from the relative timing of the hits in different layers. The CMU track segments are combined with reconstructed CMP track segments to create “CMUP”
muon candidates. For the majority of the data CMX candidates also require local CSX hits consistent with particles originating from the collision. For our $W$ and $Z$ boson samples we use a muon trigger that requires CMU or CMX $p_T > 6$ GeV matched to an XFT track with $p_T > 4$ GeV (CMUP) or $p_T > 8$ GeV (CMX).

The level 2 calorimeter reconstruction uses a more sophisticated clustering algorithm for electromagnetic objects. This improves energy measurement resolution and allows a higher threshold ($p_T > 16$ GeV) to be applied. To reduce rates, the XFT track requirement for CMUP candidates was raised to $p_T > 8$ GeV for most of the data-taking period.

At level 3, approximately 300 dual processor computers allow full track pattern recognition, muon reconstruction, and calorimeter clustering. Variables used to select electrons at level 3 are the lateral shower profile, $L_{shr}$ (Section IV B), and the distance between CES $z$ and the $z$-position of the track extrapolated to the CES ($\Delta z$). The $L_{shr}$ variable quantifies the difference between the measured energies of towers neighboring the electron in $\eta$ and the expected energies determined from electron test beam data. The trigger requirements of $L_{shr} < 0.4$ and $|\Delta z| < 8$ cm are $\approx 100\%$ efficient for electrons from $W$ and $Z$ boson decays. The high-momentum electron trigger also requires electromagnetic $p_T > 18$ GeV and track $p_T > 9$ GeV. For efficiency studies we use a separate trigger that requires electromagnetic $p_T > 25$ GeV and $p_T^{L3} > 25$ GeV, but has no quality requirements at level 3 and no trigger track requirements. At level 3, $p_T^{L3}$ is defined as the negative of the vector sum of the transverse momenta of the calorimeter towers. The high-momentum muon trigger requires a COT track with $p_T > 18$ GeV matched to a CMUP or CMX track segment.

5. Luminosity Detector

The small-angle Cherenkov luminosity counters (CLC) are used to measure the instantaneous and integrated luminosity of our data samples. The CLC consists of two modules installed around the beampipe at each end of the detector, providing coverage in the regions $3.6 < |\eta| < 4.6$. Each module contains 48 conical gas Cherenkov counters pointing to the collision region. Signals in both CLC modules coinciding in time with the bunch crossing are used to measure the instantaneous luminosity and to trigger collision events. Events collected with this trigger, known as “minimum bias” events, are used to study the detector response to generic inelastic $p\bar{p}$ collisions.
B. Detector Model

We use a parametrized model of the detector response to electrons, muons, and the hadronic recoil. The model is incorporated into a custom fast simulation that includes lepton and recoil reconstruction, event selection, and fit template generation. The simulation provides both flexibility in determining the effects of various inputs, and computing speed to allow frequent high-statistics studies. A sample of $O(10^7)$ events can be generated using a single-processor machine in one day. This is several orders of magnitude more than the $O(10^5)$ events that can be produced with the standard GEANT-based CDF simulation [40, 41].

We describe in this section the simulation of electrons and muons. Fits to the data that determine the values of some of the model parameters are described in Secs. [V] and [VI]. The detector model of hadronic recoil response and resolution is discussed in Sec. [VII]

The model components common to muons and electrons are: ionization energy loss and multiple scattering in the beam pipe and tracker volume; parametrized track hit resolutions and efficiencies; and track reconstruction. We describe these components in the muon simulation overview, and then discuss the electron- and photon-specific simulation.

1. Muon Simulation

Muon and electron tracks are reconstructed using only COT hit and beam position information (Section IV). Thus, the simulation of the silicon detector consists entirely of energy loss and multiple scattering. In the COT, hit resolutions and efficiencies are additionally simulated, and track reconstruction is performed. The total measured muon EM calorimeter energy is simulated by combining the minimum-ionizing energy deposition with energy from final-state photon radiation (Section IX D) and the recoil and underlying event [39]. Finally, the detector fiducial volume we calculate the energy loss between the COT and beam position information (Appendix A), for each of 32 radial layers. Except for the inner and outer layers, the map is finely segmented longitudinally and in azimuth to capture the material variation in the silicon detector [42]. Inside the COT fiducial volume we calculate the energy loss between each of the 96 radial sense wires.

The energy loss model is tuned using the data. We apply a global correction factor of 0.94 to the calculated energy loss in the material upstream of the COT ($r < 40$ cm), we use a three-dimensional lookup table of the material properties of the beam pipe, the silicon detector, and COT affects the resolution of the decay muons (Section V B 3).

Ionization Energy Loss

The differential ionization energy loss of muons and electrons in the tracking system is simulated according to the Bethe-Bloch equation [11]:

$$-rac{dE}{dx} = \frac{KZ}{A\beta^2} \left[\frac{1}{2} \ln \frac{2m_e\beta^2I_{\text{max}}}{(1-\beta^2)I^2} - \beta^2 - \frac{\delta}{2}\right],$$

(6)

where $K = 4\pi N_A r_e^2 m_e$, $N_A$ is Avogadro’s number, $r_e$ is the classical electron radius, $Z(A)$ is the atomic (mass) number, $\beta$ is the particle velocity, $I$ is the mean excitation energy, $I_{\text{max}}$ is the maximum kinetic energy that can be transferred to a free electron in a single collision, and $\delta$ is the material-dependent density effect as a function of $\beta$ [11]. When calculating the effect of $\delta$, we take the material to be silicon throughout.

To calculate muon energy loss in the material upstream of the COT ($r < 40$ cm), we use a three-dimensional lookup table of the material properties of the beam pipe, the silicon detector, and the wall of the aluminium can at the inner radius of the COT. The lookup table determines the appropriate $Z/A$ and $I$ values, along with the radiation length $X_0$ (Appendix A), for each of 32 radial layers. Except for the inner and outer layers, the map is finely segmented longitudinally and in azimuth to capture the material variation in the silicon detector [42]. Inside the COT fiducial volume we calculate the energy loss between each of the 96 radial sense wires.

The energy loss model is tuned using the data. We apply a global correction factor of 0.94 to the calculated energy loss in the material upstream of the COT in order to obtain a $J/\psi \rightarrow \mu \mu$ mass measurement that is independent of the mean inverse momentum of the decay muons (Section V B 3).

Multiple Coulomb Scattering

Multiple Coulomb scattering in the beampipe, silicon detector, and COT affects the resolution of the reconstructed track parameters for low-momentum tracks. We model the scattering using a Gaussian distribution for 98% of the scatters [43] with an angular resolution $\sigma_\theta$ defined by

$$\sigma_\theta = \frac{13.6 \text{ MeV}}{\beta p} \sqrt{x/X_0},$$

(7)

where $x$ is the thickness of the layer and $X_0$ is the layer’s radiation length (Section I I I B 2). Simulation of multiple scattering is implemented for each radial
layer of the three-dimensional lookup table and between each COT layer.

Based on the results of low-energy muon scattering data, we model the non-Gaussian wide-angle scatters by increasing $\sigma_\theta$ by a factor of 3.8 for 2% of the scatters.

COT Simulation and Reconstruction

The charged track measurement is modeled with a full hit-level simulation of the charge deposition in the COT and a helical track fit. The parameter resolution of reconstructed tracks is affected by the individual hit resolution, and by the distribution of the number of hits ($N_{hit}$) used in the fit.

We tune the COT hit resolution using the width of the $T \rightarrow \mu\mu$ mass distribution reconstructed with non-beam-constrained tracks. The tuned value of $[150 \pm 3(\text{stat})] \mu$m is consistent with the $149 \mu$m RMS of the observed hit residual distribution for the muon tracks in $Z \rightarrow \mu\mu$ data. We use a $150 \mu$m hit resolution for the simulation of the $T$, $W$, and $Z$ bosons.

We use a dual-resolution model to describe the narrower mass peak in the high-statistics $J/\psi \rightarrow \mu\mu$ sample, where the muons generally have lower momenta than the other samples. The $J/\psi$ mass peak width is particularly sensitive to multiple scattering and relative energy loss, and our hit-resolution model compensates for any mismodeling that affects the peak width. We find that a single-hit resolution of $155 \mu$m applied to $70\%$ of the tracks and $175 \mu$m applied to the remaining $30\%$ adequately describes the width and lineshape of the $J/\psi \rightarrow \mu\mu$ mass peak.

To describe the $N_{hit}$ distribution, we use a dual-hit-efficiency model, the larger one applied to the majority of the tracks. The lower efficiency accounts for events with high COT occupancy, where fewer hits are attached to reconstructed tracks. The two parameters are tuned to match the mean and RMS of the data $N_{hit}$ distributions. We independently tune these parameters for the $J/\psi$ sample, the $Y$ sample, and the $W$ and $Z$ boson samples.

COT hit positions from a charged track are used to reconstruct a helix with a $\chi^2$-minimization procedure. The axial helix parameters are the impact parameter with respect to the nominal beam position, $d_0$, the azimuthal angle at the closest approach to the beam, $\phi_0$, and the curvature of the track, $c$, defined to be $(2R)^{-1}$, where $R$ is the radius of curvature. The stereo helix parameters are the longitudinal position at the closest approach to the beam, $z_0$, and the cotangent of the polar angle, $\cot \theta$.

When optimizing resolution of lepton tracks from prompt resonance decays, we constrain the helix to originate from the location of the beam. The transverse size of the beam is $\approx 30 \mu$m at $z = 0$ cm and increases to $50 - 60 \mu$m at $|z| = 40$ cm. For simplicity we assume an average beam size of $[39 \pm 3(\text{stat})] \mu$m, which is determined from a fit to the width of the $Z \rightarrow \mu\mu$ mass peak. The beam constraint improves the intrinsic fractional momentum resolution by about a factor of three, to $\delta p_T/p_T \approx 0.0005 p_T/\text{GeV}$.

We perform a track fit on our simulated hits in the same manner as the data. The hits are first fit to a helix without a beam constraint; hits with large residuals (> $600 \mu$m) are dropped from the track (in order to remove spurious hits added in data pattern recognition); and the track is fit again with an optional beam constraint. This option is applied to prompt lepton tracks from $W$ and $Z$ boson decays, but not to tracks from $J/\psi$ decays, approximately $20\%$ of which are not prompt. The prompt muons from $Y$ decays are fit twice, both with and without the beam constraint, as a consistency check.

Calorimeter Response

Muons deposit ionization energy in the calorimeter. We simulate a muon’s EM energy deposition using a distribution taken from cosmic ray muons passing through the center of the detector, in events with no other track activity. An additional contribution comes from energy flow into the calorimeter from the underlying event. We model this energy using a distribution taken from cosmic ray muons passing through the center of the detector, in events with no other track activity. We use the map in the fast simulation to determine the fiduciality of a muon at the CES (Section VII B).

Detector Fiduciality

The CMUP and CMX muon systems do not have complete azimuthal or polar angle coverage. We create an $\eta - \phi$ map of each muon detector’s coverage using muons simulated with a detector geometry based on GEANT 4. We use the map in the fast simulation to determine the fiduciality of a muon at...
a given $\eta - \phi$ position.

We incorporate the relative efficiency of the CMUP to CMX triggers in the fast simulation by matching the ratio of CMUP to CMX events in the $W \rightarrow \mu\nu$ data (Section IV A).

2. Electron and Photon Simulation

The dominant calibration of the calorimeter energy measurement \(E\) of electrons uses their track momenta \(p\) and a fit to the peak of the \(E/p\) distribution. An additional calibration results from a mass fit to the \(Z\) boson resonance and reduces the calibration uncertainty by 20% relative to the \(E/p\) calibration alone.

The \(E/p\) method relies on an accurate modeling of radiative effects that reduce the track momentum measured in the COT. A given electron loses $\approx 20\%$ of its energy through bremsstrahlung radiation in the silicon detector, and this process has the most significant impact on the \(E/p\) calibration. The total amount of silicon detector material is tuned with data using highly radiative electrons (Section VI A). We additionally model processes that affect the shape of the \(E/p\) distribution: photon conversion in the tracker; energy loss in the solenoid and the time-of-flight system; electromagnetic calorimeter response and resolution; and energy loss into the hadronic calorimeter. The models of ionization energy loss and multiple scattering in the tracker, as well as the COT track simulation and reconstruction, are the same as for muons (Section III B 1).

Bremsstrahlung

The differential cross section for an electron of energy \(E_e\) to radiate a photon of energy \(E_\gamma\) is given by the screened Bethe-Heitler equation [48] over most of the \(y \equiv E_\gamma/E_e\) spectrum. In terms of the material’s radiation length \(X_0\), the differential cross section for bremsstrahlung radiation is:

\[
\frac{d\sigma}{dy} = \frac{A}{NAX_0\rho} \left[ \left( \frac{4}{3} + C \right) \left( \frac{1}{y} - 1 \right) + y \right], \quad (8)
\]

where \(C\) is a small material-dependent correction (Appendix A). Figure 9 shows the integrated thickness of material upstream of the COT, in terms of radiation lengths, traversed by reconstructed electron tracks in \(W \rightarrow e\nu\) data. The number of photons emitted per layer is given by:

\[
N_\gamma = \frac{x}{X_0} \left[ \left( \frac{4}{3} + C \right) (y_0 - \ln y_0 - 1) + \frac{1}{2} (1 - y_0)^2 \right], \quad (9)
\]

where \(x\) is the thickness of the layer and \(y_0\) is a lower threshold introduced to avoid infrared divergences. We use \(y_0 = 10^{-4}\) [49] and determine \(C = 0.0253\) using the silicon atomic number \(Z = 14\).

For each layer of the silicon or COT material, we use a Poisson distribution with mean \(N_\gamma\) to determine the number of photons radiated in that layer. For each radiated photon, we calculate \(y\) from the spectrum in Eq. (8). To correct for inaccuracies of the screened Bethe-Heitler equation at the ends of the \(y\) spectrum, we apply a suppression factor if \(y \leq 0.005\) or \(y \geq 0.8\).

For radiation of high-momentum photons \((y \gtrsim 0.8)\), the approximation of complete screening of the nuclear electromagnetic field by the atomic electrons breaks down. In this region, the full Bethe-Heitler equation for incomplete screening [48] must be used. We implement this correction by removing generated photons in the high-\(y\) region such that we match the reduced cross section from incomplete screening.

Two effects reduce the cross section for low-momentum photon radiation [50]: multiple scattering and Compton scattering. Multiple Coulomb scattering suppresses long-distance interactions, and the resulting \(LPM\) suppression [51] in low-momentum radiation can be expressed in terms of the Bethe-Heitler cross section [52]:

\[
S_{LPM} \equiv \frac{d\sigma_{LPM}/dy}{d\sigma_{BH}/dy} = \sqrt{\frac{E_{LPM}}{E_e} \frac{y}{(1 - y)}}, \quad (10)
\]
where $E_{LPM}$ depends on the material. We use $E_{LPM} = 72$ TeV, appropriate for silicon, and apply the suppression when $S_{LPM} < 1$.

Radiated photons scatter off the atomic electrons, and destructive interference of low-momentum photons suppresses this radiation [53]. The suppression factor is:

$$S_{Compton} = \frac{y^2}{y^2 + E_p^2/E_e^2}, \quad (11)$$

where $E_p = \gamma \omega_p$ is 2.4 MeV for a 40 GeV electron in silicon, using the silicon plasma frequency $\omega_p$, and $\gamma$ is the Lorentz factor.

In any given simulated event, the product of $S_{LPM}$ and $S_{Compton}$ provides the probability that a photon generated from the screened Bethe-Heitler equation with $y \leq 0.005$ survives the low-momentum suppression. For a 40 GeV electron radiating a 20 MeV (8 MeV) photon, the suppression factors are $S_{LPM} = 0.95 (0.60)$ and $S_{Compton} = 0.99 (0.92)$. Our simulated $y$ spectrum from $W$ boson decay electrons reproduces the spectrum obtained by a GEANT [40] simulation.

**Photon Conversion**

Photons can convert to an electron-positron pair by interacting with the tracker material. The differential cross section for a photon of energy $E_\gamma > 1$ GeV to convert into an electron with energy $E_e$ is given by the screened Bethe-Heitler equation [48]:

$$\frac{d\sigma}{dy} = \frac{A}{N_A X_\rho} \left[ 1 - \frac{(4/3 + C)}{y(1 - y)} \right], \quad (12)$$

where $y = E_e/E_\gamma$. Integrating over $y$ and multiplying by $\rho x N_A / A$ gives the total cross section, from which we obtain the following conversion probability at high photon energy:

$$P_{\gamma \rightarrow e^+e^-}(E_\gamma \rightarrow \infty) = 1 - e^{-(7/9 - C/6) x / X_\rho}, \quad (13)$$

We parametrize the cross section as a function of photon energy using the tables for photon cross sections in silicon given in [54]. We apply the ratio shown in Fig. 10 to the high-energy cross section when calculating the conversion probability.

For each radiated photon upstream of the COT, we integrate the material between the radiation point and the COT inner can. If the photon converts, we take the conversion point to be halfway between the radiation point and the inner can. If the photon does not convert before the COT, we integrate the material in the COT and take a converting photon to convert halfway through the COT.

**Compton Scattering**

The cross section for a low-momentum photon to scatter off an electron is similar to that of conversion into an $e^+e^-$ pair. The differential cross section with respect to the photon fractional energy loss $y$ can be approximated as (Appendix A):

$$\frac{d\sigma}{dy} \propto 1/y + y. \quad (14)$$

Using a lower bound of $y = 0.001$, this spectrum approximates the Compton energy loss distribution for photons radiated from electrons from $W$ boson decays.

We calculate the total cross section in terms of the pair production cross section using the tables for photon interactions in silicon in [54]. The ratio of cross sections as a function of energy is parametrized as (Fig. 11):

$$R_{Com} \equiv \frac{\sigma_{\gamma \rightarrow e^+e^-}(E_\gamma \rightarrow \infty)}{\sigma_{\gamma \rightarrow e^+e^-}(E_\gamma \rightarrow \infty)} = e^{F(E_\gamma)}, \quad (15)$$
where \( F(E_\gamma) = 2.35e^{-1.16E_\gamma} + 2.42e^{-15.8E_\gamma} - 5.21 - 0.151E_\gamma \), with \( E_\gamma \) in GeV, and \( \gamma \) and \( \gamma' \) are the initial- and final-state photons, respectively. We thus use the following Compton scattering probability per layer:

\[
P_{\gamma \rightarrow \gamma'} = R_{Com}(7/9 - C/6)x/X_0. \tag{16}
\]

**Energy Loss in Solenoid**

After exiting the tracker electrons and photons travel through the time-of-flight (TOF) system and the solenoid. These systems have thicknesses of \( \approx 10\% \) and \( \approx 85\% \) of a radiation length, respectively. With this much material it becomes prohibitive to model individual radiative processes, and we instead use a parametrized energy-loss model determined from a GEANT simulation \([40]\). The energy loss is defined as the difference in energy of a single particle entering the TOF and the total energy of particles exiting the solenoid.

Figure 12 shows the mean \( p_T \) loss as a function of \( \log_{10}(p_T/\text{GeV}) \) for electrons with \( p_T > 400 \text{ MeV} \) and photons traversing the time-of-flight system and solenoid.

![Energy Loss in Solenoid](image)

The calorimeter simulation models the response of the electromagnetic calorimeter as a function of each particle’s energy and position, and the fraction of shower energy leaking into the hadronic calorimeter.

The electromagnetic calorimeter response, or the average measured energy divided by the true particle energy entering the calorimeter, can depend on each particle’s energy. Possible sources of this dependence are variations in light yield as a function of calorimeter depth, attenuation in the light guide from the scintillator to the phototube, or leakage of showering particles into the hadronic calorimeter. The mean fractional energy leakage into the hadronic calorimeter for particles exiting the tracker, determined using the GEANT calorimeter simulation, is shown as a function of \( \log_{10}(p_T/\text{GeV}) \) in Fig. 13.

For a low-\( p_T \) particle exiting the tracker, the distribution of energy loss into the hadronic calorimeter is adequately described by an exponential. For high-\( p_T \) particles (\( \gtrsim 10 \text{ GeV} \)), the distribution has a peak at non-zero values of energy loss. In this energy region we model the hadronic energy loss fluctuations with the distributions shown in Fig. 13. Because a non-negligible fraction of electrons lose a significant amount of energy (5–10\%) in the hadronic calorimeter, it is important to model the energy loss spectrum in addition to the mean hadronic energy loss.

Calorimeter Response and Fiduciality
uniform response as a function of distance from the "non-linearity" of the calorimeter response. The parameter $E/p$ (Section VI). The inclusive response $S \rightarrow W p$ mean is a function of $10 \xi$. We determine the slope parameter $\kappa$ for incident particles by applying the inverse of the quadratic attenuation function to the measured EM energy. We match this procedure in the simulation.

The EM calorimeter response drops rapidly as a particle crosses the edge of the scintillator and into the dead region between towers [30]. We take the calorimeter to have zero response for any particle with $|\text{CES} x| > 23.1$ cm or $|\text{CES} z| < 4.2$ cm. For the $m_W$ measurement we only use high-energy electrons far from the dead regions (Section VII B). These measurements are incorporated into the simulation.

We apply the following smearing to the calorimeter cluster energy:

$$\sigma_E/E = \sqrt{0.135^2/p_T + \kappa^2},$$

where the constant term $\kappa$ is determined to be $0.89 \pm 0.06 \text{(stat)} \pm 0.13 \text{(sys)}$% from a fit to the width of the electron $E/p$ peak in $W$ boson decays [55]. We find further energy smearing is necessary to model the multi-particle energy clusters populating the high $E/p$ region. When a simulated $W$ or $Z$ decay electron radiates in the tracker, we apply an additional fractional resolution of $\kappa_\gamma = [8.3 \pm 2.2 \text{(stat)}]$% to each radiated particle. This smearing contributes $\approx 1.3$% to the effective constant term, and is determined from a fit to the width of the $Z$ boson mass peak reconstructed from radiative electrons ($E/p > 1.06$).

The final contribution to the electron cluster energy comes from the underlying event [39] and additional $p\bar{p}$ interactions. As with muons, we measure this energy distribution in W boson data as a function of $u_\parallel$, $u_\perp$, and the electron tower $\eta$ (Section VII B). These measurements are incorporated in the simulation.

IV. $W$ BOSON SELECTION

The $W$ boson samples are collected with triggers requiring at least one central ($|\eta| \leq 1$) lepton candidate in the event. A narrow kinematic region is defined for $W$ boson selection: $30$ GeV $< \text{lepton } p_T < 55$ GeV; $30$ GeV $< \rho_T < 55$ GeV; $60$ GeV...
< m_T(l, p_T) < 100 \text{ GeV}; and \, u_T < 15 \text{ GeV}. This
selection results in low background while retaining
events with precise m_W information. Additional
background rejection is achieved through event se-
lection targeting the removal of Z boson decays to
leptons. To minimize bias, lepton selection criteria
are required to have high efficiency or to be explic-
itly modeled by our fast simulation.

A. W → µν Selection

Muons are identified based on their reconstructed
COT track quality and production vertex, minimum
ionizing energy deposited in the calorimeter, and the
consistency of the track segments reconstructed in
the muon chambers with the COT tracks.

All charged lepton candidates from W and Z bos-
on decay are required to have fully-fiducial central
(|z_0| < 60 \text{ cm}) COT tracks with at least 5 hits
on each of ≥ 3 axial superlayers and ≥ 3 small-
angle stereo superlayers. For muon candidates we
remove background from decays of long-lived hadrons
to muons (“decays in flight”) by requiring the track
impact parameter to be small (|d_0| < 1 \text{ mm}) and the
track fit quality to be good (χ^2/\text{dof} < 3). After this
initial selection, the COT track parameters are up-
dated with an additional constraint to the transverse
position of the beam, which has a size of ≈ 30 \mu m
in the luminous region. The beam constraint results in
a factor of ≈ 3 improvement in momentum resolution
for muons from W boson decays.

Each muon candidate’s COT track is extrapolated
to the calorimeter and its energy deposition in the
electromagnetic and hadronic calorimeters is sepa-
ately measured. Muons near a tower edge in the
z direction cross two calorimeter towers, and those
tower energies are combined to determine the muon’s
total energy deposition. We require the muon’s elec-
tromagnetic energy deposition E_{EM} to be less than 2
GeV and its hadronic energy deposition E_{Had} to be
less than 6 GeV [27].

All W muon candidates must have a track segment
in either the CMU and CMP detectors, or the CMX
detector. COT tracks extrapolated to these detectors
must have \, r − \phi positions that match to within 3, 5,
or 6 cm of the CMU, CMP, or CMX track segment
positions, respectively.

The Z/γ^* → µµ process presents a significant
background to the W → µν sample. We reduce
this background by removing events with a second
opposite-charge muon candidate passing the above
selection, or passing the following looser set of cri-
teria: an opposite-charge track with p_T > 10 \text{ GeV},
|d_0| < 1 \text{ mm}, ≥ 2 axial superlayers with ≥ 5 hits,
and ≥ 2 (1) small-angle stereo superlayers with ≥ 5
hits for tracks fully (partially) fiducial to the COT;
E_{EM} < 2 \text{ GeV} and E_{Had} < 6 \text{ GeV}; and calorim-
eter isolation < 0.1. Calorimeter isolation is defined
as the calorimeter p_T in an \eta − \phi cone of radius 0.4
surrounding the muon calorimeter towers, divided by
the muon track p_T. For events with one identified
W decay muon and a second muon candidate pass-
ing the looser criteria, the identified W decay muon
must also have isolation < 0.1 for the event to be re-
jected from the W boson sample. The full W boson
sample, after kinematic selection and Z boson rejec-
tion, contains 51,128 events in (190.8 ± 11.1) pb^{-1}
of data.

The identification efficiency of muons has a small
dependence on the recoil in W → µν and Z → µµ
events, due primarily to the track χ^2 and d_0 require-
ments. We measure this dependence using Z → µµ
events, selected with one muon passing the W muon
candidate criteria and a second “probe” muon iden-
tified as a track with p_T > 30 \text{ GeV}. The two muons
must have opposite charge and reconstruct to an
invariant mass in the 81 − 101 \text{ GeV} range. The frac-
tion of probe muons passing the additional W muon
candidate selection criteria is shown in Fig. [13] as a
function of net recoil energy along the muon direction
(u_||). The observed dependence is parametrized as:

\begin{equation}
\epsilon = a[1 + b(u_|| + |u_||)],
\end{equation}

where a is a normalization factor that does not affect
the m_W measurement and \, b = [1.32 ± 0.40(stat) \times
10^{-3}]. We vary b by ±3\sigma in simulated data and fit
for m_W. Assuming a linear variation of m_W with b,
we derive uncertainties of \delta m_W = 1.6, and 13 \text{ MeV}
for the m_T, p_T, and p_T^* fits, respectively.

B. W → eν Selection

Electron identification uses information from the
COT track quality and production vertex, the match-
ing of the track to calorimeter energy and position,
and the longitudinal and lateral calorimeter energy
profiles.

An electron candidate’s COT track has the same
fiduciality and hit usage requirements as a muon can-
didate track, and utilizes the same beam-constrained
track fit. The track is required to have p_T > 18 \text{ GeV},
a kinematic region where the trigger track-finding ef-
ficiency has no p_T dependence.

The clustering of showers in the CES produces an
energy-weighted position at the electron shower max-
imum. We require the CES cluster to be well sepa-
Electrons are differentiated from hadrons by their high fraction of energy deposited in the electromagnetic calorimeter. The electron’s EM energy is measured in two neighboring towers in \( \eta \), while the energy collected in the hadronic calorimeter is measured in three towers. The ratio, \( E_{\text{Had}}/E_{\text{EM}} \), is required to be less than 0.1. Only the EM calorimeter measurement is used to determine the electron’s \( p_T \).

An electron shower will typically be confined to a single tower, with a small amount of energy flowing into the nearest tower in \( \eta \). We define an error-weighted difference between the observed and expected energies in the two towers neighboring the electron in the \( \eta \) direction [56]:

\[
L_{\text{shr}} = 0.14 \sum_i \frac{E_i^{\text{adj}} - E_i^{\text{exp}}}{\sqrt{0.14^2 E_i^{\text{adj}} + (\Delta E_i^{\text{exp}})^2}},
\]

where \( E_i^{\text{adj}} \) is the energy in a neighboring tower, \( E_i^{\text{exp}} \) is the expected energy contribution to that tower, \( \Delta E_i^{\text{exp}} \) is the RMS of the expected energy, energies are measured in GeV, and the sum is over the two neighboring towers. We require \( L_{\text{shr}} < 0.3 \), consistent with the trigger criterion (Section III A 4).

The \( Z \rightarrow ee \) background is highly suppressed by the \( u_T < 15 \) GeV requirement for the \( W \) boson sample. Residual background results from electrons passing through dead calorimeter regions, which reduces \( u_T \) and increases \( \pT \). We remove events from the \( W \) sample if a track with \( p_T > 20 \) GeV and \( |d_0| < 0.3 \) cm extrapolates to a calorimeter region with reduced response (\(|\text{CES } x| > 22 \) cm or \(|\text{CES } z| < 6 \) cm), and the track’s calorimeter isolation is < 0.1 (Section IV A). The full \( W \rightarrow e\nu \) selection results in a sample of 63,964 candidate events in \((218.1 \pm 12.6) \) pb\(^{-1}\) of integrated luminosity.

The track selection in the single-electron trigger (Section III A 4) results in an \( \eta \)-dependent trigger efficiency for reconstructed electrons (Fig. 15). We study this efficiency using \( W \) events selected with a trigger where the track requirements are replaced by a \( \pT \) threshold. The efficiency decreases as \( |\eta| \) decreases because the reduced path length reduces the ionization charge collected by each wire, thus reducing the single hit efficiency. There is an additional decrease in efficiency due to the dead region at \( |z| \lesssim 2 \) mm. Electrons crossing this region at track \( |\eta| = 0 \) are not included in the efficiency plot, since we only measure electrons with \(|\text{CES } z| > 9 \) cm. Thus, at \( |\eta| = 0 \) there is no inefficiency due to the dead COT region, and the measured efficiency increases.

We measure the \( u_|| \) dependence of the electron identification efficiency (Fig. 16) using \( Z \rightarrow ee \) events, selected with one electron passing the \( W \) electron candidate criteria and a second “probe” electron identified as an EM energy cluster with \( E/p < 2 \). Since the probe electron definition includes
an $E/p$ requirement, this cut is not included in the efficiency measurement. We instead study the unbiased $E/p < 2$ efficiency by recalculating $E$ and $u_{||}$ for towers separated in \phi from the identified electron in $W \rightarrow e\nu$ events, and find no significant $u_{||}$ dependence in this efficiency. In the simulation we use $b = 0 \pm 0.54 \times 10^{-3}$, obtained by fitting the measured efficiencies to the function in Eqn. [19]. We vary $b$ by $\pm 3\sigma$ in pseudoexperiments and assume linear variation of $m_W$ with $b$ to derive uncertainties of $\delta m_W = 3.5$, and $16$ MeV for the $m_T$, $p_T$, and $p_T^2$ fits, respectively. Since $b$ is measured with different data samples for the electron and muon channels, there is no correlation between the corresponding systematic uncertainties.

V. TRACK MOMENTUM MEASUREMENT

Muon momenta are determined from helical fits to tracks reconstructed using COT information. The momentum resolution of prompt muons is improved by constraining the helix to originate from the transverse beam position. A given muon’s transverse momentum is determined by the Lorentz equation,

$$mv^2/R = evB, \quad p_T = eB/(2|c|),$$

where $B$ is the magnetic field, $R$ is the radius of curvature, $c \equiv q/(2R)$ is the curvature of the helix, and $q$ is the muon charge. The \textit{a priori} momentum scale is determined by the measurements of the magnetic field and the radius of the tracker. At CDF, $eB/2 = 2.11593 \times 10^{-3}$ GeV/cm, where $B$ is measured using an NMR probe at a COT endplate. Measurements of the local field nonuniformities and tracker geometry were performed during construction and installation and are used to determine the positions of individual track hits. We find these measurements provide an \textit{a priori} momentum scale accuracy of $\approx 0.15\%$.

We refine the momentum scale calibration with data. Using reconstructed cosmic ray muon tracks, we align the relative positions of the tracker wires. Track-level corrections derived from $W \rightarrow e\nu$ data reduce relative curvature bias between positive and negative particles. Finally, we perform an absolute calibration of the momentum scale using high-statistics data samples of $J/\psi$, $\Upsilon$, and $Z$ boson decays to muons. The final calibration is applied as a relative momentum correction $\Delta p/p$ to the $W$ boson data and has an accuracy of $\approx 0.02\%$.

A. COT Alignment

The COT contains 30,240 sense wires for measuring the positions of charged particles passing through the detector. The position measurements rely on an accurate knowledge of the wire positions throughout the chamber. We determine these positions using a combination of alignment survey, computer modeling, and cosmic-ray muon data. Any remaining biases in track parameter measurements are studied with $J/\psi \rightarrow \mu\mu$ and $W \rightarrow e\nu$ data, from which final track-level corrections are derived.

After construction of the COT endplates, the position of each 12-wire cell was measured with an accuracy of $\pm 13$ \textmu m using a coordinate measuring machine. The effect of the load of the wire plane and field sheets was modeled with a finite element analysis (FEA) and found to cause an endplate bend towards $z = 0$ cm, with the maximum bend of $\approx 6$ mm in the fifth superlayer [29]. An equivalent load was applied to the detector and further measurements found the FEA to be accurate to within $\approx 20\%$. The FEA results were scaled to match the measurements, and the positions determined from the FEA were set as the directly-determined cell positions.

While each cell position determines the average positions of its 12 sense wires within the chamber, several effects create a non-linear wire shape as a function of $z$. Gravity has the most significant effect, causing each wire to sag $\approx 260$ \textmu m in $y$ at $z = 0$.
cm. Electrostatic deflection towards the nearest field sheet occurs for cells where the sense wire is not centered between the field sheets. By construction, the wires are slightly offset within a cell; in addition, the gravitational sag of the field sheets is larger than that of the sense wires, resulting in an electrostatic deflection that partially counteracts the sag of the sense wires. Combined, the electrostatic effects cause a \( \phi \)-dependent wire shift that has a maximum of 74 \( \mu \)m at \( \phi = 145^\circ \) and \( z = 0 \) cm. The gravitational and electrostatic effects were combined to determine the best \textit{a priori} estimate of the wire shapes.

Starting from the predicted cell and wire positions, we develop \textit{in situ} corrections based on cosmic-ray muon data taken during \( p\bar{p} \) crossings with the single muon trigger. The data are selected by requiring exactly two reconstructed tracks in the event, eliminating effects from overlapping hits from collision-induced particles. Since the two tracks on opposite sides of the COT result from a single cosmic-ray muon, we refit both tracks to a single helix and determine hit residuals with respect to this helix [57]. For each cell, we use the residuals to determine a tilt correction about its center, and a shift correction along the global azimuth (Fig. 17). We show the tilt and shift corrections for the inner superlayer of the west endplate in Fig. 18 after removing global corrections. We apply these corrections to each cell of each superlayer in each endplate. In addition, we measure a relative east-west shift and include it in each cell’s correction.

We combine the cell-based corrections with wire-based corrections for the shapes of the wires between the endplates. We measure these corrections as functions of \( z \) and radius \( R \) using the differences in the measured \( d_0 \) and curvature parameters for the helix fits on opposite sides of the COT for a cosmic ray muon. The corrections are applied as additional offsets \( \Delta \xi \) of the wires at \( z = 0 \) cm, with a parabolic wire shape as a function of \( z \). The corrections include a radial dependence,

\[
\Delta \xi = -160 + 380(R/140) - 380(R/140)^2, \tag{22}
\]

where \( R \) is measured in cm and \( \Delta \xi \) in \( \mu \)m. Figure 19 shows the gravitational and electrostatic shifts of a wire as a function of \( z \) at \( \phi = \pi \), as well as the data-based correction at \( R = 130 \) cm (the outer superlayer).

The cell- and wire-based corrections are implemented for the track-finding and fitting stage, and re-
reduce the measured hit resolution for high-momentum muons from \( \approx 180 \, \mu m \) to \( \approx 140 \, \mu m \). Final track-based corrections are applied to the measured track curvature, which is inversely related to the transverse momentum [Eq. (21)]. Expanding the measured curvature \( c \) as a function of the true curvature \( c_t \) in a Taylor series around zero,

\[
c = \epsilon_1 + (1 + \epsilon_2) c_t + \epsilon_3 c_t^2 + \epsilon_4 c_t^3 + \ldots,
\]

the terms even in \( c_t \) cause biases in positive tracks relative to negative tracks, which tend to cancel when the two are averaged. The term linear in \( c_t \) scales the true curvature and is determined by the momentum calibration. The \( \epsilon_4 c_t^3 \) term is the first to directly affect mass measurements and is suppressed by the \( c_t^3 \) factor at low curvature (high momentum).

Corrections for high-momentum tracks from \( W \) and \( Z \) decay particles are determined using the difference in \( E/p \) for \( e^+ \) and \( e^- \) from \( W \) decays, which should be zero in the absence of misalignments. This difference can be used to constrain \( \epsilon_1 \), the first term in the Taylor expansion. Figure 20 shows the differences in \( E/p \) as functions of \( \cot \theta \) and \( \phi \), before and after corrections of the following form:

\[
\delta c = a_0 + a_1 \cot \theta + a_2 \cot^2 \theta + b_1 \sin(\phi + 0.1) + b_3 \sin(3\phi + 0.5).
\]

The terms can be interpreted as arising from the following physical effects: a relative rotation of the outer edge to the inner edge of each endplate \( (a_0) \); a relative rotation of the east and west endplates \( (a_1 \cot \theta) \); and a mismeasurement of the beam position \( (b_1 \sin(\phi + 0.1)) \). The measured values of the parameters \( a_0, a_1, a_2, b_1, \) and \( b_3 \), are shown in Table III.

Varying \( a_1 \) by \( \pm 3\sigma \) in pseudexperiments and assuming linear variation of the momentum scale with \( a_1 \), we find the \( a_1 \) uncertainty results in a relative momentum scale uncertainty of \( \pm 0.07 \times 10^{-3} \) for \( W \) and \( Z \) boson mass measurements. The other parameter uncertainties, as well as residual higher-order terms, have a negligible impact on the momentum scale for the \( m_W \) measurement.

**B. \( J/\psi \to \mu \mu \) Calibration**

With a measured \( \sigma \times BR \) of \( 16.3^{+1.4}_{-1.3} \) nb [28], \( J/\psi \) mesons are the Tevatron’s most prolific source of resonant decays to muon pairs. In addition to

**FIG. 19**: The net wire shift in \( y \) as a function of \( z \) from gravitational sag only (solid), including electrostatic effects (dashed), and including data-based corrections from Eq. (22) (dotted). The shift is shown at \( \phi = \pi \) and \( R = 130 \, cm \).

**FIG. 20**: The difference between \( e^+ \) and \( e^- E/p \) as a function of \( \cot \theta \) (top) and \( \phi \) (bottom) before (triangles) and after (diamonds) track-level corrections.
TABLE III: The parameters used to correct the track curvature of electrons and muons from W and Z boson decays. The values and statistical uncertainties are determined from fits to the E/p difference between positrons and electrons.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>Value (×10⁻⁷ cm⁻¹)</th>
</tr>
</thead>
<tbody>
<tr>
<td>a₀</td>
<td>−0.66 ± 0.17</td>
</tr>
<tr>
<td>a₁</td>
<td>−1.6 ± 0.3</td>
</tr>
<tr>
<td>a₂</td>
<td>−2.1 ± 0.5</td>
</tr>
<tr>
<td>b₁</td>
<td>−2.1 ± 0.2</td>
</tr>
<tr>
<td>b₃</td>
<td>5.7 ± 1.7</td>
</tr>
</tbody>
</table>

its high statistics, the J/ψ’s precisely known mass \(m_{J/\psi} = 3096.88 ± 0.04\) MeV \[58\] and narrow width \(\Gamma_{J/\psi} = 0.0934 ± 0.0021\) MeV \[11\]) make it a key component of the track momentum calibration. We perform measurements of the J/ψ mass as a function of mean inverse muon \(p_T\) to determine a momentum scale correction and extrapolate to the high-\(p_T\) region relevant for W and Z boson decays.

1. Data Sample

The J/ψ data sample is collected with a Level 1 trigger requiring one \(p_T > 1.5\) GeV XFT track with a matching CMU track segment, and a second \(p_T > 1.5\) (2) GeV XFT track with a matching CMU (CMX) segment. At Level 3, the two corresponding COT tracks must have opposite charge and consistent z vertex positions \(|\Delta z_0| < 5\) cm, and must form an invariant mass between 2.7 and 4 GeV. The resolution on the invariant mass measurement degrades at high track momentum, so to avoid trigger bias the mass range is extended to 2 GeV < \(m_{\mu\mu}\) < 5 GeV when the \(p_T\) of the muon pair \(p_T^{\mu\mu}\) is greater than 9 GeV.

Candidate events are selected offline by requiring two COT tracks, each with \(p_T > 2\) GeV, \(|d_0| < 0.3\) cm, and \(\geq 7\) hits on each of the eight superlayers. The tracks must originate from a common vertex \(|\Delta z_0| < 3\) cm and form an invariant mass in the range \(2.95, 3.21\) GeV.

A significant fraction (≈ 20%) of the J/ψ mesons in our data sample result from decays of B hadrons, which have an average proper decay length of \(\approx 0.5\) mm. The muons from the J/ψ decay can thus originate outside the beam radius. Therefore, no beam constraint is applied in the COT track fit of muon candidates from J/ψ decays.

The total sample consists of 606,701 J/ψ candidates in \((194.1 ± 11.3)\) pb⁻¹ of integrated luminosity.

2. Monte Carlo Generation

We use PYTHIA \[59\] to generate \(J/\psi \rightarrow \mu\mu\) events, from which templates are constructed to fit the data for the momentum scale. The shape of the \(m_{\mu\mu}\) distribution from \(J/\psi\) decays is dominated by the \(p_T\)-dependent detector resolution. We therefore model the \(p_T^{J/\psi}\) distribution as well as the \(p_T\) and relative \(p_T\) of the muons in a \(J/\psi\) decay. To obtain an adequate model, we empirically tune the generated \(J/\psi\) kinematics to describe the relevant data distributions for the \(J/\psi\) mass fits.

To tune the \(p_T^{J/\psi}\) distribution, we boost the \(J/\psi\) momentum by changing its rapidity \((y_{J/\psi})\) along its direction of motion \(\hat{p}_{J/\psi}\). In 50% of the generated events we multiply \(y_{J/\psi}\) by 1.215, and in the other 50% we multiply it by 1.535. The decay angle \(\theta^*\) in the \(J/\psi\) rest frame relative to \(\hat{p}_{J/\psi}\) is tuned by multiplying cot \(\theta^*\) by 1.3. After tuning, the simulation matches the relevant background-corrected data distributions, as shown in Fig. 21.

The PYTHIA event generator does not include energy loss due to final-state photon radiation from the muons in \(J/\psi\) decays. To simulate this effect, we scale each muon’s momentum by a factor \(x\) determined from the following leading-log probability distribution for soft photon radiation \[59, 60\]:

\[
f(x) = \beta(1 - x)^{\beta - 1},
\]

with

\[
\beta = \frac{\alpha_{EM}}{\pi} \ln(Q^2/m_{\mu\mu}^2) - 1
\]

and \(Q^2 = m_{J/\psi}^2\).

3. Momentum Scale Measurement

The momentum scale is calibrated using \(J/\psi\) decays by fitting the dimuon mass as a function of mean inverse \(p_T\) of the two muons, and then extrapolating to high \(p_T\) \((p_T^{-1}) \approx 0\) GeV⁻¹. This procedure results in a track momentum calibration accuracy of 0.025%.

The momentum scale calibration requires an accurate modeling of the muon ionization energy loss in the tracker. Each muon passing through the silicon and COT detectors loses on average 9 MeV at normal incidence. The combined effect on the reconstructed \(m_{\mu\mu}\) is about 0.6% of \(m_{J/\psi}\), a factor of ≈ 20 larger than our total uncertainty. Since the ionization energy loss \(E_I\) varies only logarithmically with \(p_T\) (Section [III.B.1]), the relative effect on the reconstructed
Thus, in a linear fit of $\Delta m/m$ as a function of mean inverse $p_T$, a non-zero slope approximately corresponds to $E_I$. Since we model the ionization energy loss based on the known detector material, this slope should be zero. We however find that we need to scale down the ionization energy loss from the detector parametrization (Section III.B.1) by 6% to achieve a zero slope. We show the result of this tuning in Fig. 22 replacing $\Delta m/m$ on the $y$-axis with the relative momentum correction $\Delta p/p$ to be applied to the data in order to measure $m_{J/\psi} = 3096.88$ MeV. The tuning is based on a $\langle p_T^{-1} \rangle$ region of (0.1, 0.5) GeV$^{-1}$, divided into eight bins. We find a scale correction of $\Delta p/p = [-1.64 \pm 0.06 \text{ (stat)}] \times 10^{-3}$ from a linear fit to $\Delta p/p$ as a function of $\langle p_T^{-1} \rangle$.

Each $\Delta p/p$ value in Fig. 22 is extracted via a binned likelihood fit to the $m_{\mu\mu}$ distribution for each $\langle p_T^{-1} \rangle$ bin. Since the mass resolution varies significantly with $\langle p_T^{-1} \rangle$, the fit ranges are adjusted from 3.08 ± 0.13 GeV for $\langle p_T^{-1} \rangle = (0.1, 0.15)$ GeV$^{-1}$ to 3.08 ± 0.08 GeV for $\langle p_T^{-1} \rangle = (0.45, 0.5)$ GeV$^{-1}$. The background is modeled as a linear function of $m_{\mu\mu}$, with normalization and slope determined from upper and lower sideband regions whose combined width is equal to that of the mass fit window. The results of the fits in the $\langle p_T^{-1} \rangle = (0.15, 0.2)$ GeV$^{-1}$ and $\langle p_T^{-1} \rangle = (0.25, 0.3)$ GeV$^{-1}$ ranges are shown in Fig. 23.

The $J/\psi$ momentum calibration includes corrections to the curvature $c$ derived from the measured dimuon mass as a function of $\Delta \cot \theta$ between the positive and negative muons from the $J/\psi$ decay. Biases linear in $\Delta \cot \theta$ are removed with a curvature correction linear in $\cot \theta$:

$$\delta c = [(-7 \pm 1) \times 10^{-7} \text{ cm}^{-1}] \cot \theta,$$

where the uncertainty is statistical only. Biases quadratic in $\Delta \cot \theta$ are removed with the following correction to the absolute length scale of the COT along the $z$ axis (statistical uncertainty only):

$$\delta \cot \theta = [(-3.75 \pm 1.00) \times 10^{-4}] \cot \theta.$$
4. Momentum Scale Uncertainties

Systematic uncertainties on the momentum scale correction extracted from $J/\psi \rightarrow \mu\mu$ decays (Table IV) are dominated by the incompleteness of the QED and energy loss models. At low muon $p_T$ (high $\langle p_T^{-1} \rangle$), the mass fits become increasingly sensitive to QED and energy loss modeling because of the better mass resolution and higher statistics. Since we only model the mean ionization energy loss, our modeling of the mass region below the peak is imperfect. Additionally, our neglect of higher-order QED corrections affects the modeling of this region. We study possible bias from our incomplete model by changing the $Q^2$ value in the photon radiation probability function [Eq. (28)] such that the $\chi^2$ of the inclusive $m_{\mu\mu}$ fit is minimized. We find that this change affects $\Delta p/p$ by $0.2 \times 10^{-3}$.

If there is a relative tilt between the solenoid and the tracker axes, the extracted momentum scale correction will have a linear dependence on $\Delta \cot \theta$. In addition, incomplete corrections of the magnetic field nonuniformities near the ends of the solenoid can cause a quadratic $\cot \theta$ dependence. We study the $\cot \theta$ dependence of $\Delta p/p$ using $J/\psi$ decays where both muons are measured in the same $\cot \theta$ region ($|\Delta \cot \theta (\mu\mu)| < 0.1$). We find that if we correct for the observed quadratic dependence, the extracted $\Delta p/p$ changes by $0.1 \times 10^{-3}$.

The uncertainty on the material correction propagates to a momentum scale uncertainty of $0.06 \times 10^{-3}$ when extrapolated to high momentum, as shown in Fig. 22. An additional statistical uncertainty of $0.01 \times 10^{-3}$ on the scale is determined by fixing the material correction and fitting for the scale.

The statistical uncertainties on the $J/\psi$ alignment corrections [Eq. 28 and 29] have a $0.05 \times 10^{-3}$ effect on $\Delta p/p$. We test our model of the $m_{\mu\mu}$ lineshape by changing the fit range by $\pm 20\%$, and find a $\pm 0.05 \times 10^{-3}$ change in $\Delta p/p$.

We apply the same $p_T$ thresholds offline as in the trigger for muons with CMU segments. Since we do not model a $p_T$-dependent trigger efficiency, any inefficiency could cause a bias in the reconstructed $m_{\mu\mu}$. We investigate this possibility by varying the offline $p_T$ thresholds by $\pm 5\%$, and find a $\Delta p/p$ variation of $\pm 0.04 \times 10^{-3}$.

The quality of the fit is highly sensitive to the hit resolution model, but the momentum scale correction is not. Changing the simulated COT hit resolution by $\pm 10 \mu m$, which corresponds to a $> 10\sigma$ statistical variation, results in a $\pm 0.03 \times 10^{-3}$ change in $\Delta p/p$. We include this in our systematic uncertainty estimate.
A $\pm 0.03 \times 10^{-3}$ uncertainty on $\Delta p/p$ from the background model is determined by changing its linear dependence on $m_{\mu\mu}$ to a constant. Finally, the world-average $J/\psi$ mass value used in this measurement contributes $\pm 0.01 \times 10^{-3}$ to the uncertainty on $\Delta p/p$.

The final momentum scale correction derived from $J/\psi$ data is:

$$\Delta p/p = (-1.64 \pm 0.25) \times 10^{-3}. \quad (30)$$

C. $\Upsilon \rightarrow \mu\mu$ Calibration

The $b\bar{b}$ resonance $\Upsilon$ provides a complementary momentum scale calibration tool to the $J/\psi$. Its precisely measured mass $m_\Upsilon = (9460.30 \pm 0.26)$ MeV [11] is three times larger than that of the $J/\psi$, so an $\Upsilon$ momentum scale calibration is less sensitive to the material and energy loss model than that of the $J/\psi$. Because the $b\bar{b}$ resonances are the highest mass mesons, long-lived hadrons do not decay to the $\Upsilon$ and the muons from $\Upsilon$ decay effectively originate from the collision point. We improve the accuracy of the muon measurements by constraining their tracks to the beam position, which is the same procedure applied to the $W$ and $Z$ decay lepton tracks.

The $\Upsilon$ data sample is based on the same Level 1 trigger as the $J/\psi$ sample (Section IV.B.1). The Level 3 requirements are: one reconstructed track with $p_T > 4$ GeV and matching CMU and CMP track segments (CMUP); a second track with opposite charge to the first, $p_T > 3$ GeV, and a matching CMU or CMX track segment; and a reconstructed mass of the two tracks between 8 and 12 GeV. Offline, the $p_T$ thresholds are increased to 4.2 (3.2) GeV for the track with a CMUP (CMU or CMX) track segment, and each track must have $|d_0| < 0.3$ cm and at least 5 hits in at least 3 axial and 3 stereo superlayers. The two tracks are required to have a common vertex ($|\Delta z_0| < 3$ cm).

We model $\Upsilon$ production and decay using PYTHIA 59, to which we apply the same tuning procedure as for $J/\psi$ generation. The data $p_T^\Upsilon$ distribution is matched in simulation by boosting the rapidity of each decay muon by $0.07y_T$ along $p_T$, where $y_T$ is the $\Upsilon$ rapidity. Radiation of photons from the final state muons is simulated using the probability distribution of Eqsns. 25 and 26. The $p_T^{\mu\mu}$ distribution is shown in Fig. 24 after subtracting background from the data.

We test any possible beam-constraint bias by separately reconstructing charged muon tracks from $\Upsilon$ decays with and without incorporating the beam constraint. For the sample with beam-constrained tracks we fit for $m_\Upsilon$ in the region 9.28 GeV < $m_{\mu\mu} < 9.58$ GeV, while for the sample with non-beam-constrained tracks we fit the region 9.25 GeV < $m_{\mu\mu} < 9.61$ GeV. In $(190.8 \pm 11.1)$ pb$^{-1}$ of integrated luminosity, we have 34,618 $\Upsilon$ candidates with beam-constrained tracks and 35,622 candidates with non-beam-constrained tracks. The two momentum scale measurements are shown in Fig. 25 and are consistent at the 2$\sigma$ level when correlations are taken into account. We define the $\Upsilon$ result to be the mean of the two values, and take half their difference ($\Delta p/p = 0.06 \times 10^{-3}$) as a systematic uncertainty on the measurement.

The remaining systematic uncertainties on the momentum scale measurement with $\Upsilon$ decays are common to those of the measurement with $J/\psi$ decays. We use the same procedures as with the $J/\psi$ calibration to estimate the sizes of the uncertainties, with one exception. Since the $\Upsilon$ sample has < 10% of the statistics of the $J/\psi$ sample, the QED and energy loss
model cannot be tested with the $\chi^2$ of the $\Upsilon \to \mu\mu$ mass fit. Instead, we change $Q$ in the photon radiation probability by the amount estimated for the $J/\psi$ systematic uncertainty (Section V B 4). We find that this variation affects $\Delta p/p$ by $\pm 0.13 \times 10^{-3}$ in the $\Upsilon$ calibration.

The final result of the $\Upsilon$ calibration is:

$$\Delta p/p = (-1.44 \pm 0.21) \times 10^{-3}. \quad (31)$$

We have verified that this result has no time dependence, at the level of the statistical precision of $\pm 0.13 \times 10^{-3}$. When combined with the momentum scale correction from the $J/\psi$ calibration, we obtain:

$$\Delta p/p = (-1.50 \pm 0.19) \times 10^{-3}. \quad (32)$$

### D. $Z \to \mu\mu$ Calibration

Given the precise momentum scale calibration from the $J/\psi$ and $\Upsilon$ decays, we measure the $Z$ boson mass and compare it to the world-average value $m_Z = (91187.6 \pm 2.1) \text{ MeV}$. We then use the world-average $m_Z$ to derive an additional $\Delta p/p$ calibration and combine it with that of the $J/\psi$ and $\Upsilon$ decays.

The systematic uncertainties of the $m_Z$ measurement are correlated with those of the $m_W$ measurement, so a momentum scale calibration with $Z$ bosons can reduce systematic uncertainties on the $m_W$ measurement. However, the statistical uncertainty from the $Z \to \mu\mu$ sample is significantly larger than the calibration uncertainty from $J/\psi$ and $\Upsilon$ decays. Thus, the main purpose of the $m_Z$ measurement is to confirm the momentum scale calibration and test our systematic uncertainty estimates.

The $Z$ boson data sample is selected using the same single-muon trigger and offline muon selection as for the $W$ boson sample (Sections II A 4 and IV A), with the exception that we remove the requirement of a track segment in a muon detector for one of the muons from the $Z$ boson decay. Removing this requirement significantly increases detector acceptance while negligibly affecting background. $Z$ boson candidates are defined by $66 \text{ GeV} < m_\mu\mu < 116 \text{ GeV}$, $p_T^{\mu\mu} < 30 \text{ GeV}$, $|\Delta t_0(\mu, \mu)| < 3$ ns, and oppositely charged muons. A muon track’s $t_0$ is defined as the time between the $pp$ bunch crossing and the muon’s production, and should be $(0 \pm 1)$ ns for $Z \to \mu\mu$ production and decay. The track $t_0$ is measured using the time information from the track hits in the COT by incorporating $t_0$ into the helical fit. The $|\Delta t_0| < 3$ ns requirement effectively removes cosmic ray muons passing through the detector. An additional cosmic ray identification algorithm reduces this background to a negligible size. After applying all selection criteria, the $Z \to \mu\mu$ sample contains 4,960 events in $(190.8 \pm 11.1) \text{ pb}^{-1}$ of integrated luminosity.

We model $Z$ boson production and decay using the RESBOS event generator and a next-to-leading order QED calculation of photon radiation from the final-state muons (Section IX). For $m_\mu\mu$ near the $Z$ boson resonance, the photon propagator and $Z/\gamma^*$ interference make small contributions to the shape of the $m_\mu\mu$ distribution. We separately simulate these components and include them as fixed “background” to the $Z$ lineshape. We measure $m_Z$ using a binned likelihood template fit to the data in the range $83 \text{ GeV} < m_\mu\mu < 99 \text{ GeV}$. Our measurement of $m_Z = [91.184 \pm 0.043(\text{stat})] \text{ GeV}$ is

<table>
<thead>
<tr>
<th>Source</th>
<th>$J/\psi \times 10^{-3}$</th>
<th>$\Upsilon \times 10^{-3}$</th>
<th>Common $\times 10^{-3}$</th>
</tr>
</thead>
<tbody>
<tr>
<td>QED and energy loss model</td>
<td>0.20</td>
<td>0.13</td>
<td>0.13</td>
</tr>
<tr>
<td>Magnetic field nonuniformities</td>
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</tr>
<tr>
<td>Beam constraint bias</td>
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<td>0</td>
</tr>
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<td>Ionizing material scale</td>
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<td>0.03</td>
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<td>0.03</td>
</tr>
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<td>Fit range</td>
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<td>0.02</td>
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<tr>
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<td>0</td>
</tr>
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<td>0</td>
</tr>
<tr>
<td>Total</td>
<td>0.25</td>
<td>0.21</td>
<td>0.17</td>
</tr>
</tbody>
</table>

**TABLE IV:** Uncertainties on the momentum scale correction derived from the $J/\psi$ and $\Upsilon$ mass measurements.
in good agreement with the world-average value of $m_Z = (91.188 \pm 0.002)$ GeV [11].

Systematic uncertainties on $m_Z$ are due to the momentum scale calibration (17 MeV), alignment correction uncertainties (7 MeV), and incomplete modeling of higher-order QED corrections (14 MeV). The combined statistical and systematic uncertainty is 49 MeV.

Given the precise world-average measurement of $m_Z$, we use the $Z$ boson resonance as an additional calibration input to $\Delta p/p$. We find that adding the $m_Z$ information reduces $\Delta p/p$ and its uncertainty by less than $0.01 \times 10^{-3}$ each.

Incorporating the alignment uncertainty (Section VA) into $\Delta p/p$ from Eqn. 32 gives the momentum scale correction applicable to the $W$ boson sample:

$$\Delta p/p = (-1.50 \pm 0.21) \times 10^{-3}. \quad (33)$$

The corresponding uncertainty on the $m_W$ fits in the muon channel is 17 MeV.

VI. ELECTRON ENERGY MEASUREMENT

An electron’s energy is measured from its shower in the electromagnetic calorimeter. We perform an initial data calibration by scaling the measured energy such that a Gaussian fit to the reconstructed dielectron mass in a region dominated by $\Upsilon \rightarrow \mu\mu$ decays (86-98 GeV) gives a mean of 91 GeV. This is slightly below the world-average $m_Z$ because the Gaussian fit is biased by the energy lost to final-state photon radiation (Section IX D). This initial data calibration is accurate to $\approx 0.15\%$.

To model the data, the simulated calorimeter energy is scaled to match the measured $E/p$ distribution of electrons in $W \rightarrow e\nu$ events. A calibrated data measurement would result in an $E/p$ of unity for electrons that do not radiate before entering the calorimeter, and deposit all of their energy in the EM calorimeter. We verify that the $E/p$ calibration is unbiased by using it to measure $m_Z$ in dielectron events. Given consistency of the measured $m_Z$ with the world-average value, we incorporate the $m_Z$ fit into the calibration. The final calibration has an accuracy of 0.037%.

FIG. 25: The momentum scale correction $\Delta p/p$ derived from binned likelihood fits to the $m_{\mu\mu}$ data distribution (points) in the region dominated by $\Upsilon \rightarrow \mu\mu$ decays. The small difference between fits using tracks with (top) and without (bottom) a beam constraint is incorporated into the systematic uncertainty. The arrows indicate the fit region and the uncertainties are statistical only.

FIG. 26: The $m_Z$ fit to data (points) in the 83 GeV $< m_{\mu\mu} < 99$ GeV region (denoted by arrows). The uncertainty is statistical only.
**A. $E/p$ Calibration**

We transfer the precise tracker calibration to the calorimeter using the ratio of electron calorimeter energy to track momentum, $E/p$. The material from the beam pipe to the inner COT wall causes bremsstrahlung that affects the measured position of the $E/p$ peak, and this material is scaled such that the simulation matches the data in the high $E/p$ region. The non-linearity of the energy scale is removed by applying a correction to the simulation scale as a function of the calorimeter shower shape properties. Offline, candidate events are required to have $\text{Had}/\text{EM} > 0.05$ and $E + p > 22$ GeV to remove any trigger bias. Using the mean of the $E/p$ range $0.8 - 1.25$, we apply relative corrections of $\mathcal{O}(3\%)$ to remove variations as functions of tower and time.

Because of bremsstrahlung radiation in the tracker, the mean $E/p$ correction has a small bias that depends on the electron path length. Since the path length increases as $|\eta|$ increases, we perform a final $|\eta|$-dependent calibration of the data. Using template fits to the $E/p$ peak region of the $W \rightarrow e\nu$ sample in bins of $|\eta|$, we derive a relative correction for each bin. This calibration removes $\approx 2\%$ residual variation in the calorimeter energy response.

![Figure 27: The maximum likelihood fit to the tracker radiation length correction factor using the high $E/p$ region ($1.19 \leq E/p < 1.85$) in $W \rightarrow e\nu$ data (points).](image-url)
FIG. 28: The energy scale as a function of measured electron calorimeter $p_T$ for $W$ (top) and $Z$ (bottom) boson decays. The fast simulation incorporates a per-particle non-linear response correction [Eq. (17)]. The combination of the linear fits to $\xi$ results in no energy dependence of the energy scale.

With the complete set of corrections applied to the data and simulation, we calibrate the simulation energy scale using $W \to e\nu$ events. The fit for $S_E$ [Eq. (17)] to the $E/p$ peak region (Fig. 29) has a statistical uncertainty of 0.025%. Including systematic uncertainties due to $S_{mat}$ ($\pm 0.011\%$) and the tracker momentum scale ($\pm 0.021\%$), we obtain a total uncertainty of 0.034% on the $E/p$ calibration of the electron energy scale.

The $E/p$ calibration requires an accurate simulation of electron radiation in the tracker. We test the track simulation by measuring $m_Z$ (Section IV B) using electron track information only. The measurement is a binned likelihood fit to the region 75 GeV < $m_{ee}$ < 99 GeV (Fig. 30), with $m_Z$ as the fit parameter. Because of the significant radiated energy loss, the test is less precise than the measurement using the calorimeter (Fig. 31). Nevertheless, we obtain good consistency with the world-average $m_Z$, verifying that we do not have any significant mismodeling of electron radiation in the tracker.

FIG. 29: The energy scale calibration using the peak $E/p$ region (0.93 ≤ $E/p$ < 1.11, denoted by arrows) in $W \to e\nu$ data (points).

FIG. 30: The $m_Z$ fit to data (points) in the 75 GeV < $m_{ee}$ < 99 GeV region (denoted by arrows), using reconstructed track information only. The uncertainty is statistical only.

FIG. 30: The $m_Z$ fit to data (points) in the 75 GeV < $m_{ee}$ < 99 GeV region (denoted by arrows), using reconstructed track information only. The uncertainty is statistical only.

B. $Z \to ee$ Calibration

Using the $E/p$-based calorimeter energy calibration, we measure the $Z$ boson mass from its decay
and (Section IX). We include the virtual photon exchange operator and a next-to-leading order QED calculation (Section V D). We use the measured

\[\langle m_{ee}\rangle \approx 67\text{ GeV}\]

fraction to be less than that of the Z boson sample in the fit region, we estimate any corresponding bias on the measured \(m_{Z}\) to be \(\lesssim 10\text{ MeV}\).

The model for Z boson production and decay to electrons is the same as for the muon decay channel (Section V D). We use the RESBOS [61] event generator and a next-to-leading order QED calculation of photon radiation from the final-state electrons [62] (Section IX). We include the virtual photon exchange and \(Z/\gamma^*\) interference contributions as fixed “backgrounds” to the Z boson lineshape, and determine \(m_{Z}\) from a binned likelihood fit to the data in the range 81 GeV < \(m_{ee}\) < 101 GeV (Fig. 31).

Systematic uncertainties on the \(m_{Z}\) measurement result from the \(E/p\) calibration (29 MeV), calorimeter non-linearity measurement (23 MeV), and higher-order QED radiation (14 MeV). The measured \(m_{Z} = (91.190 \pm 0.067\text{(stat)})\text{ GeV}\) is consistent with the world-average value \(m_{Z} = (91.188 \pm 0.002)\text{ GeV} [11]\), given the total uncertainty of 78 MeV on the measurement.

The uncorrelated uncertainties in the combination of the \(m_{Z}\) and \(E/p\) calibrations are the uncertainty on the non-linearity parameter \(\xi\), the statistical uncertainty on the \(m_{Z}\) measurement (0.073%), and the uncertainty on the \(E/p\) calibration (0.034%). Since the \(m_W\) fit relies predominantly on the shape of the Jacobian edge of the \(m_T\) distribution, the relevant electron transverse energies are in the \(\approx 40 - 45\text{ GeV}\) range. The uncertainty on the energy dependence of the scale from the \(Z\) boson mass is negligible, as the \(\langle p_T^e\rangle\) is about 42 GeV in this sample. The \(E/p\)-based calibration involves an extrapolation from \(\langle p_T^e\rangle = 39\text{ GeV}\), so it receives an additional uncertainty contribution of 23 MeV to the \(m_W\) measurement from the non-linearity parameter \(\xi\). Combining the two calibrations, we obtain a total electron energy measurement uncertainty of 30 MeV on the \(m_W\) measurement in the electron channel. Of this uncertainty, we take 17 MeV to be 100% correlated with the muon channel through the momentum scale uncertainty.

\[\langle m_{ee}\rangle = 91.190\text{ GeV}\]

\[\chi^2/dof = 34 / 38\]

The recoil \(\vec{u}_T\) (Fig. 5) in a W boson event results from quark or gluon radiation in the initial state, and from photon radiation in the initial and final states. A quark or gluon typically fragments into multiple hadrons, which are detected in the calorimeter. Additional energy from the underlying event is also measured in the calorimeter and obscures the recoil measurement. Rather than rely on detailed modeling of the underlying event, we develop an empirical model of the recoil \(\vec{u}_T\) using Z boson events, where the four-momentum of the Z boson is measured precisely using its leptonic decays. The model of the recoil energy measurement is tuned with these decays and applied to W boson events.

We measure the recoil energy using all calorimeter towers except those with ionization or shower energy from the charged leptons. To reduce potential bias and facilitate our model parametrization, we correct the measured energy in each tower for acceptance differences resulting from an uncentered beam. In addition, we improve the measurement resolution by correcting for response differences between the central and plug calorimeters.

\[\langle m_{ee}\rangle = 91.190\text{ GeV}\]

\[\chi^2/dof = 34 / 38\]

\[m_{Z} = (91190 \pm 67)\text{ MeV}\]

\[\chi^2/dof = 34 / 38\]
of an inelastic collision (Section III A 5).

This correction is determined using events collected by applying a multiplicative correction factor to each tower as a function of $|\eta|$. The decreasing curve corresponds to towers closest to the beam line, and the increasing curve corresponds to towers on the opposite side. The correction factors reduce the $\phi$ variation in acceptance due to the beam line displacement from the detector axis.

A. Data Corrections

The data used in this analysis have a relative offset of about 4 mm between the beam line and the center of the CDF II detector. This results in a variation in calorimeter acceptance as a function of $\phi$ such that the calorimeter towers closest to the beam line have a larger acceptance, and thus a larger average measured energy per tower. The variation is largest in the forward region, where the towers are in closest proximity to the beam line. We suppress this azimuthal energy variation by applying a threshold on the combined EM and hadronic tower $p_T$ of 5 GeV for towers with detector $|\eta| > 2.6$. The threshold strongly suppresses the forward tower energy variation in $W$ boson events, while retaining the energy from high-$p_T$ hadronic jets in multijet events. All other towers have EM and hadronic energy thresholds of 20 MeV each.

We reduce the residual azimuthal energy variation by applying a multiplicative correction factor to each measured tower energy according to the following empirical function (Fig. 32):

$$S_{\text{tower}} = 1 - 0.6(0.32|\eta|)^{4.74}\sin(\phi - 0.47).$$  (34)

This correction is determined using events collected by a minimum bias trigger, which requires evidence of an inelastic $p\bar{p}$ collision (Section III A 5). The relative energy scale between the central and forward calorimeters is initially determined from the calibration of high-$p_T$ hadronic jets. The relative response has a significant energy dependence, however, and the initial calibration is not optimized for the low $p_T$ particles relevant to the $W$ boson recoil measurement. Using the $E/p$ distribution of charged pions from minimum bias events, we find that a relative energy scale of $\approx 12\%$ between central and forward calorimeters is appropriate for particles with $p_T \lesssim 2$ GeV, the momentum region of a typical recoil particle. To maintain the mean recoil energy scale, we scale the central (forward) calorimeter tower energies up (down) by 5% (7%). This calibration improves the recoil resolution, and thus the statistical precision of the $m_W$ fits. It also minimizes the sensitivity of the recoil model to differences in phase space sampled by the selected $W$ and $Z$ boson decays.

B. Lepton Tower Removal

The recoil $\vec{u}_T$ is measured as the sum of corrected $\vec{p}_T$ in all calorimeter towers (Sec. VII A), excluding the towers in which the lepton(s) deposit energy. The exclusion of these towers also removes some recoil energy from the measurement, thus causing a bias in $u_\perp$. We measure this bias from the data and incorporate it in the simulation.

An electron shower typically distributes energy to two calorimeter towers, but can also contribute to a third tower if the electron is near a tower edge. We remove each tower neighboring the electron’s tower, as well as the corner towers closest to the electron’s CES position (Fig. 33). A muon near a tower edge can cross two towers, so we remove the two towers in $\eta$ neighboring the muon’s tower (Fig. 34). The tower window definitions are motivated by the presence of excess energy in a given tower above the background energy from the underlying event.

We estimate the recoil energy flow into the excluded towers, denoted by $\Delta u_{||}$, using equivalent windows separated in $\phi$ from the lepton in $W \rightarrow l\nu$ events. When simulating a $W$ or $Z$ boson event, we correct the simulated $\vec{u}_T$ by a $\Delta u_{||}$ taken from the measured distribution. The simulated $\Delta u_{||}$ incorporates its measured dependence on $u_\perp$ and $u_\perp$, and lepton $|\eta|$. These dependencies are shown for $W \rightarrow \mu\nu$ events in Fig. 35 and similar functions are defined for electrons. The incorporation of these functions preserves $\langle \Delta u_{||} \rangle$, which is 260 MeV for electrons and 112 MeV for muons (with negligible statistical uncertainty).

To estimate the systematic uncertainty associated with modeling the tower removal, we study the varia-
FIG. 33: The average energy collected in the electromagnetic (top) and hadronic (bottom) calorimeters in the vicinity of the electron shower in \( W \) boson decays. The differences \( \Delta \phi \) and \( \Delta \eta \) are signed such that positive differences correspond to towers closest to the electron position at shower maximum. The central seven towers inside the box are removed from the recoil measurement. Statistical uncertainties on the values outside the box are \( \mathcal{O}(1 \text{ MeV}) \).

FIG. 34: The average energy collected in the electromagnetic (top) and hadronic (bottom) calorimeters in the vicinity of the muon in \( W \) boson decays. The differences \( \Delta \phi \) and \( \Delta \eta \) are signed such that positive differences correspond to towers closest to the muon position at shower maximum detector. The central three towers inside the box are removed from the recoil measurement. Statistical uncertainties on the values outside the box are \( \mathcal{O}(1 \text{ MeV}) \).

C. Recoil Model Parametrization

The recoil consists of three separate components: radiation in the \( W \) or \( Z \) boson production; radiation from the spectator partons; and energy from additional \( p\bar{p} \) collisions in a given bunch crossing. We use the RESBOS [41] generator to predict the net \( p_T \) distribution of radiation in the \( W \) or \( Z \) boson production, and minimum bias data for the \( p_T \) distribution from spectator partons and additional interactions. The parameters for the detector response to the recoil are measured in \( Z \) boson events.

To facilitate tuning of the recoil model, we define axes such that quark and gluon radiation lies predominantly along one axis, denoted as the “\( \eta \)” axis (Fig. 36). This axis is chosen to be the angular bisector of the two leptons, whose angles are precisely measured. The orthogonal axis is denoted as the “\( \zeta \)” axis.

1. Recoil Energy Scale

We tune the simulation to match the observed detector response to the recoil radiation. The recoil re-
FIG. 35: The scales applied to the energy in the three removed muon towers in the simulation (solid lines), as functions of $u_\parallel$ (top), $u_\perp$ (middle), and $\eta$ (bottom). The points show the scales measured using towers separated in $\phi$ from the muon in $W \rightarrow \mu \nu$ data. The scaling functions for removed electron towers have similar shapes.

The response is defined as $R \equiv \vec{u}_T^{\text{meas}} \cdot \frac{\vec{u}_T^{\text{true}}}{u_T^{\text{true}}}$, where $\vec{u}_T^{\text{true}} = -\vec{p}_T^Z$ is the generated net $\vec{p}_T$ of the initial-state-radiation, and $\vec{u}_T^{\text{meas}}$ is the reconstructed vector of this transverse momentum.

To simulate the measured recoil, we parametrize the response as

$$R(A, B) = A \ln(u_T^{\text{true}} + B) / \ln(15 + B),$$  \hspace{1cm} (35)$$

where $u_T^{\text{true}}$ is in units of GeV, and $A$ and $B$ are constants determined from the data. Figure 37 shows $-\vec{u}_T^{\text{meas}} \cdot \vec{p}_T^{\mu \nu} / p_T^{\mu \nu}$, which approximates $R$, for $Z$ boson decays to muons. The response $R$ is less than 1 due to calorimeter energy loss from particles curling in the tracker, particles passing through calorimeter cracks, and non-linearity of the hadronic calorimeter response.

Projecting the lepton momenta and the recoil along the $\eta$ axis to obtain $p_{\eta}^{ll}$ and $u_\eta$, the sum $p_{\eta}^{ll} + u_\eta$ is sensitive to $R$. This sum is zero for $R = 1$, and positive for $R < 1$. We measure $A = 0.635 \pm 0.007\text{(stat)}$ and $B = 6.68 \pm 1.04\text{(stat)}$ by minimizing the combined $\chi^2$ of the electron and muon $(p_{\eta}^{ll} + u_\eta)$ distributions as a function of $p_{\eta}^{ll}$ (Fig. 35). We determine $A$ and $B$ with the $(p_{\eta}^{ll} + u_\eta)$ distribution rather than the distribution of Fig. 37 because $(p_{\eta}^{ll} + u_\eta)$ is well-defined as $p_T^{Z} \rightarrow 0$ GeV, while $R$ is not. The parameters $A$ and $B$ are statistically uncorrelated by construction. We apply $R(A, B)$ to the generated recoil $\vec{u}_T$ in simulated $W$ and $Z$ boson events.

2. Spectator and Additional $p\bar{p}$ Interactions

The net $\vec{p}_T$ flow from spectator quarks and additional interactions is negligible due to momentum conservation. However, detector resolution causes its measurement to generally be non-zero. The resolution is predominantly determined by the energy sampling in the calorimeter, and we expect it to increase as the square root of the scalar sum $\sum p_T$ of the calorimeter tower $p_T$. We plot the width of the $p_T$ distribution, projected along the $x$ and $y$ axes, as a function of the $\sum p_T$ in minimum bias data. We parametrize the dependence as a power law, with the
FIG. 37: The mean $-\vec{u}_T^{\text{meas}} \cdot \vec{p}_T^{\mu\mu} / p_T^{\mu\mu}$, which approximates the recoil response $R$, as a function of dimuon $p_T$ for $m_{\mu\mu}$ in the $Z$ boson mass region. The simulation (solid) uses parameters fit from the electron and muon $p_T^{\mu\mu} + u_n$ distributions, and models the data (circles) well.

fitted result:

$$\sigma_{x,y} = 0.3842 \left( \sum p_T \right)^{0.5333} \text{GeV}, \quad (36)$$

where $\sum p_T$ is defined in units of GeV. The distribution of $\sum p_T$ from additional interactions, denoted $P_{MB}$, is parametrized as (Fig. 38):

$$P_{MB} \left( \sum p_T \right) \propto \left( \sum p_T \right)^{0.325} e^{-\sum p_T/19.98}, \quad (37)$$

with constants obtained from a fit to the minimum bias data. In our simulation, we draw a value of $\sum p_T$ from this distribution, for the fraction of events containing at least one $p\bar{p}$ collision beyond that producing the $W$ or $Z$ boson. This fraction is calculated from the average instantaneous luminosity of $2.137 \times 10^{31} \text{ cm}^{-2} \text{ s}^{-1}$ for $W$ and $Z$ boson data in the muon (electron) channel, and the assumed instantaneous luminosity per additional collision ($3.3 \times 10^{31} \text{ cm}^{-2} \text{ s}^{-1}$).

The observed $\sum p_T$ from spectator partons in the $p\bar{p} \rightarrow W$ or $Z$ boson interaction is modeled from the minimum bias data, which correspond to one or more $p\bar{p}$ collisions. We deconvolute the $\sum p_T$ spectrum of Eqn. 37 with the distribution of the number of collisions in minimum bias data to derive the following single-collision $\sum p_T$ distribution $P_{1\text{-col.}}$, applicable to $W$ or $Z$ production (Fig. 38):

$$P_{1\text{-col.}} \left( \sum p_T \right) \propto \left( \sum p_T \right)^{0.345} e^{-\sum p_T/14.27}. \quad (38)$$

The $\sum p_T$ produced in a single minimum bias collision can be different from that produced by the spectator partons in $W$ or $Z$ boson production. In order to allow for a difference, we scale the $\sum p_T$ drawn from the single-collision spectrum by a parameter $N_{W,Z}$, which we tune on the $Z$ boson data.

With this model, the $\sum p_T$ in a simulated event is obtained by adding the contributions from the spectator partons and the additional interactions. The corresponding recoil resolution is generated according to Eqn. 36 with a single tunable parameter $N_{W,Z}$.

3. Recoil Energy Resolution

The measurement of the quark and gluon radiation is affected by detector energy resolution, which in
turn affects the measured recoil direction. We model the recoil angular resolution as a Gaussian distribution with \( \sigma_\phi = 0.14 \pm 0.01 \) (stat), determined from fits to the \( \Delta \phi(u_T, -p_T^\mu) \) distribution in Z boson events (Fig. 40). Since the lepton directions are precisely measured, the width of the peak at \( \Delta \phi = 0 \) is dominated by the recoil angular resolution.

The energy resolution of the quark and gluon radiation is predominantly determined by stochastic fluctuations in the hadronic calorimeter, which motivate the functional form \( \sigma_u \propto \sqrt{u_T^{\text{true}}} \). We measure the proportionality constant \( s_{\text{hard}} \) using Z boson data. To tune \( s_{\text{hard}} \) and \( N_{W,Z} \), we project the momentum imbalance \( p_T^{\mu\ell} + u_T \) along the \( \eta \) and \( \xi \) axes in Z boson decays (Fig. 41). The width of these projections as a function of \( p_T^{\mu\ell} \) provides information on \( N_{W,Z} \) and \( s_{\text{hard}} \). At low \( p_T^{\mu\ell} \) the resolution is dominantly affected by \( N_{W,Z} \), with the \( s_{\text{hard}} \) contribution increasing as the boson \( p_T \) increases. We compare the widths of the data and simulation projections as a function of \( p_T^{\mu\ell} \) and compute the \( \chi^2 \). Minimizing this \( \chi^2 \), we obtain \( N_{W,Z} = 1.167 \pm 0.026 \) (stat) and \( s_{\text{hard}} = [0.828 \pm 0.028 \) (stat) \] GeV\(^{1/2} \). The tuning is performed such that the statistical uncertainties on these parameters are uncorrelated.

### D. Recoil Model Cross-Checks

The full recoil model, with parameters tuned from Z boson events, is applied to the simulated W boson sample. We compare the data to the predictions of distributions that can affect the final mass measurement: the projections of the recoil along \( u_{||} \) and perpendicular to \( u_{\perp} \) the charged lepton; and the total recoil \( u_T \).

The \( u_{||} \) distribution is directly affected by the measurements of lepton efficiency as a function of \( u_{||} \) (Figs. 14 and 16) and the modeling of lepton tower removal (Figs. 33 and 34). The \( u_{||} \) is also sensitive to the boson \( p_T \) (Sec. 11.3) and decay angular distributions, and to the recoil response and resolutions.

Since \( u_T \) is much less than the charged lepton \( p_T \) for our event selection, \( \bar{p}_T \approx |p_T + u_{||}| \). Thus, \( m_T \) can be written as:

\[
m_T \approx 2p_T \sqrt{1 + u_{||}/p_T} \approx 2p_T + u_{||} \tag{39}
\]

To a good approximation, any bias in \( u_{||} \) directly en-

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**FIG. 39:** The \( \sum_i p_{T,i} \) distribution for minimum bias data (solid) and a single pp collision (dashed), as derived from the minimum bias distribution.

**FIG. 40:** The angle between the measured recoil and the direction opposite to \( p_T^\mu \), for simulation (histogram) and data (circles) events where the Z boson decays to muons (top) or electrons (bottom). The \( \chi^2 \) from the \( Z \to \mu\mu \) sample is minimized in the fit to the recoil angular resolution. The corresponding uncertainty on \( m_W \) is negligible.
FIG. 41: The simulation (solid) and data (circles) $\eta_{\mu\mu}$ (top) and $\eta_{ee}$ (bottom) resolutions for $Z$ boson decays to muons and electrons. The sum of the four $\chi^2$ values is minimized in the fit for the recoil resolution parameters $N_{W,Z}$ and $s_{\text{hard}}$. Since there are four distributions and two fit parameters, each distribution contributes half a degree of freedom to the fit.

FIG. 42: The simulation (solid) and data (circles) $u_{\parallel}$ distributions for $W$ boson decays to $\mu\nu$ (top) and $e\nu$ (bottom). The simulation uses parameters fit from $Z$ boson data, and the uncertainty on the simulation is due to the statistical uncertainty on these parameters. The data mean ($\mu$) and RMS ($\sigma$) are well-modeled by the simulation.

ters as a bias in the $m_T$ fit. We compare the $u_{\parallel}$ distributions in data and simulation $u_{\parallel}$ in Fig. 42 and observe no evidence of a bias at the level of the data statistics and simulation systematics derived from the recoil model parameters. All backgrounds (Section VIII) are included in the comparison, except $W \rightarrow \tau\nu$, which has similar distributions to the other $W$ leptonic decays.

The $u_{\perp}$ distribution is dominantly affected by the recoil resolution, with a smaller contribution from the recoil response. The simulation models this distribution well for both $W \rightarrow e\nu$ and $W \rightarrow \mu\nu$ samples (Fig. 43).

The mean of the $u_T$ distribution is sensitive to the
FIG. 43: The simulation (solid) and data (circles) $u_\perp$ distributions for $W$ boson decays to $\mu\nu$ (top) and $e\nu$ (bottom). The simulation uses parameters fit from $Z$ boson data, and the uncertainty on the simulation is due to the statistical uncertainty on these parameters. The data mean ($\mu$) and RMS ($\sigma$) are well-modeled by the simulation.

FIG. 44: The simulation (solid) and data (circles) $u_T$ distributions for $W$ boson decays to $\mu\nu$ (top) and $e\nu$ (bottom). The simulation uses parameters fit from $Z$ boson data, and the uncertainty on the simulation is due to the statistical uncertainty on these parameters. The data mean ($\mu$) and RMS ($\sigma$) are well-modeled by the simulation.

VIII. BACKGROUNDS

The event selection criteria (Section IV) result in $W$ boson samples with high purity. However, the small residual backgrounds affect the distributions used for the $m_W$ fits. Both the $W \rightarrow e\nu$ and $W \rightarrow \mu\nu$ samples receive contributions from: $Z/\gamma^* \rightarrow ll$. The recoil response and the boson $p_T$, and is affected to a lesser extent by the resolution. The reverse is the case for the RMS of the $u_T$ distribution. Both are modeled well by the simulation for both $W \rightarrow e\nu$ and $W \rightarrow \mu\nu$ samples (Fig. 44).

The uncertainties on the $m_W$ fits from the recoil parameters (Table IV) are determined by varying each parameter by $\pm 3\sigma$ and assuming linear variation of the fit $m_W$ with the parameter. Since all uncertainties are uncorrelated, we add them in quadrature to obtain total recoil model uncertainties of 12, 17, and 34 MeV on $m_W$ from the $m_T$, $p_T$, and $p_T$ fits, respectively. The uncertainties are the same and 100% correlated for the electron and muon channels, since the recoil parameters are obtained from combined fits to $Z \rightarrow ee$ and $Z \rightarrow \mu\mu$ data. The uncertainty on the $p_T$ fit arises predominantly from the modeling of the $u_T < 15$ GeV threshold used to select $W$ boson events (Section IV).
where one lepton is not detected; $W \rightarrow \tau \nu$, where the \( \tau \) decay products are reconstructed as a charged lepton; and multijet production, where at least one jet is misreconstructed. The $W \rightarrow \mu \nu$ sample also contains backgrounds from cosmic rays, where a muon passing through the COT is reconstructed on only one side of the COT, and long-lived hadrons decaying to $\mu \nu X$, where the muon momentum is misreconstructed.

### Table V: Signed shifts in the $m_W$ fits due to $1\sigma$ increases in the recoil model parameters.

<table>
<thead>
<tr>
<th>Parameter</th>
<th>$m_T$</th>
<th>$p_T$</th>
<th>$p_T$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$A$</td>
<td>$-9$</td>
<td>$-8$</td>
<td>$2$</td>
</tr>
<tr>
<td>$B$</td>
<td>$-2$</td>
<td>$15$</td>
<td>$15$</td>
</tr>
<tr>
<td>$N_{W,Z}$</td>
<td>$5$</td>
<td>$0$</td>
<td>$22$</td>
</tr>
<tr>
<td>$s_{hard}$</td>
<td>$5$</td>
<td>$-3$</td>
<td>$21$</td>
</tr>
</tbody>
</table>

We model the $W \rightarrow \tau \nu$ and $Z/\gamma^* \rightarrow ee$ backgrounds using events generated with PYTHIA \[59\] and simulated with a full GEANT-based detector simulation \[40, 41\]. The full simulation models global detector inefficiencies and is thus more appropriate for predicting background normalizations than the custom fast simulation. The multijet background is estimated using a data-based approach.

In the standard model the branching ratio for $W \rightarrow ev$ is the same as for $W \rightarrow \tau \nu$, neglecting lepton masses. Measurements from LEP \[19\] test this prediction with a precision of 2.9\%, and a slight discrepancy from the standard model is observed with a significance of 2.6\%. In estimating the $W \rightarrow \tau \nu$ background, we assume the standard model prediction and determine the ratio of $W \rightarrow \tau \nu$ to $W \rightarrow ev$ events from the ratio of acceptances of these two processes, as determined by the full GEANT-based detector simulation. We include an uncertainty of 2.9\%, corresponding to the statistical precision of the tests of this assumption. We estimate the $W \rightarrow \tau \nu$ background to be $(0.93 \pm 0.03)\%$ of the $W \rightarrow ev$ sample.

The $Z/\gamma^*$ background is determined from the ratio of $Z/\gamma^* \rightarrow ee$ to $W \rightarrow ev$ acceptances determined from the GEANT-based detector simulation, multiplied by the corresponding ratio of cross sections times branching ratios. The ratio $\sigma \cdot BR(Z \rightarrow ee)/\sigma \cdot BR(W \rightarrow ev)$ has been calculated in the standard model to be $10.69 \pm 0.08$ \[25, 26\], and measurements are consistent with this value \[26, 50\].

![Figure 45](image-url)  
**FIG. 45:** The $p_T$ distribution of the $W \rightarrow ev$ candidate sample (triangles) and prediction (solid), with the $p_T$ and $m_T$ selection cuts removed. We fit for the normalization of the hadronic jet background (dotted) after fixing the normalization of the $W \rightarrow ev$ distribution (dashed) in the peak. Not shown are the $Z/\gamma^* \rightarrow ee$ and $W \rightarrow \tau \nu$ backgrounds, whose relative normalizations are fixed from the simulation.

We take an uncertainty of $\pm 0.43$ on this value from the CDF Run I measurement, and estimate the $Z/\gamma^* \rightarrow ee$ background in the $W \rightarrow ev$ candidate sample to be $(0.24 \pm 0.01)\%$.

Multijet background enters the signal data sample when a hadronic jet is misreconstructed as an electron and a second jet results in large $p_T$ through energy misreconstruction or the semi-leptonic decay of a hadron. To estimate this background, we remove the $p_T$ threshold in our signal event selection to include the background-dominated kinematic region of low $p_T$. We then fit the observed $p_T$ spectrum to the combination of the hadronic jet, $W \rightarrow ev$, $Z/\gamma^* \rightarrow ee$, and $W \rightarrow \tau \nu$ components, floating only the hadronic jet shape normalization (Fig. 45).

In this fit, the shapes and normalizations for the $W \rightarrow ev$, $Z/\gamma^* \rightarrow ee$, and $W \rightarrow \tau \nu$ components are determined from the GEANT-based simulation. The shape of the $p_T$ spectrum of the hadronic jet background is determined from the single-electron events that pass an anti-electron identification requirement based on a neural network discriminant $NN$. The discriminant is determined by combining the electron quality variables (Section IV B) into a neural network \[65\] trained with single-electron data events, using $p_T$ to separate signal from background.

Electron candidates in the $W \rightarrow ev$ sample with low $NN$ values have a high probability to be jets.
misreconstructed as electrons. Events with such candidates provide a $p_T$ distribution characteristic of hadronic jet production. We apply a small correction to this distribution to account for the expected contribution from $W \rightarrow e\nu$ decay electrons with low $N\bar{N}$ values.

This method relies on the assumption that the hadronic jet background has a $p_T$ distribution that is independent of the electron identification variables. As a test of this assumption, we perform the same fit for the jet background normalization, using only the isolation variable (Section IV A) instead of the $N\bar{N}$ to select a hadronic jet subsample. We take a weighted average of the two fitted background normalizations, and assign an uncertainty that covers the range of the two results. The resulting background estimate is $(0.25 \pm 0.15)\%$ of the $W \rightarrow e\nu$ sample.

The $m_T$, $p_T$, and $p_T$ distributions are obtained from the GEANT-based simulation for $W$ and $Z$ boson backgrounds, and from events in the $W \rightarrow e\nu$ sample with low-$N\bar{N}$ electron candidates for the hadronic jet background. We fit these distributions (Fig. 46) and include their shapes and relative normalizations in the $m_W$ template fits. The uncertainties on the background estimates result in uncertainties of 8, 9, and 7 MeV on $m_W$ from the $m_T$, $p_T$, and $p_T$ fits, respectively (Table VI).

<table>
<thead>
<tr>
<th>Background</th>
<th>$%$ of $W \rightarrow e\nu$ data</th>
<th>$\delta m_W$ (MeV)</th>
</tr>
</thead>
<tbody>
<tr>
<td>$W \rightarrow \tau\nu$</td>
<td>0.93 $\pm$ 0.03</td>
<td>$m_T$ fit 2</td>
</tr>
<tr>
<td>Hadronic jets</td>
<td>0.25 $\pm$ 0.15</td>
<td>$p_T$ fit 8</td>
</tr>
<tr>
<td>$Z/\gamma^* \rightarrow ee$</td>
<td>0.24 $\pm$ 0.01</td>
<td>$p_T$ fit 1</td>
</tr>
<tr>
<td>Total</td>
<td>1.42 $\pm$ 0.15</td>
<td>$p_T$ fit 8</td>
</tr>
</tbody>
</table>

TABLE VI: The percentages of the various backgrounds in the $W \rightarrow e\nu$ data set, and the corresponding uncertainties on the $m_T$, $p_T$ and $p_T$ fits for $m_W$.

B. $W \rightarrow \mu\nu$ Backgrounds

The $W \rightarrow \tau\nu$ and $Z/\gamma^* \rightarrow \mu\mu$ backgrounds are modeled using events generated with PYTHIA [54] and simulated with the GEANT [40]-based detector simulation. We use the data to estimate backgrounds from cosmic rays, multijets, and hadrons decaying in flight to $\mu\nu X$.

Backgrounds from $W \rightarrow \tau\nu$ and $Z/\gamma^* \rightarrow \mu\mu$ to the $W \rightarrow \mu\nu$ sample are modeled in the same manner as for the $W \rightarrow e\nu$ sample (Section VIII A). We determine the ratio of the acceptance for $W \rightarrow \tau\nu$
or $Z/\gamma^* \rightarrow \mu\mu$ events to the acceptance for $W \rightarrow \mu\nu$ events using the GEANT-based detector simulation. We assume equal branching ratios for the two $W$ boson decay modes, and use the ratio $\sigma \cdot BR(Z \rightarrow \mu\mu)/\sigma \cdot BR(W \rightarrow \mu\nu) = 10.69 \pm 0.43$ (Section VIII A). We estimate the fraction of $W \rightarrow \tau\nu (Z/\gamma^* \rightarrow \mu\mu)$ events in the $W \rightarrow \mu\nu$ candidate sample to be $(0.89 \pm 0.02)\%$ $(6.6 \pm 0.3)\%$. The $Z/\gamma^* \rightarrow \mu\mu$ background is large because our event selection does not identify muons with $|\eta| \gtrsim 1.2$. The tracker and muon detectors have incomplete or no coverage in the forward rapidity region, and the muons deposit only a few GeV of energy in the calorimeter. Thus, a $Z/\gamma^* \rightarrow \mu\mu$ event with one central and one forward muon is measured as a single-muon event with large $p_T$.

Cosmic-ray muons passing close to the beam line are a source of background to the $W \rightarrow \mu\nu$ sample when the muon track is reconstructed on only one side of the COT. The cosmic-ray identification algorithm searches for unreconstructed tracks and removes cosmic rays with high efficiency. The residual cosmic-ray background is estimated using the reconstructed interaction time $t_0$ and impact parameter $d_0$ from the COT track fit. Figure 47 compares the $t_0$ distributions of the $W \rightarrow \mu\nu$ candidate sample, $Z/\gamma^* \rightarrow \mu\mu$ candidates, and identified cosmic rays. The cosmic ray fraction is fit by minimizing the $\chi^2$ of the sum of the $Z/\gamma^* \rightarrow \mu\mu$ and cosmic ray distributions with respect to the $W \rightarrow \mu\nu$ distribution. We obtain an alternative background estimate by comparing the $d_0$ distribution of identified cosmic rays to the $d_0$ distribution of $W \rightarrow \mu\nu$ candidates with the $d_0$ selection cut removed. The high impact parameter region of the $W$ boson sample is enriched with cosmic rays, and is used to estimate the cosmic ray background within the selection region $|d_0| < 1$ mm. We take the cosmic-ray background to be $(0.05 \pm 0.05)\%$, where the uncertainty covers the range of results from the two estimates.

Decay of a long-lived meson to a muon can result in a reconstructed track with high momentum and large event $p_T$. A low-momentum pion or kaon ($\lesssim 10$ GeV) that decays in the tracking chamber can be reconstructed as a high-momentum muon if the decay is in an azimuthal direction opposite the meson’s curvature (i.e., a kink in the trajectory). Such misreconstruction typically results in a poor COT track $\chi^2$ and a large impact parameter. For each of these quantities we obtain a prompt muon distribution from $Z$ boson decays and a meson decay-in-flight distribution from the $W$ boson sample by requiring either high COT track $\chi^2$ or high impact parameter. We fit for the background fraction by summing the prompt muon distribution with the decay-in-flight distribution, and minimizing the $\chi^2$ with respect to the muon distribution from the $W$ boson sample. We obtain a background fraction of $(0.3 \pm 0.2)\%$, where the uncertainty covers the range of the estimates obtained using the COT track $\chi^2$ and impact parameter distributions.

A separate class of hadronic background results from high-momentum muons from short-lived hadronic decays, or energetic hadrons penetrating the calorimeter to the muon detectors. These background muon candidates are typically accompanied by significant hadronic energy due to an associated hadronic jet, and can be separated using a muon isolation variable. Two such variables are determined by using either calorimeter energy or track momenta in an $\eta - \phi$ cone of size 0.4 surrounding the muon candidate. Using the low $p_T$ region to select a jet-dominated sample, we fit the track and calorimeter isolation distributions of the $W$ boson candidate sample to the sum of the expected distributions from $Z \rightarrow \mu\mu$ events and jet-dominated events. As a third method, we fit the $p_T$ distribution, using muon candidates with high-isolation values to provide the $p_T$ distribution of the hadronic-jet background. From the range of results of the three methods, we obtain a jet background estimate of $(0.1 \pm 0.1)\%$.

The distributions for the $m_{\mu\tau}$ fit variables are obtained from the GEANT-based simulation for $W$ and $Z$ boson backgrounds, from identified cosmic ray events.
for the cosmic ray background, and from events in the $W \to \mu\nu$ sample with high-$\chi^2$ (isolation) muons for the decay-in-flight (hadronic jet) background. Including uncertainties on the shapes of the distributions, the total uncertainties on the background estimates result in uncertainties of 9, 19, and 11 MeV on $m_W$ from the $m_T$, $p_T$, and $p_T/\eta$ fits, respectively (Table VII).

<table>
<thead>
<tr>
<th>Background</th>
<th>$W \to \mu\nu$ data</th>
<th>$m_T$ fit</th>
<th>$p_T$ fit</th>
<th>$p_T/\eta$ fit</th>
</tr>
</thead>
<tbody>
<tr>
<td>$Z/\gamma^* \to \mu\mu$</td>
<td>6.6 ± 0.3</td>
<td>6</td>
<td>11</td>
<td>5</td>
</tr>
<tr>
<td>$W \to \tau\nu$</td>
<td>0.89 ± 0.02</td>
<td>1</td>
<td>7</td>
<td>8</td>
</tr>
<tr>
<td>Decays in flight</td>
<td>0.3 ± 0.2</td>
<td>5</td>
<td>13</td>
<td>3</td>
</tr>
<tr>
<td>Hadronic jets</td>
<td>0.1 ± 0.1</td>
<td>2</td>
<td>3</td>
<td>4</td>
</tr>
<tr>
<td>Cosmic rays</td>
<td>0.05 ± 0.05</td>
<td>2</td>
<td>2</td>
<td>1</td>
</tr>
<tr>
<td>Total</td>
<td>7.9 ± 0.4</td>
<td>9</td>
<td>19</td>
<td>11</td>
</tr>
</tbody>
</table>

TABLE VII: The percentages of the various backgrounds in the $W \to \mu\nu$ data set, and the corresponding uncertainties on the $m_T$, $p_T$, and $p_T/\eta$ fits for $m_W$.

IX. PRODUCTION AND DECAY MODELS

The measurement of the $W$ boson mass relies on a complete model of $W$ and $Z$ boson production and decay. The production process is described by perturbative QCD and a parametrization of non-perturbative QCD effects, with parameters determined from global fits to hadron-hadron and lepton-hadron collision data. $W$ and $Z$ boson decay are modeled using a next-to-leading-order electroweak calculation and includes QCD corrections for the lepton angular distributions, as a function of boson $p_T$. The most important process in the decay is photon radiation off the final-state charged lepton, which has been calculated at next-to-leading order [62].

A. Parton Distribution Functions

The longitudinal momentum of the produced $W$ or $Z$ boson depends on the momenta of the interacting partons. These momenta, generally expressed in terms of the fractions $x_i$ of the colliding (anti-)proton energies, are not known on an event-by-event basis. The $x_i$ parton distribution functions (PDFs) are however well constrained by hadron-hadron and lepton-hadron collision data. The distributions have been parametrized as simple functional forms for the quarks, antiquarks, and gluons inside a proton. Two
independent fits to the global data, performed by the
MRST \cite{22} and CTEQ \cite{23} collaborations, constrain
the parameters in these PDFs.

We model the quark momentum fractions using the
next-to-leading-order CTEQ6M parton distribution
functions. The CTEQ parametrization \cite{23} for most
of the distribution functions inside the proton is:

\[ x_p f_a(x_p, Q_0) = A_0 x_p^{A_1} (1 - x_p)^{A_2} e^{A_3 x_p} (1 + A_4 x_p)^{A_5}, \]

where \( f_a \) are the distributions of a particular quark or
glue combination \( a \), \( A_i \) are the fit parameters, and
\( Q_0 \) is the energy scale at which the parameters are
defined. The functions at a particular energy scale \( Q \)
determine the uncertainty due to that eigenvector.
We calculate the total PDF uncertainty using
the quadrature sum of all eigenvector contributions
\cite{24}:

\[ \delta m_W^{PDF} = \frac{1}{1.6} \left[ \frac{1}{2} \sum_i (m_W^{+i} - m_W^{-i})^2 \right], \]

where \( m_W^{+i} \) represents the mass fits for the \( \pm 1.6 \sigma \)
shifts in eigenvector \( i \). These fits are performed using
templates and simulated pseudoexperiments both
generated with PYTHIA \cite{59}. The resulting \( \delta m_W^{PDF} \)
are 11, 9, and 13 MeV, for the \( m_T, p_T, \) and \( \rho_T \) fits,
respectively. A fit to pseudodata using the MRST
PDF set results in \( m_W \) shifts smaller than these uncer-
tainties.

B. W and Z Boson \( p_T \)

Because mass is a Lorentz invariant, the \( W \) boson
transverse mass is only weakly sensitive to the \( W \)
boson transverse momentum \( p_T^W \). However, the decay
lepton \( p_T \) spectra are more significantly affected by
the \( p_T^W \) distribution.

At the Tevatron, the \( p_T \) spectra of \( W \) and \( Z \) bosons
peak at a few GeV (Fig. 19), where the shapes are
predominantly determined by non-perturbative QCD
interactions. We model the distribution with the
RESBOS generator \cite{61}, which uses the Collins-Soper-
Sterman (CSS) \cite{69} resummation formalism and a
parametrized non-perturbative form factor. In this
formalism, the cross section for \( W \) boson production is
written as:

\[
\frac{d\sigma(p\bar{p} \rightarrow W + X)}{d\hat{s}d^2\hat{p}_T^W dy} = \frac{1}{2\pi^2} \delta(\hat{s} - m_W^2) \int d^2\hat{p}_T^W \hat{b} \times \hat{W}_{jk}(\hat{b}, \hat{s}, x_i) + Y(p_T^W, \hat{s}, x_i),
\]

where \( x_i \) are the parton energy fractions of the (anti-)
proton, \( y = 0.5 \ln(x_p/x_{\bar{p}}) \) is the boson rapidity, \( \hat{b} \)
is the relative impact parameter of the partons in the
collision, \( Y \) is a function calculable at fixed order, and
\( \hat{W} \) can be separated into its perturbative and non-perturbative
components. We use the Brock-Landry-Nadolsky-Yuan (BLNY)
form for the non-perturbative component:

\[
\hat{W}_{jk}^{NP} = e^{g_1 - g_2 \ln(\frac{x_p}{x_{\bar{p}}}) - g_3 \ln(100x_p x_{\bar{p}})},
\]

where \( Q_0 = 1.6 \text{ GeV} \) and \( g_i \) are parameters suggested
by the CSS formalism to be universal to processes with
initial state quarks and colorless objects in the
final state \cite{61}.

The \( g_2 \) parameter affects the position of the most
probable \( p_T^W \) and is the most relevant for the \( m_W \)
measurement. We use \( g_1 = 0.21 \text{ GeV}^2 \), \( g_2 = 0.68 \text{ GeV}^2 \),
and \( g_3 = -0.60 \), which are determined from
fits to global Drell-Yan data \cite{61}. We verify the app-
licability of these values to our data by fitting the
dilepton \( p_T \) distribution (Fig. 19) for \( g_2 \). We find
\( g_2 = (0.685 \pm 0.048 \text{ (stat)}) \text{ GeV}^2 \), consistent with the
global fits. Varying \( g_2 \) by \( \pm 3\sigma \) in pseudoexperiments
and taking the fit \( m_W \) to be linearly dependent on
\( g_2 \), we find that the uncertainty of \( \delta g_2 = 0.048 \text{ GeV}^2 \)
results in uncertainties of 3, 9, and 5 MeV, on \( m_W \) for
the \( m_T, p_T, \) and \( \rho_T \) fits, respectively. These uncer-
tainties are the same and 100\% correlated between
the electron and muon channels, since \( g_2 \) is fit us-
ing \( Z \rightarrow ee \) and \( Z \rightarrow \mu\mu \). Neglecting correlations
between PDFs and the \( g_i \) parameters, we find that
uncertainties on the other \( g_i \) do not significantly af-
fect the \( m_W \) measurement.

C. W Boson Decay

The \( m_W \) measurement is sensitive to the charged
lepton decay angle relative to the boson \( p_T \). The
mismodeling of this angle can bias the projection of the
recoil along the lepton (\( u_{ij} \)), which in turn affects
\( m_W \) measured from the \( m_T \) fit (Section VII D).

The lepton decay angle is predicted by the matrix
element calculation in the RESBOS generator, which
computes the differential cross section $\frac{d^4\sigma}{dp_T^2\ d\psi\ dQ^2\ d\Omega}$. 

The angular distributions are defined in the Collins-Soper rest frame of the $W$ boson [71]. In this frame, the $z$-axis is defined to bisect the angle between the proton momentum and the opposite of the antiproton momentum.

The angular component of the differential cross section can be written as [71],

\[
\frac{d\sigma}{d\Omega} \propto (1 + \cos^2 \theta) + \frac{1}{2} A_0(1 - 3 \cos^2 \theta) 
+ A_1 \sin 2\theta \cos \phi + \frac{1}{2} A_2 \sin^2 \theta \cos 2\phi
+ A_3 \sin \theta \cos \phi + A_4 \cos \theta + A_5 \sin^2 \theta \sin 2\phi
+ A_6 \sin 2\theta \sin \phi + A_7 \sin \theta \sin \phi, \tag{44}
\]

where the $A_i(p_T^W, y_W)$ have been calculated to NLO in $\alpha_s$. Because of the $V-A$ structure of the electroweak interaction, for leading-order valence quark interactions all $A_i$ are zero except $A_4 = 2$. The $A_i$ can be determined experimentally through a moments analysis [72] of the lepton angle in the Collins-Soper frame.

We have performed a moments analysis to extract the $A_i$ from the RESBOS generator, and compared the results to those obtained [72] from the DYRAD event generator [73], which produces $W+$ jet events to order $\alpha_s^2$. The two generators give consistent results in the overlapping region 15 GeV < $p_T^W$ < 100 GeV.

D. Photon Radiation

The quarks, the $W$ boson, and charged lepton have non-zero electromagnetic charge and can radiate photons in the $W$ boson production process. Radiation off the initial-state quarks and the $W$ boson propagator have a negligible effect on the invariant mass distribution of the $W$ boson. Radiation off the final-state charged lepton reduces the measured transverse mass (relative to the $W$ boson mass) and must be accurately modelled.

We study photon radiation using the WGRAD event generator [62], which models the full next-to-leading-order (NLO) electroweak physics. The generator allows an independent study of photon radiation from the initial-state quarks (ISR), the $W$ boson propagator, and the final-state charged lepton (FSR). Interference between the contributing diagrams can also be studied independently. We verify that the initial-state, propagator, and interference effects do not affect the measured $W$ boson mass, within the 5 MeV statistical uncertainty of the simulation.

We simulate final-state photon radiation in our RESBOS-generated $W$ and $Z$ boson events by generating a photon for each charged lepton. The energy and angular spectra are taken from the WGRAD generator using the appropriate boson mass. To avoid the infrared divergence that arises when the photon momentum goes to zero, we require $E_\gamma > \delta_s \sqrt{\delta_s}/2$, where $\delta_s = 10^{-4}$. We find that increasing $\delta_s$ to $10^{-3}$.
FIG. 50: The cube root of the fraction of electron (solid) or muon (dashed) momentum contained in the radiated photon (top), and the square root of the angle $\Delta R$ (bottom) between the radiated photon and the electron (solid) or muon (dashed).

The energy of a photon in a given event is calculated from the fraction $y = E_\gamma / E_l$. The photon angle $\Delta R = \sqrt{(\Delta \eta)^2 + (\Delta \phi)^2}$ is taken with respect to the charged lepton. To improve the phase space sampling, we sample from a two-dimensional distribution of the variables $\sqrt{\Delta R}$ and $y^{1/3}$ when selecting a photon. The individual distributions of these variables are shown in Fig. 50.

We validate our photon simulation by fitting a sample of events generated with FSR using WGRAD to templates generated with leading-order WGRAD and photons simulated according to our model. We find our FSR model to be consistent with that of WGRAD at the level of the 5 MeV statistical precision of the test. The total effect of including FSR is shown in Table VIII Since FSR reduces the charged lepton momentum, the shift is largest for the $p_T$ fit and smallest for the $p_T$ fit. The effects are smaller for electrons than for muons because the electron calorimeter energy measurement recovers much of the energy of FSR photons.

We approximate the effect of next-to-next-to-leading-order FSR by increasing the photon’s momentum fraction ($y$) by 10%, consistent with the results of a study of higher-order photon radiation [74]. We take half the correction as a systematic uncertainty to account for higher-order QED effects.

The total uncertainty due to photon radiation is the quadrature sum of: uncertainties on ISR, interference between ISR and FSR, and radiation off the propagator (5 MeV); uncertainty due to the infrared cutoff of the FSR photon (5 MeV); the FSR model (5 MeV); and uncertainties on higher-order FSR corrections (7 MeV for the electron and 8 MeV for the muon $m_T$ fits). The total uncertainties are 12 (11), 13 (13), and 10 (9) MeV, for the muon (electron) $m_T$, $p_T$, and $p_T$ fits, respectively.

### TABLE VIII: The mass shifts obtained by fitting events generated with our simulation of single-photon radiation to templates generated without final-state photon radiation. The shifts are for the W boson $m_T$, $p_T$, and $p_T$ fits, and for the Z boson $m_{ll}$ fit. The shifts have statistical uncertainties of 7 MeV each.

<table>
<thead>
<tr>
<th>Fit</th>
<th>$\delta m_{W,Z}(\mu)$ (MeV)</th>
<th>$\delta m_{W,Z}(e)$ (MeV)</th>
</tr>
</thead>
<tbody>
<tr>
<td>$m_T$</td>
<td>-158</td>
<td>-138</td>
</tr>
<tr>
<td>$p_T$</td>
<td>-206</td>
<td>-186</td>
</tr>
<tr>
<td>$p_T$</td>
<td>-77</td>
<td>-59</td>
</tr>
<tr>
<td>$m_{ll}$</td>
<td>-196</td>
<td>-215</td>
</tr>
</tbody>
</table>

X. **W BOSON MASS FITS**

We fit the W boson data distributions to a sum of background and simulated signal templates of the $m_T$, $p_T$, and $p_T$ distributions, fixing the normalization of the sum to the number of data events. The fit minimizes the negative log likelihood (Section II D) as a function of the template parameter $m_W$, which is defined by the relativistic Breit-Wigner mass distribution [11]:

$$\frac{d\sigma}{dm} \propto \frac{m^2}{(m^2 - m_W^2)^2 + m^4 \Gamma_W^2 / m_W^4},$$

where $m$ is the invariant mass of the propagator. The
likelihood is calculated in $m_W$ steps of 1 MeV. We use the standard model $W$ boson width $\Gamma_W = 2.094$ GeV, which has an accuracy of 2 MeV and is calculated for $m_W = 80.393$ GeV. Using pseudoexperiments, we find the input $\Gamma_W$ affects the fit $m_W$ according to the relation $dm_W/d\Gamma_W = 0.14 \pm 0.04$.

### A. Fit Results

The results of the $m_T$ fits are shown in Fig. 51 and Table IX gives a summary of the 68% confidence level uncertainties associated with the fits. We fit for $m_W$ in the range $65$ GeV < $m_T$ < $90$ GeV, where the fit range has been chosen to minimize the total uncertainty on $m_W$. The $p_T$ and $\not{p}_T$ distributions are fit in the range $32$ GeV < $p_T$ < $48$ GeV (Figs. 52 and 53, respectively) and have uncertainties shown in Tables X and XI, respectively. We show the individual fit results in Table XII, and the negative log-likelihoods of all fits in Fig. 54.

<table>
<thead>
<tr>
<th>Source</th>
<th>$W \rightarrow \mu \nu$</th>
<th>$W \rightarrow e \nu$</th>
<th>$W \rightarrow \mu \nu$ and $W \rightarrow e \nu$ Correlation</th>
</tr>
</thead>
<tbody>
<tr>
<td>Tracker Momentum Scale</td>
<td>17</td>
<td>17</td>
<td>100%</td>
</tr>
<tr>
<td>Calorimeter Energy Scale</td>
<td>0</td>
<td>25</td>
<td>0%</td>
</tr>
<tr>
<td>Lepton Resolution</td>
<td>3</td>
<td>9</td>
<td>0%</td>
</tr>
<tr>
<td>Lepton Efficiency</td>
<td>1</td>
<td>3</td>
<td>0%</td>
</tr>
<tr>
<td>Lepton Tower Removal</td>
<td>5</td>
<td>8</td>
<td>100%</td>
</tr>
<tr>
<td>Recoil Scale</td>
<td>9</td>
<td>9</td>
<td>100%</td>
</tr>
<tr>
<td>Recoil Resolution</td>
<td>7</td>
<td>7</td>
<td>100%</td>
</tr>
<tr>
<td>Backgrounds</td>
<td>9</td>
<td>8</td>
<td>0%</td>
</tr>
<tr>
<td>PDFs</td>
<td>11</td>
<td>11</td>
<td>100%</td>
</tr>
<tr>
<td>$W$ Boson $p_T$</td>
<td>3</td>
<td>3</td>
<td>100%</td>
</tr>
<tr>
<td>Photon Radiation</td>
<td>12</td>
<td>11</td>
<td>100%</td>
</tr>
<tr>
<td>Statistical</td>
<td>54</td>
<td>48</td>
<td>0%</td>
</tr>
<tr>
<td>Total</td>
<td>60</td>
<td>62</td>
<td></td>
</tr>
</tbody>
</table>

**TABLE IX:** Uncertainties in units of MeV on the transverse mass fit for $m_W$ in the $W \rightarrow \mu \nu$ and $W \rightarrow e \nu$ samples.

We combine results from the $W \rightarrow \mu \nu$ and $W \rightarrow e \nu$ fits using the Best Linear Unbiased Estimator (BLUE) [7]. The BLUE algorithm defines a procedure for constructing a complete covariance matrix using the derivative of $m_W$ with respect to each model parameter [18]. We construct this matrix assuming each source of systematic uncertainty is independent of any other source of uncertainty. The resulting covariance matrix (Table XIII) is then used to combine all six $m_W$ fits. When combining any subset of fits, the appropriate smaller covariance matrix is used.

The result of combining the $m_W$ fits to the $m_T$ distribution in the $W \rightarrow \mu \nu$ and $W \rightarrow e \nu$ channels is

$$m_W = 80.417 \pm 0.048 \text{ GeV.} \quad (46)$$

The $\chi^2$/dof of the combination is 3.2/1 and the probability that two measurements of the same quantity would have a $\chi^2$/dof at least as large as this is 7%.

The combination of the fits to the $p_T$ distribution yields

$$m_W = 80.388 \pm 0.059 \text{ GeV,} \quad (47)$$
with a $\chi^2$/dof of 1.8/1 and an 18% probability for the two measurements to obtain a $\chi^2$/dof $\geq 1.8$.

The results of the fits to the $p_T$ distribution gives

$$m_W = 80.434 \pm 0.065 \text{ GeV}, \quad (48)$$

with a 43% probability of obtaining a $\chi^2$/dof at least as large as observed (0.6/1).

Combining the $m_T$, $p_T$, and $p_T$ fits within the individual decay channels gives $m_W = (80.352 \pm 0.060)$ GeV with a $\chi^2$/dof of 1.4/2 for the $W \rightarrow \mu\nu$ channel and $m_W = (80.477 \pm 0.062)$ GeV with a $\chi^2$/dof of 0.8/2 for the $W \rightarrow e\nu$ channel.

We combine the six fits with the BLUE procedure to obtain our final result of

$$m_W = 80.413 \pm 0.048 \text{ GeV}, \quad (49)$$

which has statistical and systematic uncertainties of 34 MeV each. The statistical correlations between the fits, determined from simulation pseudoexperiments, are shown in Table XIV. The relative weights of the fits are 47.7% (32.3%), 3.4% (8.9%), 0.9% (6.8%) for the $m_T$, $p_T$, and $p_T$ fit distributions, respectively, in the muon (electron) channel. The combination establishes an a priori procedure to incorporate all the information from individual fits, and yields a $\chi^2$/dof of 4.8/5. The probability to obtain a
FIG. 53: The simulation (solid) and data (points) $p_T$ distributions for $W$ boson decays to $\mu\nu$ (top) and $e\nu$ (bottom). The simulation corresponds to the best-fit $m_W$, determined using events between the two arrows. The uncertainty is statistical only.

FIG. 54: The negative log of the likelihood ratio $L/L_0$, where $L_0$ is the maximum likelihood, as a function of $m_W$ for the $m_T$ (solid), $p_T$ (dashed), and $p_T^*$ (dotted) fits in the electron (top) and muon (bottom) channels.

<table>
<thead>
<tr>
<th>Distribution</th>
<th>$m_W$ (GeV)</th>
<th>$\chi^2$/dof</th>
</tr>
</thead>
<tbody>
<tr>
<td>$m_T(e,\nu)$</td>
<td>80.493 ± 0.048 ± 0.039</td>
<td>86/48</td>
</tr>
<tr>
<td>$p_T(e)$</td>
<td>80.451 ± 0.058 ± 0.045</td>
<td>63/62</td>
</tr>
<tr>
<td>$\not{p}_T(e)$</td>
<td>80.473 ± 0.057 ± 0.054</td>
<td>63/62</td>
</tr>
<tr>
<td>$m_T(\mu,\nu)$</td>
<td>80.349 ± 0.054 ± 0.027</td>
<td>59/48</td>
</tr>
<tr>
<td>$p_T(\mu)$</td>
<td>80.321 ± 0.066 ± 0.040</td>
<td>72/62</td>
</tr>
<tr>
<td>$\not{p}_T(\mu)$</td>
<td>80.396 ± 0.066 ± 0.046</td>
<td>44/62</td>
</tr>
</tbody>
</table>

<table>
<thead>
<tr>
<th>$m_T(e,\nu)$</th>
<th>$m_T(\mu,\nu)$</th>
<th>$p_T(e)$</th>
<th>$p_T(\mu)$</th>
<th>$\not{p}_T(e)$</th>
<th>$\not{p}_T(\mu)$</th>
</tr>
</thead>
<tbody>
<tr>
<td>$m_T(e,\nu)$</td>
<td>64$^2$</td>
<td>27$^2$</td>
<td>61$^2$</td>
<td>27$^2$</td>
<td>61$^2$</td>
</tr>
<tr>
<td>$m_T(\mu,\nu)$</td>
<td>61$^2$</td>
<td>27$^2$</td>
<td>59$^2$</td>
<td>28$^2$</td>
<td>59$^2$</td>
</tr>
<tr>
<td>$p_T(e)$</td>
<td>75$^2$</td>
<td>35$^2$</td>
<td>51$^2$</td>
<td>32$^2$</td>
<td></td>
</tr>
<tr>
<td>$p_T(\mu)$</td>
<td>77$^2$</td>
<td>32$^2$</td>
<td>53$^2$</td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\not{p}_T(e)$</td>
<td>81$^2$</td>
<td>43$^2$</td>
<td></td>
<td></td>
<td></td>
</tr>
<tr>
<td>$\not{p}_T(\mu)$</td>
<td>81$^2$</td>
<td></td>
<td></td>
<td></td>
<td></td>
</tr>
</tbody>
</table>

TABLE XII: The results of the fits for $m_W$ to the $m_T$, $p_T$, and $\not{p}_T$ distributions in the electron and muon decay channels. The first uncertainty is statistical and the second is systematic.

TABLE XIII: The complete covariance matrix for the $m_T$, $p_T$, and $\not{p}_T$ fits in the electron and muon decay channels, in units of MeV$^2$. The matrix is symmetric.

$\chi^2$ at least as large as this is 44%.

B. Cross-Checks

Figures 55-60 show the differences between data and simulation, divided by the statistical uncertainties on the predictions, for the $m_T$, $p_T$ and $\not{p}_T$ distributions. Figures 58-60 show the variations
TABLE XIV: The statistical correlations between the \( m_T, p_T, \) and \( \bar{p}_T \) fits in the electron and muon decay channels.

<table>
<thead>
<tr>
<th>Correlation ( W \rightarrow \mu \nu )</th>
<th>( W \rightarrow e \nu )</th>
</tr>
</thead>
<tbody>
<tr>
<td>( m_T - p_T )</td>
<td>70% 68%</td>
</tr>
<tr>
<td>( m_T - \bar{p}_T )</td>
<td>72% 63%</td>
</tr>
<tr>
<td>( p_T - \bar{p}_T )</td>
<td>38% 17%</td>
</tr>
</tbody>
</table>

The largest systematic uncertainty in this difference arises in the muon channel from the uncertainty on the alignment parameters \( a_0 \) and \( a_2 \) (Table III). The uncertainties on the mass difference due to these parameters are 49 MeV and 56 MeV, respectively, for a total uncertainty of 75 MeV. Any bias in these parameters affects the positive and negative lepton fits in opposite directions, and thus has a negligible net effect when the two are combined.

XI. SUMMARY

We have performed a measurement of the \( W \) boson mass using 200 pb\(^{-1}\) of data collected by the CDF II detector at \( \sqrt{s} = 1.96 \) TeV. From fits to \( m_T, p_T, \) and \( \bar{p}_T \) distributions of the \( W \rightarrow \mu \nu \) and \( W \rightarrow e \nu \) data...
FIG. 57: The difference between the data and simulation, divided by the statistical uncertainty on the data points, for the $p_T$ distributions in the muon (top) and electron (bottom) channels. The arrows indicate the fit region.

Fit difference $W \rightarrow \mu \nu$ $W \rightarrow e \nu$

$\Delta m_W(l^+) - \Delta m_W(l^-)$ $286 \pm 152$ $257 \pm 117$

$\Delta m_W(\phi_l > 0) - \Delta m_W(\phi_l < 0)$ $0 \pm 133$ $116 \pm 117$

$m_W$(Mar, 2002-Apr, 2003) - $m_W$(Apr, 2003-Sep, 2003) $75 \pm 135$ $-107 \pm 117$

TABLE XV: Differences of $m_W$ in the $p_T$ fits between positively and negatively charged leptons, leptons in the upper and lower halves of the detector, and early and late data. The units are MeV.

samples, we obtain

$$m_W = 80.413 \pm 0.048 \text{ GeV},$$

which is the single most precise determination of $m_W$ to date. The uncertainty includes statistical and systematic contributions of 34 MeV each.

Combining this result with the Run I Tevatron measurements using the method in \[76\], we obtain a CDF Run I/II combined result of

$$m_W = 80.418 \pm 0.042 \text{ GeV},$$

and a combined Tevatron result of

$$m_W = 80.429 \pm 0.039 \text{ GeV}. \quad (52)$$

In these combinations, we take the uncertainties due to PDFs and photon radiation to be fully correlated between our measurement and the previous Tevatron measurements. In the BLUE combination method \[75\], each uncertainty source contributes its covariance matrix, and all covariance matrices are summed
to obtain the total covariance matrix. We evaluate
an individual contribution to the uncertainty on our
result by ignoring its respective covariance matrix
and repeating the six-fold combination of our individ-
ual electron and muon channel $m_T$, $p_T$ and $\not{p}_T$ fits.
The difference in quadrature between the total un-
certainty, including and excluding a given covariance
matrix contribution, is taken to be the uncertainty
due to that source. Following this procedure, we ob-
tain the systematic uncertainty contributions due to
PDFs and QED radiative corrections to be 12.6 MeV
and 11.6 MeV respectively (Table XVI), for a com-
bined uncertainty of 17.2 MeV.

Assuming no correlation between the Tevatron and

LEP measurements, we obtain a new world average of

$$m_W = 80.398 \pm 0.025 \text{ GeV.} \quad (53)$$

Our measurement reduces the world uncertainty to
31 parts in $10^5$, and further constrains the properties
of the Higgs boson and other new particles coupling
to the $W$ and $Z$ bosons. Within the context of the
standard model, fits made to high energy precision
electroweak data in 2006 gave $m_H = 85^{+39}_{-28}$ GeV, with
$m_H < 166$ GeV at the 95% confidence level [19]. The
values used for the top quark and $W$ boson masses
in these fits were $m_t = (171.4 \pm 2.1)$ GeV and $m_W =
(80.392 \pm 0.029)$ GeV, respectively. Updating these
the radiation length, the cross section is \( \rho \) where \( X_a \) and \( \alpha \).

\[
\rho \sim \frac{1}{X_a} \equiv \left( \frac{4}{3} + \frac{C}{y} \right) \frac{1}{y} - 1 + y,
\]

where

\[
\psi_1(Z) = \frac{Z^2}{Z \ln(1194Z - 2/3)} - f + \frac{Z \ln(1194Z - 2/3)}{Z \ln(1194Z - 2/3)},
\]

\[
\psi_2(Z) = Z^2 + Z,
\]

\[
f = a^2 [(1 + a^2) - 0.20206 - 0.0369a^2 + 0.0083a^4 - 0.002a^6],
\]

and \( a = \alpha_{EM}Z \). We define a material’s radiation length \( X_0 \) according to [48]:

\[
X_0^{-1} = 4\alpha_{EM}r_e^2 NA \rho \psi_1(Z)/A,
\]

where \( \rho \) is the density of the material. In terms of the radiation length, the cross section is

\[
\frac{d\sigma}{dy} = \frac{A}{NA X_0 \rho} \left[ \left( \frac{4}{3} + C \right) \left( \frac{1}{y} - 1 \right) + y \right],
\]

where

\[
C \equiv \frac{\psi_2(Z)}{\psi_1(Z)}.
\]

The conversion cross section takes a similar form, since the relevant Feynman diagram is a rotation of the bremsstrahlung process [48]:

\[
\frac{d\sigma}{dy} = \frac{A}{NA X_0 \rho} \left[ 1 - (4/3 + C) y(1 - y) \right].
\]

The Compton scattering cross section as a function of scattering angle is given by the Klein-Nishina formula [54]:

\[
\frac{d\sigma}{d\Omega} = \frac{r_e^2}{2} \left[ \frac{1 + \cos^2 \theta}{[1 + k(1 - \cos \theta)]^2} + \frac{k^2(1 - \cos \theta)^2}{[1 + k(1 - \cos \theta)]^5} \right],
\]

where \( k \equiv E_\gamma/m_e \). The scattering angle is kinematically related to the energy loss by [54]

\[
y = k'/k = [1 + k(1 - \cos \theta)]^{-1},
\]

where \( k' \) is the energy of the photon after scattering, in units of \( m_e \). Using this equation, the differential cross section with respect to \( y \) can be written as [10]:

\[
\frac{d\sigma}{dy} \propto 1/y + y,
\]
ignoring terms containing $1/k$.

**ACKNOWLEDGMENTS**

We thank the Fermilab staff and the technical staffs of the participating institutions for their vital contributions. This work was supported by the U.S. Department of Energy and National Science Foundation; the Italian Istituto Nazionale di Fisica Nucleare; the Ministry of Education, Culture, Sports, Science and Technology of Japan; the Natural Sciences and Engineering Research Council of Canada; the National Science Council of the Republic of China; the Swiss National Science Foundation; the A.P. Sloan Foundation; the Bundesministerium für Bildung und Forschung, Germany; the Korean Science and Engineering Foundation and the Korean Research Foundation; the Science and Technology Facilities Council and the Royal Society, UK; the Institut National de Physique Nucleaire et Physique des Particules/CNRS; the Russian Foundation for Basic Research; the Comisión Interministerial de Ciencia y Tecnología, Spain; the European Community’s Human Potential Programme; the European Commission under the Marie Curie Programme; the Slovak R&D Agency; and the Academy of Finland.


[22] There is in principle a correlation due to final-state photon radiation and the $W$ boson width, but these are sufficiently small that they are ignored when combining $m_W$ measurements.


[27] In order to obtain thesis results, two of the authors removed the blinding offset at two different points in the analysis. Access to these theses was denied to all authors outside of the University of Toronto.


[36] For a small fraction of the data, only 10 hits out of 12 sense wires were required to create XFT segments.


[39] The underlying event refers to the spectator parton and additional inelastic pp interactions that produce low pt particles roughly uniform in phase space.


[42] The map is divided into 333 longitudinal and 120-1000 azimuthal sections, with the number of azimuthal sections increasing as radius increases.


[45] The radial distribution of the hits has a small impact on parameter resolution, with the importance depending on the parameter. We do not attempt to model the radial hit distribution.


[49] Increasing γp from 10^{-4} to 10^{-5} has about a 50 MeV effect on the E/p calibration for electrons from W boson decays; reducing it to 5 x 10^{-5} has less than a 5 MeV effect.


[55] Because the model of the width of the E/p peak depends on the model of the momentum resolution, there is a 0.13% systematic uncertainty in γ associated with the momentum resolution model.


[63] We do not require the correction factor for the number of radiation lengths (0.4 ± 0.9)% to be the same as that for the ionization energy loss (-6 ± 2)% because of our assumption that copper cable needs to be added to the GEANT simulation to model the radial distribution of photon conversions. Since the bremsstrahlung and conversion cross sections scale as X_e^{-1} \propto Z^2 (Appendix A) and the ionization energy loss scales as Z (Section III B 1), the correction factors can be different if the type of material is incompletely known.

[64] The statistical significance of the slope in the Z boson sample is 0.5σ. Averaging with the W boson sample results in zero slope.


