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Intrinsic electrophysiological properties of entorhinal cortex stellate cells and their contribution to grid cell firing fields

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INTRODUCTION

Stellate cells in layer II of the medial entorhinal cortex (MEC) are a crucial component of the neural circuit for representation of space. These neurons receive synaptic input from diverse cortical areas and send axonal projections to the dentate gyrus of the hippocampus (Steward and Scoville, 1976; Schwartz and Coleman, 1981; Ruth et al., 1982; Ruth and Collier, 1988; Insauti et al., 1997; Dolorfo and Amaral, 1998; Burwell, 2000; van Groen et al., 2003; Witter, 2007). They are defined by a stellate dendritic architecture and distinctive electrophysiological properties (Ramón y Cajal, 1995; Alonso and Llinás, 1989; Alonso and Klink, 1993; Jones, 1994; Klink and Alonso, 1997a; Heinemann et al., 2000; Erchova et al., 2004; Burton et al., 2008; Garden et al., 2008). Stellate cells received considerable attention initially because of their possible role in theta-frequency (4–12 Hz) oscillations that are associated with exploratory spatial behaviors (Mitchell and Ranck, 1980; Alonso and García-Austt, 1987; Dickson et al., 1995; Buzsáki, 1996; White et al., 1998; Hasselmo et al., 2002).

The discovery that cells in layer II of the MEC have grid-like spatial firing fields has given further motivation to investigate of the functional properties of stellate cells (Fyhn et al., 2004; Hafting et al., 2005). Several observations suggest that grid cells in layer II of the MEC are stellate cells. Stellate cells are the main excitatory neuron type in the layer where grid cells are most frequently found (Alonso and Klink, 1993; Hafting et al., 2005; Sargolini et al., 2006), and they were recently shown to encode spatial information during navigation on linear tracks (Burgalossi et al., 2011). The topographical organization of the spatial resolution of grid firing fields along the dorsal-ventral axis of the MEC (Hafting et al., 2005; Sargolini et al., 2006; Barry et al., 2007; Brun et al., 2008; Fyhn et al., 2008) is mirrored by differences between stellate cells in their intrinsic theta-frequency activity and resonance (Giocomo et al., 2007; Boehlen et al., 2010; Dodson et al., 2011; Yoshida et al., 2011), and in their integration of synaptic responses (Garden et al., 2008). The intrinsic electrophysiological properties of stellate cells may, therefore, offer vital clues and experimental targets for the investigation of cellular mechanisms involved in the neuronal representation of space (Giocomo et al., 2011b; O’Donnell and Nolan, 2011).

Keywords: ion channel, grid cell, HCN, synaptic integration, oscillation, theta, resonance
With this goal in mind we review the known intrinsic electrophysiological properties of stellate cells in layer II of the MEC and introduce new data obtained using perforated patch-clamp methods. We begin by considering the identity and nature of the ion channels that stellate cells express. We then examine how these channels influence membrane properties of stellate cells typically measured with recordings from brain slices. We compare the voltage dependence of key membrane properties of stellate cells (Figures 1A,B) with that of their ion channels (Figure 1C). We show that new data obtained with the perforated patch-clamp method can reconcile differences between results obtained with whole-cell and sharp electrode recordings (Figures 2–9). Finally, we consider evidence for how the intrinsic properties of stellate cells influence neuronal computations important for the representation of space in behaving animals. We argue that models that require stellate cells to generate periodic oscillations are not consistent with existing data. We suggest that establishing how stellate cells integrate their synaptic input (Garden et al., 2008) will be crucial for understanding the cellular basis for grid cell firing fields.

RESULTS AND DISCUSSION
ION CHANNELS DETERMINING THE INTRINSIC EXCITABILITY OF STELLATE CELLS
Sodium channels

Three types of sodium conductance have been identified in stellate cells (Figure 1C). Depolarization to membrane potentials above approximately –50 mV activates a transient sodium current (NaT) with rapid kinetics (White et al., 1993; Magistretti and Alonso, 1999b, 2007; Magistretti et al., 1999b). This current consists of a component that is sensitive to the classic Na+ channel blocker tetrodotoxin (TTX) and a component that is relatively insensitive (White et al., 1993). The kinetics and voltage dependence of these two components are indistinguishable, but the TTX-resistant component of NaT is more sensitive to block by Cd2+, La3+, and Zn2+ (White et al., 1993). The sensitivity to Zn2+ may be of functional relevance as Zn2+ is found at high levels in the neuropil of the MEC (Haug, 1976; Holm and Geneser, 1989; Slomianka, 1992). The TTX-insensitive component of NaT is smaller than the TTX-sensitive component (White et al., 1993).

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**FIGURE 1** | Summary of intrinsic electrophysiological properties of stellate cells from layer II of the medial entorhinal cortex. (A) Example of membrane potential activity at rest and during injection of constant current. Shading indicates the range of typical resting membrane potentials (orange), peri-threshold theta-frequency activity (purple), and threshold for initiation of action potential firing (blue). Subthreshold activity is shown on an expanded voltage-scale below. (B) Examples of excitatory postsynaptic potentials evoked by stimuli of increasing amplitude (left to right). Each block shows five consecutive responses to stimuli of identical intensity. Data in A and B are obtained with whole-cell patch-clamp recordings from the same stellate cell in the presence of picrotoxin (50 μM) and CGP55845 (1 μM) to block inhibitory synaptic transmission. (C) Voltage dependence of stellate cell conductances. Boltzmann fits for activation (black) and inactivation (gray). From left: NaT; activation (V1/2 = –29.9 mV, k = 3.8 mV) and inactivation (V1/2 = –58.2 mV, k = 3 mV) [Magistretti and Alonso (1999a)]; NaP; activation (V1/2 = –44.4 mV, k = 5.2 mV), inactivation (V1/2 = –48.8 mV, k = 19 mV) [Magistretti and Alonso (1999a)]; KA; activation (V1/2 = –34.7 mV, k = 12 mV), inactivation (V1/2 = –80.3 mV, k = 12 mV) where V1/2 parameters were calculated in [Eder and Heinemann (1996)] and the slopes were estimated from Figure 2A [Eder and Heinemann (1996)]; KDR; activation (V1/2 = –38.4 mV, k = 11 mV), inactivation (V1/2 = –43.3 mV, k = 15 mV) where V1/2 given in [Eder and Heinemann (1996)] and the slopes were estimated from Figure 2B [Eder and Heinemann (1996)]; KLeak; activation (V1/2 = –45 mV, k = 8 mV), estimated from threshold and peak amplitude data in [Bruehl and Wadman (1999)]; CaLVA activation (V1/2 = –89 mV, k = 6.9 mV) [Bruehl and Wadman (1999)]; CaHVA activation (V1/2 = –11.1 mV, k = 8.4 mV) [Castelli and Magistretti (2008)]; HCN activation (V1/2 = –83.1 mV, k = 8.1 mV] [Nolan et al. (2007)].
In addition to NaT, stellate cells express a prominent voltage-gated, persistent sodium current (NaP) (Alonso and Llinás, 1989; Jones, 1994; Magistretti and Alonso, 1999a,b, 2002, 2007; Magistretti et al., 1999a,b, 2003). The threshold for activation of NaP is approximately 10–15 mV lower than NaT (Magistretti et al., 1999b) and in contrast to NaT it inactivates relatively slowly with a half-maximal inactivation voltage of approximately −60 mV (Magistretti and Alonso, 1999a,b) (Figure 1C). These properties suggest that in stellate cells NaP is active at membrane potentials below the threshold for action potential firing (Figures 1A,B).

A series of elegant studies have investigated the properties of NaP using single channel recordings (Magistretti and Alonso, 1999a, 2002, 2007; Magistretti et al., 1999a,b, 2003). This approach has the advantages that it is not subject to voltage errors associated with whole-cell recordings, provides information about channel location and gives useful insights into channel gating mechanisms. Single channels with properties of NaP and NaT are found on the dendrites as well as the cell bodies of stellate cells (Magistretti et al., 1999a). NaP is mediated by relatively high conductance single channels (∼20 pS) compared to NaT (∼16 pS) (Magistretti et al., 1999b). The voltage-dependence of activation of the macroscopic NaP current appears to depend on the strong voltage dependence of single channel opening times during bursts of channel activity (Magistretti and Alonso, 2002). Increases in closure duration between such bursts may underlie the slow inactivation of the current (Magistretti and Alonso, 2002). Further detailed analysis suggests that ion channels mediating NaP exist in one of two preferred gating modes (Magistretti et al., 2003) and have at least six different sub-states between fully open and closed (Magistretti and Alonso, 2007). Given the likely importance of stochastic gating of NaP for computation by stellate cells (White et al., 1998; Dorval and White, 2005), it will be important in the future to generate complete kinetic schemes of NaP gating that fully account for these observations.

Potassium channels

Stellate cells express at least four major types of potassium current: leak, delayed rectifier, A-type, and calcium-activated (Figure 1C). They are distinguished by their kinetics, voltage dependence, and pharmacology. An M-type current has also been investigated (Yoshida and Alonso, 2007; Heys et al., 2010), but is yet to be isolated with voltage-clamp experiments.

A substantial potassium leak conductance found in stellate cells is relatively independent of membrane potential and is at least in part blocked by Ba2+, low pH, and quinidine (Deng et al., 2007; Garden et al., 2008), suggesting that it is mediated by two-pore domain (K2P) potassium channels. Consistent with this idea, immunohistochemical data demonstrate that the two-pore domain channels TASK-1 (Kcnk3), TASK-3 (Kcnk9), TWIK-1 (Kcnk1), TREK-1 (Kcnk2), and TREK-2 (Kcnk10) are all present in layer II of MEC (Deng et al., 2009). At present it is not completely clear which of these channels are open at rest, but pharmacological data point to a role for TASK channels (Deng et al., 2007, 2009).

A delayed rectifier potassium current (KDR) activates in stellate cells on depolarization to potentials above approximately

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**FIGURE 2** | Resting membrane properties. (A,C) Examples of perforated patch-clamp recordings of membrane potential responses to current steps (C, lower) recorded from dorsal (A) and ventral (C) stellate cells. Arrows indicate the membrane potential “sag.” (B,D) Peak (closed triangles) and steady-state (open circles) membrane potential responses plotted as a function of current step amplitude for data from dorsal (B) and ventral (D) stellate cells in (A,C).
indirectly following depolarization as a result of Ca\(^{2+}\) from stellate cells has a large relatively fast component and a positive to approximately 45 mV (Eder et al., 1991; Eder and Heinemann, 1996) (Figure 1C). K\(_{DR}\) shows considerable rundown (~80% within 60 min) in the whole-cell recording configuration, with small shifts in steady-state activation and much larger shifts in steady-state inactivation, suggesting that it is under metabolic control. Rundown can be avoided with the use of perforated patch-clamp techniques (Eder et al., 1991). This current has a greater density in stellate cells than pyramidial cells (Eder and Heinemann, 1994). Its molecular basis is unclear.

An A-type current (K\(_{A}\)) recorded from stellate cells is distinguished from K\(_{DR}\) by its strong inactivation (Eder et al., 1991; Eder and Heinemann, 1996) (Figure 1C). Consequently, activation of K\(_{A}\) first requires removal of inactivation by hyperpolarization of the membrane to potentials below ~60 mV. Depolarization to voltages positive to ~50 mV then activates K\(_{A}\). The density of K\(_{A}\) is lower in stellate cells than in neighboring pyramidial cells (Eder and Heinemann, 1994). Its molecular basis is again not clear.

Other voltage-gated potassium channels are much less prominent in stellate cells. M-current blockade has small effects on stellate cell properties (Yoshida and Alonso, 2007; Heys et al., 2010). Nevertheless, there is reasonably high expression of mRNAs for Kv7 subunits in layer II in Allen Brain Atlas data (Garden et al., 2008), although this could be due to expression in non-stellate cells, where M current blockade has a greater effect (Yoshida and Alonso, 2007). There is evidence that stellate cells express inward-rectifying potassium channels. This comes from a small, Ba\(^{2+}\)-sensitive rectification in the instantaneous current-voltage relationship at voltages negative to ~80 mV (Dickson et al., 2000b). However, this current has not been studied further and these findings could also be consistent with properties of some K\(_{2P}\) channels.

Stellate cells also express calcium-activated potassium channels (K\(_{Ca}\)) (Khawaja et al., 2007). These channels are activated indirectly following depolarization as a result of Ca\(^{2+}\) influx through voltage-gated Ca\(^{2+}\) channels. The K\(_{Ca}\) current recorded from stellate cells has a large, fast enhancement component and a smaller slow component (Khawaja et al., 2007). Neither component is sensitive to apamin, which blocks small conductance K\(_{Ca}\) channels, but the slow component can be modulated by the cAMP pathway, as forskolin reduces the amplitude of K\(_{Ca}\) currents (Khawaja et al., 2007). The molecular basis for these currents is unclear. MEC layer II neurons express RNA for apamin-sensitive K\(_{Ca}\)2.2 (SK2) and K\(_{Ca}\)2.3 (SK3) subunits and the apamin-insensitive K\(_{Ca}\)2.1 (SK1) subunit of small-conductance channels (Stocker and Pedarzani, 2000; D’Hoedt et al., 2004). It is either possible that the apamin-sensitive subunits are expressed by non-stellate cells, or that they are expressed by stellate cells, but have functions that have not yet been investigated in stellate cells, for example to control local synaptic depolarizations (Ngo-Anh et al., 2005).

Calcium channels

Two calcium currents have been described in MEC stellate cells (Figure 1C). A high-voltage-activated calcium current (Ca\(_{HV}\)) is mediated by channels that open on depolarization to potentials positive to approximately ~50 mV (Bruehl and Wadman, 1999; Visan et al., 2002; Castelli and Magistretti, 2006), whereas a low-voltage activated current (Ca\(_{LV}\)) is mediated by channels that open with depolarization above approximately ~60 mV (Bruehl and Wadman, 1999; Visan et al., 2002). The two channel populations are further distinguished by the membrane potentials at which they inactivate. Both currents slowly inactivate on depolarization, but Ca\(_{LV}\) inactivates at more depolarized potentials compared with Ca\(_{HV}\) (Bruehl and Wadman, 1999; Visan et al., 2002). This difference is important as it means that depolarization from potentials in the region of ~50 mV will activate Ca\(_{HV}\), but not Ca\(_{LV}\). In contrast, depolarization from more hyperpolarized voltages will activate both currents. Ca\(_{LV}\) currents appear to have similar properties in stellate and pyramidial cells (Bruehl and Wadman, 1999; Castelli and Magistretti, 2006), suggesting that they are not responsible for the distinct electrophysiological properties of stellate cells (Castelli and Magistretti, 2006). There is some evidence that the density of Ca\(_{HV}\) currents is larger in pyramidial cells (Bruehl and Wadman, 1999), but this may be specific to older animals and could be confounded by inclusion of layer III cells in the analysis (Castelli and Magistretti, 2006).

Several calcium channel types appear to contribute to the Ca\(_{HV}\) current. These include L-type channels (~50%), N-type (~23–30%), P/Q-type (~22–24%) and a component that is insensitive to pharmacological block of these channel types (11–13%) (Castelli and Magistretti, 2006). This residual current is reduced by 57–62% by the Ca\(_{2.3}\) channel blocker SNX-482 (Castelli and Magistretti, 2006). Ca\(_{HV}\) has properties consistent with T-type channels, but again the molecular identity of the current remains to be established.

HCN channels

Hypermolarization activated cation currents (I\(_{h}\)) have been extensively studied in stellate cells. Unlike most other voltage-gated currents, I\(_{h}\) activates when the membrane is hyperpolarized and de-activates on depolarization (Robinson and Siegelbaum, 2003). Its presence in stellate cells was initially established from the observation of a sag-like response during injection of negative current steps (Alonso and Llinás, 1989; Jones, 1994). Subsequent voltage-clamp experiments have investigated the kinetics, voltage dependence, and pharmacology of I\(_{h}\) in stellate cells (Richter et al., 1997, 2000; Dickson et al., 2000b; Nolan et al., 2007; Giocomo and Hasselmo, 2008). Interestingly, the density of I\(_{h}\) increases with age from newborn to adult (Richter et al., 1997), with a peak approximately three weeks after birth (Burton et al., 2008), suggesting its importance for function of the mature MEC circuit.

Typically, I\(_{h}\) in stellate cells activates at voltages negative to ~50 mV (Dickson et al., 2000b; Richter et al., 2000; Nolan et al., 2007), shows half-maximal activation at ~80 mV (Dickson et al., 2000b; Nolan et al., 2007) and has a reversal potential of approximately ~20 mV (Dickson et al., 2000b). It is blocked by ZD7288 (Richter et al., 1997; Dickson et al., 2000b; Nolan et al., 2007) and Cs\(^{+}\) (Alonso and Llinás, 1989; Klink and Alonso, 1993; Dickson et al., 2000b; Richter et al., 2000). I\(_{h}\) in stellate cells, but not MEC pyramidial cells, can be activated by the cAMP analog 8-Bromo-cAMP (Richter et al., 2000). The kinetics of I\(_{h}\) are relatively slow compared to transient potassium and sodium currents. The time
course of activation of $I_h$ is best fit with a sum of two exponentials (Dickson et al., 2000b; Nolan et al., 2007). When measured at room temperature to better isolate currents, the fast component has a time constant of approximately $\sim 80$ ms on hyperpolarization to $-90 \text{ mV}$, whereas the slow component has a time constant of approximately $\sim 400$ ms (Dickson et al., 2000b; Nolan et al., 2007). The relative amplitude of the fast component increases with hyperpolarization (Dickson et al., 2000b; Nolan et al., 2007). Deactivation is also relatively slow, with kinetics that speed-up on depolarization (Dickson et al., 2000b; Nolan et al., 2007). While these kinetics are accelerated at physiological temperatures (Nolan et al., 2007), it is important to recognize that they are nevertheless much slower than the voltage-gated channels underlying the action potential. These kinetics are important for certain models of grid cell firing fields discussed further below.

The molecular basis for $I_h$ has been investigated in more detail than that of other ion channels expressed by stellate cells. Four genes (HCN1–4) encode channels with properties of $I_h$ (Robinson and Siegelbaum, 2003). A direct role for HCN1 is evident from experiments using mice in which the HCN1 gene is deleted. This deletion reduces the amplitude of $I_h$ in stellate cells by approximately two-thirds and abolishes the fast activation of the current (Nolan et al., 2007). The residual $I_h$ has slower kinetics and a more negative voltage-dependence than the wild-type current (Nolan et al., 2007). This suggests that the two components of the wild-type current are unlikely to be a simple arithmetic sum of independent currents mediated by HCN1 and other HCN subunits. Rather, a substantial fraction of the wild-type current is likely to be mediated by heteromers containing HCN1 and one or more additional subunits. The kinetics of the wild-type current most closely resemble those of heteromers of HCN1 and HCN2 (Santoro et al., 2000; Chen et al., 2001; Ulen and Tytgat, 2001). Nevertheless, more direct experiments will be required to clearly establish the molecular identity of the additional HCN subunits expressed by stellate cells.

Consistent with these electrophysiological data, gene expression data from the Allen Brain Atlas indicates that mRNA levels of HCN1 and HCN2 are particularly high in layer II of the MEC. Antibody labeling also suggests strong HCN1 expression in superficial layers of the MEC (Notomi and Shigemoto, 2004; Nolan et al., 2007), whilst HCN2 and HCN3 show moderate expression (Notomi and Shigemoto, 2004). However, antibody labeling could reflect HCN1 channels expressed in the dendrites of pyramidal cells with somata in layers III and V (Shah et al., 2004; Rosenkranz, 2006). The use of HCN1 channel knockout mice to establish links between cellular properties of neurons and spatial behavior (Nolan et al., 2004; Giocomo et al., 2011a) is discussed below.

Other ion channels

A non-specific cation current, $I_{\text{NCM}}$, has also been identified as a potentially important contributor to stellate cell function. It is activated during muscarinic receptor-dependent depolarization and is sensitive to Ca$^{2+}$ influx through voltage-gated channels (Klink and Alonso, 1997a; Shalinsky et al., 2002; Magistretti et al., 2004). It possesses a transient tail and sustained plateau component (Magistretti et al., 2004), the former of which may be of functional relevance to action potential firing patterns (Yoshida and Alonso, 2007). Interestingly, this current is analogous to that mediated by TRP channels (Shalinsky et al., 2002) and there is evidence that TRP channels mediate $I_{\text{NCM}}$ in layer V neurons (Zhang et al., 2011). Immunohistochemical analysis shows that the TRPC5 and TRPC1 channels are present in layer II of MEC (Bohlen und Halbach et al., 2010).

INTRINSIC EXCITABLE PROPERTIES OF STELLATE CELLS

Since initial pioneering investigations (Alonso and Llinás, 1989), it has been clear that intrinsic voltage-gated ion channels play an important role in determining the membrane potential trajectory of stellate cells prior to initiation of action potential firing. Ion channels that activate at membrane potentials below the threshold for action potential firing may contribute to the generation of intrinsic activity that is independent of synaptic input (Figures 1A,C), and will also shape integration of synaptic responses as the membrane potential moves toward the action potential threshold (Figures 1B,C). Ion channels that open at more depolarized voltages influence action potential initiation and repolarization (Figures 1A-C). We will consider below key integrative properties of stellate cells at sub-threshold and supra-threshold membrane potentials.

Because there are differences between data obtained from stellate cells with the two techniques most frequently used for investigating their membrane properties—sharp electrode and whole-cell patch-clamp recording (Erchova et al., 2004; Nolan et al., 2007; Boehlen et al., 2010)—we will also introduce new data obtained with the perforated patch-clamp recording method (Figures 2–9). With this technique a high resistance seal is formed between a patch-clamp electrode and the recorded cell. Electrical access to the inside of the cell is then obtained through small pores in the membrane formed by an antibiotic included in the pipette solution (Horn and Marty, 1988; Mathias et al., 1991). Perforated patch-clamp recordings overcome a potential limitation of whole-cell recording by avoiding dialysis of the cell and possible rundown or washout of membrane currents (Horn and Marty, 1988). At the same time, because a tight seal is maintained between the recording electrode and the cell membrane, errors associated with sharp electrode recordings are also avoided (Spruston and Johnston, 1992).

Resting properties

In the absence of significant synaptic input stellate cells do not fire action potentials and have a stable resting membrane potential. At physiological temperatures in brain slices from mature (>30 day) rodents the resting membrane potential of a stellate cell in typically in the range of $-60 \text{ mV}$ to $-70 \text{ mV}$ (Alonso and Llinás, 1989; Alonso and Klink, 1993; Jones, 1994; Erchova et al., 2004; Nolan et al., 2007; Garden et al., 2008; Boehlen et al., 2010). In anesthetized rats the resting potential is approximately $-70 \text{ mV}$ (Quilichini et al., 2010). Input resistance is typically between 20 and 80 MΩ and the membrane time constant is in the range $5–15 \text{ ms}$ (Alonso and Llinás, 1989; Alonso and Klink, 1993; Jones, 1994; Erchova et al., 2004; Nolan et al., 2007; Garden et al., 2008; Boehlen et al., 2010). These properties appear similar in rats and mice (cf. (Nolan et al., 2007; Garden et al., 2008;
Boehlen et al., (2010), change during development (Burton et al., 2008) and are substantially different when measured at room temperature (Dickson et al., 2000b). A large part of the within-study variation in input resistance and time constant can be explained by differences in neuronal location along the dorsal-ventral axis of the MEC (Garden et al., 2008; Boehlen et al., 2010) (see below).

A further source of variation between studies is the recording method used to investigate stellate cell properties. After accounting for a neuron’s position, it appears that sharp electrode recordings produce membrane potentials that are typically more negative, input resistance that is lower and membrane time constants that are shorter compared with data from whole-cell recordings (Boehlen et al., 2010). We find that the direction of the difference between perforated patch-clamp and whole-cell measures of resting potential and input resistance (Table 1) is opposite to that between sharp electrode and whole-cell measures (Boehlen et al., 2010). Therefore, as for hippocampal neurons (Spruston and Johnston, 1992), sharp electrode recordings may under-estimate input resistance and the membrane time constant of stellate cells. Nevertheless, whole-cell recordings also appear to modify the resting membrane potential and input resistance of stellate cells, suggesting the importance of validating key results with perforated patch-clamp methods.

The ionic basis for the stellate cell resting potential involves a balance between a depolarizing drive from leak potassium channels and a hyperpolarizing drive from leak potassium channels. \( I_h \) depolarizes the resting membrane potential by approximately 10 mV (Dickson et al., 2000b; Haas et al., 2007; Nolan et al., 2007) and appears to contribute approximately 50% of the resting membrane conductance, with a large component requiring HCN1 channels (Nolan et al., 2007). Leak potassium channels that are open at rest in stellate cells have properties consistent with the \( \text{Ba}^{2+} \)- and acid-sensitive TASK sub-type of \( K_{\text{v}} \) channel and may account for the majority of the remaining resting membrane conductance (Deng et al., 2007; Garden et al., 2008). It is possible that the resting membrane potential is also influenced by a depolarizing leak \( Na^+ \) conductance (Garden et al., 2008).

Control of the resting potential by leak potassium currents is important for actions of neuromodulators on stellate cells. Serotonin and GABA modulate \( K_{\text{v}} \) channels through 5-HT\(_{1A} \) and GABA\(_B \) receptors, respectively. GABA\(_B \) receptors are highly expressed in stellate cells (Mizukami et al., 2002). Baclofen, a GABA\(_B \) receptor agonist, hyperpolarizes stellate cells and reduces their input resistance by activating TREK-2 channels (Deng et al., 2009). This modulation may be critical in spatial learning (Deng et al., 2009). Serotonin modulation occurs through activation of the quinidine-sensitive channel TWIK-1, which depends on 5-HT\(_{1A} \) receptor activation (Deng et al., 2007). HCN channels may also be targets for modulation (Heys et al., 2010).

### Membrane potential sag

Small perturbations of the membrane potential of a stellate cell are opposed after a short delay, leading to a characteristic “sag” response (Alonso and Llinás, 1989; Alonso and Klink, 1993; Klink and Alonso, 1993, 1997b; Jones, 1994; van der Linden and Lopes da Silva, 1998; Dickson et al., 2000b; Nolan et al., 2007). The sag response is prominent in previous whole-cell and sharp electrode recordings, and using perforated patch-clamp recordings (Figure 2 and Table 1). The sag response is due to a change in the depolarizing inward current provided by \( I_h \) (Dickson et al., 2000b; Nolan et al., 2007). During responses to negative current steps the membrane potential hyperpolarizes and \( I_h \) slowly activates, thus causing a return depolarization. Conversely, when the membrane potential depolarizes \( I_h \) slowly deactivates, causing the membrane potential to then hyperpolarize. \( I_h \) block by \( Cs^+ \) or ZD7288 abolishes the sag (Klink and Alonso, 1993; Jones, 1994; Dickson et al., 2000b; Haas et al., 2007; Nolan et al., 2007), while deletion of HCN1 reduces the amplitude and slows the kinetics of the sag (Giocomo and Hasselmo, 2009; Nolan et al., 2007).

### Theta-frequency resonance

A complementary approach to investigating the integrative properties of stellate cells is through their response to sinusoidally modulated current inputs. Membrane potential responses are largest to inputs that vary at frequencies in the theta range (4–12 Hz) (Haas and White, 2002; Erchova et al., 2004; Schreiber et al., 2004) (Figure 3). This theta-frequency resonance distinguishes stellate cells from other nearby neurons (Erchova et al., 2004). At lower frequencies the membrane potential response also has a small phase advance relative to the injected current, whereas at higher frequencies the membrane potential lags behind the injected current (Erchova et al., 2004; Nolan et al., 2007; Heys et al., 2010) (Figure 3).

Resonant properties also appear to be sensitive to recording method. Two properties are often used to describe resonance: the frequency of the resonance peak (F) and the relative amplitude of the peak (Q) (Hutcheon et al., 1996). Previous studies indicate that F is greater for sharp-electrode compared (Erchova et al., 2004; Boehlen et al., 2010) with whole-cell recordings (Nolan et al., 2007).

| Table 1 | Sub-threshold membrane properties. |
|---|---|---|---|
| | Perforated patch (\( n = 11 \)) | Whole-cell (\( n = 14 \)) | \( p \) | Test |
| Resting membrane potential (mV) | \(-64.2 \pm 0.7 \) | \(-67.1 \pm 0.8 \) | 0.047 | ANCOVA |
| Input resistance (+ve) (\( \Omega \)) | 55.5 ± 9.6 | 34.9 ± 3.5 | 0.0002 | ANCOVA |
| Input resistance (−ve) (\( \Omega \)) | 53.8 ± 9.4 | 32.5 ± 3.0 | 0.0004 | ANCOVA |
| Membrane time constant (ms) | 11.2 ± 0.9 | 12.7 ± 0.9 | 0.64 | ANCOVA |
| Sag coefficient | 0.66 ± 0.02 | 0.58 ± 0.01 | 0.0005 | ANCOVA |
| Location (µm) | 1096 ± 192 | 1114 ± 206 | 0.92 | t-test |
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FIGURE 3 | Membrane potential resonance. (A,D) Examples of perforated patch-clamp recordings of membrane potential responses from resting potential (black) and from depolarized membrane potential (red) to ZAP current waveforms (D, lower), from dorsal (A) and ventral (D) stellate cells. (B) Membrane impedance plotted as a function of frequency. (C) Membrane phase plotted as a function of frequency. (B,C) and (E,F) are calculated from data in (A) and (D), respectively.

et al., 2007; Boehlen et al., 2010; Heys et al., 2010). We find that F is similar, at either resting membrane potentials or peri-threshold potentials (approximately −50 to −55 mV), when measured with whole-cell or perforated patch-clamp methods (Table 2). Q also does not differ significantly between the two methods (Table 2), although with perforated patch-clamp recordings Q is significantly smaller at peri-threshold compared to resting membrane potentials (Table 3). Our finding here that with whole-cell recordings Q is similar when measured at resting potential (−65 mV) and peri-threshold potentials (−50 to −55 mV), while previous results indicate that Q is reduced on depolarization from −70 to −60 mV (Nolan et al., 2007), may be explained by resonance at peri-threshold potentials resulting from a different mechanism to that observed at resting potentials. Together, these data suggest that sharp electrode recordings may over-estimate the F at both resting and peri-threshold potentials, while whole-cell methods appear to provide estimates of resonant properties that are comparable to perforated patch-clamp data.

At resting membrane potentials the attenuation of low frequency signals that is essential for resonance in stellate cells requires HCN1 channels (Nolan et al., 2007). Deletion of HCN1 or pharmacological block of Ih abolishes resonance by increasing the amplitude of responses to current inputs with frequencies less than 4 Hz (Haas et al., 2007; Nolan et al., 2007). The phase advance is also abolished (Nolan et al., 2007). Thus, resonance at resting membrane potentials can be explained by Ih opposing slow changes in the membrane potential (Hutcheon and Yarom, 2000; Nolan et al., 2007; Dudman and Nolan, 2009; Dodson et al., 2011). Roles for other ion channels in resonance are less clear. In other cell types persistent sodium currents and M-type potassium currents can contribute to resonance at more depolarized voltages (Hu et al., 2007). In stellate cells the M-current blocker XE991 reduces the frequency and amplitude of the resonance peak at more depolarized voltages, but not at resting membrane potentials, although the frequency shift is on average less than 1 Hz and the reduction in amplitude appears quite small (Heys et al., 2010).
Theta-frequency resonance is also observed when stellate cells are probed with non-periodic frozen-noise inputs (Schreiber et al., 2004). However, the relationship between resonance and stellate cell responses to more physiological inputs is not fully understood. For excitatory glutamatergic synaptic input, stimulation at theta frequency produces little or no temporal summation of the response, whereas stimulation at gamma frequencies (40–80 Hz) causes strong temporal summation (Garden et al., 2008). This frequency dependence may reflect the fact that glutamatergic synaptic currents are very short compared with the membrane time constant (Garden et al., 2008) and so fast synaptic inputs will have very different temporal structure to the continuous sinusoidal inputs that are typically used to investigate resonant properties.

**Peri-threshold theta-frequency membrane potential activity**

When the membrane potential of a stellate cell is depolarized from rest to close to the threshold for action potential firing it becomes unstable, appearing to oscillate with frequency in the theta range (4–12 Hz) (Figure 1A) (Alonso and Llinás, 1989; Alonso and Klink, 1993; Dickson et al., 2000b; Erchova et al., 2004; Nolan et al., 2007; Dodson et al., 2011; Yoshida et al., 2011). This is of considerable interest first because it suggests a cellular correlate of the network theta rhythm recorded from the MEC and other hippocampal regions in vivo (Stewart et al., 1992; Dickson et al., 2000a; Mizuseki et al., 2009), and second because it has been argued that theta-frequency oscillations may be a substrate for encoding of spatial information by grid cells (Burgess et al., 2007; Giocomo et al., 2007).

Peri-threshold theta-frequency activity is consistently recorded from stellate cells with both sharp electrode and whole-cell techniques. However, the frequency of membrane potential activity during whole-cell recordings (Giocomo et al., 2007; Nolan et al., 2007; Giocomo and Hasselmo, 2009; Boehlen et al., 2010; Dodson et al., 2011) appears to be lower than during sharp electrode recordings (Alonso and Llinás, 1989; Klink and Alonso, 1993; White et al., 1998; Erchova et al., 2004; Boehlen et al., 2010). We find with perforated patch-clamp recordings that all stellate cells generate significant theta-frequency activity at membrane potentials close to spike threshold (Figure 4). In spectrograms the activity is characterized by brief epochs of power at theta frequencies (Figures 4B,G). Use of long windows for spectrograms as in some previous studies (Yoshida et al., 2011) may cause transient activity to appear sustained, whereas with shorter windows spectral peaks are revealed as non-stationary (cf Figures 4B–C, G–H). When activity is analyzed with Lomb periodograms of adjacent segments of activity (see Dodson et al., 2011 for details), the instability is manifest as variability in the frequency of the most significant activity peak and as multiple peaks in each segment (Figures 4D,J). This variability is clear from representative data (Figures 4E,J). Variability in the frequency of activity is also apparent during periods between action potential clusters (Figure 4K). Considering only the most significant peak in each of five epochs of duration 3 s, the

### Table 2 | Resonant properties.

<table>
<thead>
<tr>
<th></th>
<th>Perforated patch (n = 5)</th>
<th>Whole-cell (n = 14)</th>
<th>p</th>
<th>Test</th>
</tr>
</thead>
<tbody>
<tr>
<td>Q (resting)</td>
<td>2.0 ± 0.2</td>
<td>1.7 ± 0.1</td>
<td>0.1</td>
<td>ANCOVA</td>
</tr>
<tr>
<td>Q (peri-threshold)</td>
<td>1.6 ± 0.1</td>
<td>1.8 ± 0.1</td>
<td>0.3</td>
<td>ANCOVA</td>
</tr>
<tr>
<td>$F_{\text{Z, max}}$ (Hz) (resting)</td>
<td>5.9 ± 1.0</td>
<td>6.1 ± 0.5</td>
<td>0.84</td>
<td>ANCOVA</td>
</tr>
<tr>
<td>$F_{\text{Z, max}}$ (Hz) (peri-threshold)</td>
<td>4.7 ± 0.9</td>
<td>4.4 ± 0.4</td>
<td>0.56</td>
<td>ANCOVA</td>
</tr>
</tbody>
</table>

<table>
<thead>
<tr>
<th></th>
<th>Resting</th>
<th>Peri-threshold</th>
<th>p</th>
<th>Test</th>
</tr>
</thead>
<tbody>
<tr>
<td>Q (perforated patch)</td>
<td>2.0 ± 0.2</td>
<td>1.6 ± 0.1</td>
<td>0.03</td>
<td>Paired t-test</td>
</tr>
<tr>
<td>Q (whole-cell)</td>
<td>1.7 ± 0.1</td>
<td>1.8 ± 0.1</td>
<td>0.4</td>
<td>Paired t-test</td>
</tr>
<tr>
<td>$F_{\text{Z, max}}$ (Hz) (perforated patch)</td>
<td>5.9 ± 1.0</td>
<td>4.7 ± 0.9</td>
<td>0.02</td>
<td>Paired t-test</td>
</tr>
<tr>
<td>$F_{\text{Z, max}}$ (Hz) (whole-cell)</td>
<td>6.1 ± 0.5</td>
<td>4.4 ± 0.4</td>
<td>2.8e-6</td>
<td>Paired t-test</td>
</tr>
</tbody>
</table>

### Table 3 | Action potential properties.

<table>
<thead>
<tr>
<th></th>
<th>Perforated patch (n = 11)</th>
<th>Whole-cell (n = 13)</th>
<th>p</th>
<th>Test</th>
</tr>
</thead>
<tbody>
<tr>
<td>Proportion spikes in clusters (Pc)</td>
<td>0.85 ± 0.06</td>
<td>0.26 ± 0.1</td>
<td>1.1e-5</td>
<td>ANCOVA</td>
</tr>
<tr>
<td>Intra-cluster spike frequency (Hz)</td>
<td>11.8 ± 1.3</td>
<td>7.8 ± 0.3</td>
<td>0.004</td>
<td>ANCOVA</td>
</tr>
<tr>
<td>Overall spiking frequency (Hz)</td>
<td>2.1 ± 0.1</td>
<td>2.3 ± 0.1</td>
<td>0.17</td>
<td>ANCOVA</td>
</tr>
<tr>
<td>Spikes per cluster</td>
<td>2.6 ± 0.1</td>
<td>3.8 ± 0.7</td>
<td>0.1</td>
<td>ANCOVA</td>
</tr>
<tr>
<td>Rheobase (mA)</td>
<td>161.8 ± 18.8</td>
<td>246.2 ± 21.6</td>
<td>0.005</td>
<td>ANCOVA</td>
</tr>
<tr>
<td>Spike max (mV)</td>
<td>41.9 ± 1.7</td>
<td>45.8 ± 1.0</td>
<td>0.06</td>
<td>ANCOVA</td>
</tr>
<tr>
<td>Spike width at half-height (ms)</td>
<td>0.59 ± 0.02</td>
<td>0.6 ± 0.01</td>
<td>0.65</td>
<td>ANCOVA</td>
</tr>
<tr>
<td>Spike threshold (mA)</td>
<td>−44.7 ± 1.45</td>
<td>−42.5 ± 0.7</td>
<td>0.16</td>
<td>ANCOVA</td>
</tr>
<tr>
<td>AHP minimum (mV)</td>
<td>−60.9 ± 0.7</td>
<td>−59.2 ± 0.6</td>
<td>0.08</td>
<td>ANCOVA</td>
</tr>
<tr>
<td>AHP width at half-height (ms)</td>
<td>64.3 ± 3.9</td>
<td>74.9 ± 4.1</td>
<td>0.03</td>
<td>ANCOVA</td>
</tr>
</tbody>
</table>
mean frequency is $3.6 \pm 0.5$ Hz, while the range of significant frequencies across all five epochs is $3.2 \pm 0.55$ Hz. These results are similar to data from previous whole-cell recordings (Dodson et al., 2011).

The mechanism responsible for peri-threshold theta-frequency activity has been subject to debate. One class of ionic mechanism proposes that theta-frequency activity is generated by a periodic oscillator. This follows from the consideration that addition of an amplifying conductance, such as NaP, to a resonator can lead to the generation of oscillations (Hutcheon and Yarom, 2000). Consistent with this idea, models of stellate cells that include NaP and Ih can produce periodic membrane potential oscillations at peri-threshold potentials (Dickson et al., 2000b; Fransén et al., 2004). An important test of this oscillator model is whether the theta activity is periodic. In this case power spectra of the activity will have a single spectral peak that will have a stable frequency in consecutive analysis windows (Hajimiri and Lee, 1998; Dodson et al., 2011). The theta-frequency activity, recorded with perforated patch-clamp methods (Figure 4), and in previous whole-cell recordings (Nolan et al., 2007; Dodson et al., 2011), does not match this prediction (Figures 4B–D, G–I). Instead, the theta activity during sub-threshold sweeps and between spikes during suprathreshold sweeps (Figure 4K) has a frequency that changes apparently unpredictably. These observations are evidence against some models of grid cell firing fields that we discuss below. While these properties could perhaps be explained by variation in the period of an oscillator (Zilli et al., 2009), the variation would be so high that this would amount to an oscillatory system with dynamics that are essentially swamped by stochastic fluctuations.

A second class of ionic mechanism that can account for peri-threshold theta-frequency activity is based on evidence that stochastic gating of membrane ion channels can cause the membrane potential to fluctuate at theta frequencies (White et al., 1998). Direct support for this mechanism comes from experiments in which theta-frequency activity is abolished by pharmacologically blocking NaP. When a dynamic clamp is then used to reintroduce a purely deterministic version of NaP theta-frequency fluctuations remain absent (Dorval and White, 2005). In contrast, reintroduction of stochastic NaP channels restores theta-frequency activity (Dorval and White, 2005). Consistent with this mechanism, detailed models of stellate cells, in which all ion channels gate stochastically, reproduce spectral properties of theta-frequency activity (Dudman and Nolan, 2009; Dodson et al., 2011). Nevertheless, because these models represent stellate cells as a single compartment, whereas the electrical properties of stellate cells are primarily determined by their dendrites (Garden et al., 2008), it will be important to establish whether theta-frequency activity is also accounted for by models that include stochastic gating of dendritic ion channels (Cannon et al., 2010).

Considering these data together, it is unlikely that theta-frequency membrane potential fluctuations reflect the output of a periodic oscillatory process. Instead, they appear most

![FIGURE 4](http://www.frontiersin.org) Peri-threshold theta-frequency activity. (A,F) Examples of perforated patch-clamp recordings of peri-threshold (black) and threshold (gray) membrane potential responses to current steps. (B,G) Spectrograms from peri-threshold responses in A,F, calculated with 1 s windows. (C,H) Spectrograms from the same traces calculated using 6.5 s windows. (D,I) Frequency of most significant component (red) and all other significant frequencies (black) of peri-threshold membrane potential activity in A,F. Frequencies are obtained from Lomb periodograms of contiguous 3 s segments of data. (E,J) Examples of consecutive 2 s segments of data from peri-threshold activity in A,F (K) Examples of 1.5 s segments of membrane potential activity triggered by final spikes of three consecutive clusters of action potentials in A. Note the different frequency of the theta activity following each spike cluster.
consistent with filtering of membrane potential noise resulting from stochastic ion channel gating. In this scenario, resonance due to voltage-gated ion channels likely plays an important role in shaping the frequency of the noise-driven membrane potential activity (White et al., 1998; Haas et al., 2007). It remains to be determined whether this intrinsic theta-frequency activity originates in the cell body or dendrites of stellate cells. For example, membrane potential changes at the cell body may reflect much larger stochastic events originating in the dendrites (Cannon et al., 2010).

The action potential and its after-polarization
Stellate cells generate action potentials when their membrane potential is depolarized above approximately −50 mV (Alonso and Klink, 1993; Jones, 1994). When action potentials are initiated during injection of constant positive current they are followed immediately by an after-polarization with several components (Figure 5). First, there is a shallow and rapid AHP, which is followed by a brief after-depolarization (ADP). Finally, at typical physiological spiking rates (<∼20 Hz) (Sargolini et al., 2006), there is a slower AHP which maintains the membrane potential below its peri-threshold value (Alonso and Klink, 1993). These properties are seen with whole-cell and sharp electrode recordings, but their relative prominence may differ (Boehlen et al., 2010). We find that with perforated patch-clamp recordings these characteristic features of the action potential and its AHP are also maintained (Figure 5). We note that when action potentials are triggered by brief synaptic input, the waveform of the after-polarization differs (Figure 1B), but while this difference may be important for physiological activation of stellate cells it has so far received very little attention.

Initiation of the action potential is blocked by TTX (Alonso and Llinás, 1989), but not by riluzole (Dorval and White, 2005), indicating that it requires NaT, but not NaP. As for many other neuron types, the duration of the action potential is increased by blocking voltage-gated potassium channels (Klink and Alonso, 1993). This is consistent with KDR channels mediating rapid repolarization. Ca2+-dependent activation of KCa is important for the later stages of repolarization and for the AHP (Klink and Alonso, 1993; Khawaja et al., 2007). The recovery of the membrane potential to rest following the peak of the AHP is accelerated by activation of Ih mediated by HCN1 channels (Nolan et al., 2007). Further experimental investigation of the action potential AHP is likely to be important, particularly as the AHP and ADP are difficult to fully account for in current biophysical models of stellate cells (Fransén et al., 2004; Dudman and Nolan, 2009) and may play important roles in generation of grid cell firing fields (Navratilova et al., 2012).

Clustering of action potentials
Compared with other neuron types, stellate cells generate distinctive clustered patterns of action potentials during maintained supra-threshold depolarization (Alonso and Klink, 1993; Klink and Alonso, 1993; Nolan et al., 2007; Engel et al., 2008; Fernandez and White, 2008) (Figures 1A, 5). During whole-cell recordings, spikes within a cluster have similar inter-spike intervals that are independent of the overall firing frequency and that are often in the high theta range (about 8–12 Hz) (Nolan et al., 2007; Fernandez and White, 2008). Clustered patterns of spikes are most robust when the overall spike frequency is less than 5 Hz, with the inter-cluster interval sometimes being a second or longer (Nolan et al., 2007). At higher overall spike frequencies clustered patterns of spikes are no longer detectable. Spike trains show very little frequency adaptation at firing rates below the typical physiological maximum (∼20 Hz) (Alonso and Klink, 1993; Sargolini et al., 2006). Importantly, while clustered patterns of action potential firing have peaks in their power spectra in the theta-frequency range, the underlying mechanism is distinct from...
that of sub-threshold theta-frequency resonance (Nolan et al., 2007; Fernandez and White, 2008).

When recorded with perforated-patch methods clustered firing is exceptionally robust (Figures 5A,C). The probability that a spike is within a cluster is substantially greater when recorded with perforated patch compared to whole-cell methods (Table 3). The frequency of spikes within clusters is also higher for perforated patch recordings, whereas the number of spikes per cluster is similar to measurements with whole-cell recordings (Nolan et al., 2007). These data suggest that during whole-cell recordings there may be washout of a conductance that is not necessary for clustered firing patterns, but that increases their probability of occurrence.

We took advantage of the relatively large differences between whole-cell and perforated patch-clamp methods in the duration of the AHP and the probability that a spike is part of a cluster to estimate the time course of changes that takes place during intracellular dialysis associated with whole-cell recording from stellate cells (Figure 6). We found that even within 2 min of break in to the whole-cell configuration—the shortest interval within which we could reliably estimate the threshold current to trigger spike firing—the probability that spikes are part of a cluster differed significantly from that measured with perforated patch-clamp recording (Figure 6). This property and the AHP half-duration continued to change during the first 30 min following break-in (Figure 6). These data suggest that signaling pathways sensitive to washout during whole-cell recording regulate the firing properties of stellate cells.

Ion channels that determine the amplitude and duration of the AHP appear to be critical for determining the pattern of spiking activity generated by stellate cells (Fransén et al., 2004; Nolan et al., 2007; Fernandez and White, 2008; Dudman and Nolan, 2009). Deletion of HCN1 or block of $I_h$ greatly reduces spike clustering (Nolan et al., 2007). This can be explained by a model in which action potentials are triggered by stochastic fluctuations in membrane potential (Dudman and Nolan, 2009). In this model HCN channels engaged during the AHP provide a depolarizing drive that transiently increases the probability that membrane potential noise will trigger an action potential (Dudman and Nolan, 2009). Alternatively, termination of spike clusters could be achieved through incremental elevation of intracellular Ca$^{2+}$ following each action potential (Fransén et al., 2004). However, with this mechanism it is difficult to explain the probabilistic properties of clustered spike firing that are accounted for by stochastic models (Dudman and Nolan, 2009). A model based on inductor-dependent resonance can also reproduce statistics of clustered spiking trains (Engel et al., 2008). However, this model appears inconsistent with experiments showing that the dominant frequency of spike train power spectral densities can be manipulated independently from that of peri-threshold membrane potential activity (Fernandez and White, 2008).

**SYNAPTIC INTEGRATION**

A typical stellate cell has 6–10 primary dendrites that are densely populated with dendritic spines (Klink and Alonso, 1997b; Buckmaster et al., 2004; Garden et al., 2008). These spines are likely to be the main sites of excitatory synaptic input to stellate cells, but while the experiments considered so far evaluate stellate cell output in response to current injected through an experimental electrode, much less are known about how stellate cells respond to synaptic inputs. Stimulation of layer I evokes excitatory postsynaptic potentials (EPSPs) in stellate cells that can trigger action potentials (Figure 1B) (Alonso et al., 1990; Garden et al., 2008). The origin of these inputs is not clear, but...
might include axons from neurons in the neocortex (Dolorfo and Amaral, 1998; Burwell, 2000). Synaptic potentials can also be evoked in stellate cells by stimulation in the parasubiculum or deeper layers of the MEC (Jones, 1994). Recent experiments using localized photo-uncaging of glutamate to activate inputs from neurons at anatomically defined locations have begun to establish principles for the organization of local microcircuits containing stellate cells (Kumar et al., 2007; Beed et al., 2010). These experiments suggest that stellate cells receive local excitatory input from neurons in layer II of the MEC as well as from deeper layers. However, dual intracellular recordings from over 100 pairs of layer II principle cells did not find evidence for excitatory connections between stellate cells, suggesting that if recurrent excitation is present within layer II it is very sparse (Dhillon and Jones, 2000). The neurotransmitter receptors that mediate synaptic responses of stellate cells appear to have properties similar to other central neurons. EPSPs have fast components mediated by AMPA (non-NMDA) receptors and slower NMDA receptor responses (Garden et al., 2008). Both channel types determine the time window for detection of coincident inputs and summation of inputs activated at gamma frequencies (Garden et al., 2008). If either channel is blocked the coincidence detection time window becomes wider and summation of gamma frequency inputs increases. In contrast, the interval between synaptic responses evoked at theta frequency is too long for interactions between successive responses to be influenced by these integrative mechanisms (Garden et al., 2008). Unlike for some other neuron types (Stuart and Sakmann, 1994; Schiller et al., 2000), it is not yet clear whether synaptic input to dendrites of stellate cells can initiate local spikes. The role of N_{AP} is also unclear, although experiments in other neuron types suggest it may be important for amplification of EPSPs (Stuart and Sakmann, 1995).

DORSAL-VENTRAL ORGANIZATION OF MEMBRANE PROPERTIES

Grid cells are organized topographically according to the spatial resolution of their firing fields, such that more dorsal cells have smaller firing fields that are spaced closer together compared with more ventral cells (Hafting et al., 2005; Sargolini et al., 2006; Barry et al., 2007; Brun et al., 2008; Fyhn et al., 2008). Several of the integrative properties of stellate cells that we describe above follow a similar dorsal-ventral organization. This correlation is of potential importance as it may reflect cellular computations important for generation of grid fields (Giocomo et al., 2011b; O’Donnell and Nolan, 2011). However, to test whether such correlations play a causal role in establishing grid cell firing fields, it is essential to first understand their underlying mechanisms. We describe below evidence for dorsal-ventral organization of ion channel properties, resting membrane properties, frequency selectivity, action potential clustering, and synaptic integration. We will suggest how these observations can be explained by increasing density of currents mediated by HCN and leak potassium channels in stellate cells located at dorsal compared to more ventral locations (Garden et al., 2008).

Ion channel function

There is evidence for a dorsal-ventral gradient in the density of a leak potassium conductance (Garden et al., 2008) and in either the density (Garden et al., 2008) or kinetics of I_{h} (Giocomo and Hasselmo, 2008). In one study, when pharmacologically isolated I_{h} was measured by two independent methods—from tail currents and from currents sensitive to the blocker ZD7288—the amplitude of I_{h} appeared to be smaller for neurons located at progressively more ventral locations (Garden et al., 2008). Differences in morphology between cells at different locations were insufficient to explain the differences in the measured currents (Garden et al., 2008). These data are consistent with in situ hybridization data in the Allen Brain Atlas, which suggests that mRNA levels of HCN1 and HCN2 follow a dorsal-ventral gradient in expression in layer II of the MEC. In this study there was very little variability in the time constant of I_{h} (Garden et al., 2008). In contrast to these data, a different study suggests that time constants, but not the amplitude of I_{h}, vary across the dorsal-ventral extent of the MEC (Giocomo and Hasselmo, 2008). Because these differences lead to very different functional predictions for the roles of I_{h} in stellate cells, it will be important for future work to clarify their cause, for example young (Giocomo and Hasselmo, 2008) vs. mature animals (Garden et al., 2008), or rats (Giocomo and Hasselmo, 2008) vs. mice (Garden et al., 2008).

Resting membrane properties

Stellate cells at dorsal locations have lower input resistance and faster membrane time constants than those at ventral locations (Garden et al., 2008; Boehlen et al., 2010). Consistent with these data, in perforated patch-clamp experiments we also find a significant dorsal-ventral organization of the input-resistance and membrane time constant, but not the resting membrane potential (Figures 7A–C). This organization can be accounted for by differences in the density of both leak potassium channels and HCN channels that mediate I_{h} (Garden et al., 2008). Thus, in dorsal neurons a greater density of leak potassium and HCN conductances lowers the input resistance and membrane time constant compared with more ventral neurons.

Frequency selectivity

Stellate cells exhibit a gradient in their frequency selectivity, with dorsal cells having a resting resonance peak at higher frequencies than cells from more ventral locations (Giocomo et al., 2007; Boehlen et al., 2010). Results from perforated patch-clamp experiments are consistent with these data (Figures 7D–E). Differences in the density of currents through HCN channels could account for this organization of resonance frequencies, as greater I_{h} amplitude in dorsal cells will more effectively oppose slow changes.
FIGURE 7 | Dorsal-ventral organization of resting and resonant membrane properties. (A–E) Resting membrane potential (A), input resistance (B), membrane time constant (C), resonance amplitude (D) and resonant frequency (E) are plotted as a function of the location of the recorded stellate cell. Red lines indicate fits to the perforated patch-clamp data. Grey dashed lines indicate fits to the whole-cell data. Adjusted $R^2$ value and significance of the fit for the perforated patch-clamp data are stated above each plot. The $R^2$ and significance values for fits of the whole-cell data are as follows: $V$, $R^2 = 0.04$, $p = 0.24$; $R$, $R^2 = 0.59$, $p = 0.0009$; $\tau_m$, $R^2 = 0.64$, $p = 0.0003$; $Q_{res}$, $R^2 = 0.22$, $p = 0.052$; $Q_{peri}$, $R^2 = 0.006$, $p = 0.32$; $F_{rest}$, $R^2 = 0.48$, $p = 0.005$; $F_{peri}$, $R^2 = 0.5$, $p = 0.003$. Location refers to distance from the dorsal border of the MEC.

in membrane potential, thus causing greater attenuation of low frequency inputs (Nolan et al., 2007; Heys et al., 2010). While not found in mature neurons (Garden et al., 2008), differences in the kinetics of $I_h$ could in principle also account for the dorsal-ventral organization of membrane resonance (Giocomo and Hasselmo, 2008), but this would not explain the dorsal-ventral organization of input resistance (Garden et al., 2008).

Just as for the resonant response to injected current, the properties of peri-threshold theta-frequency activity also follow a dorsal-ventral gradient (Giocomo et al., 2007; Giocomo and Hasselmo, 2008, 2009; Dodson et al., 2011; Yoshida et al., 2011). With perforated patch-clamp recordings the frequency of the most significant peak in a 15 s window varies steeply as a function of location (Figures 8A–C). These data are consistent with previous observations from whole-cell and sharp electrode recordings (Giocomo et al., 2007; Giocomo and Hasselmo, 2008, 2009; Dodson et al., 2011; Yoshida et al., 2011). The dependence of the significant frequencies on the length of the analysis time window (Figures 8A–B) and their variability between consecutive time windows (Figure 8C) is consistent with the idea that theta-frequency activity reflects filtering of a noise source (Dodson et al., 2011).

The ionic basis for the dorsal-ventral organization of theta-frequency activity appears to involve differences in $I_h$. Knockout of the HCN1 subunit flattens the frequency gradient of the fluctuations (Giocomo and Hasselmo, 2009), although following pharmacological block of $I_h$ a gradient nevertheless remains (Dodson et al., 2011). If the theta-frequency activity was periodic then these differences could be explained by models in which ion channels gate deterministically and the kinetics of $I_h$ vary with location (Yoshida et al., 2011). However, as we discuss above, models of this kind do not appear to account for the stochastic nature of theta-frequency activity. In contrast, both the variability and the frequency of the largest amplitude activity are accounted for by models in which all ion channels gate stochastically and in which the density of HCN and leak potassium channels follows a dorsal-ventral organization (Dodson et al., 2011). It is, therefore, possible that the dorsal-ventral organization of intrinsic theta-frequency activity is a secondary consequence of control of the resting integrative properties of stellate cells by these ion channels.

Clustering of action potentials

Stellate cells also demonstrate a dorsal-ventral gradient in their pattern of action potential firing (Figures 9A–C). This organization of stellate cell firing patterns has not previously been described. With perforated patch-clamp recordings we find that, although spike clustering remains high in stellate neurons along
FIGURE 8 | Dorsal-ventral organization of theta-frequency activity. (A) The frequency of the most significant peak in a Lomb periodogram of 15 s of peri-threshold membrane potential activity is plotted as a function of the location of the recorded neuron. (B) The mean frequency (left) and the range of frequencies (right) of the most significant peak of Lomb periodograms, obtained from five consecutive 3 s duration segments of peri-threshold activity, is plotted as a function of the location of the recorded neuron. (C) The number of significant peaks (left), the mean frequency (middle) of all significant peaks, and the range of frequencies of all significant peaks (right), obtained from five consecutive 3 s duration segments of peri-threshold activity, is plotted as a function of the location of the recorded neuron. Data analyzed were from 5 s to 20 s after the onset of the largest amplitude current step that did not trigger action potential firing. Adjusted $R^2$ value and significance of the fit for the perforated patch-clamp data are stated above each plot. Location refers to distance from the dorsal border of the MEC.

FIGURE 9 | Dorsal-ventral organization of firing properties. (A–E) Probability that a spike is part of a cluster (A), frequency of spikes throughout the 15 s duration analysis window (open circles) or within a cluster (closed circles) (B), number of spikes per cluster (C), duration of the action potential after-hyperpolarization (D), and rheobase (E), are plotted as a function of the location of the recorded neuron. Data analyzed were from 5 s to 20 s after the onset of the current steps that triggered action potential firing at frequencies in the range 1–3 Hz. Adjusted $R^2$ value and significance of the fit for the perforated patch-clamp data are stated above each plot. The $R^2$ and significance values for fits of the whole-cell data are as follows: $F_{\text{cluster}}, R^2 = 0.35, p = 0.02; F_{\text{intracluster}}, R^2 = 0.56, p = 0.03; F_{\text{intracluster}}, R^2 = 0.19, p = 0.08; \text{Spikes/cluster}, R^2 = 0.38, p = 0.08; \text{AHP}, R^2 = 0.19, p = 0.08; \text{Rheobase}, R^2 = 0.61, p = 0.0009. Location refers to distance from the dorsal border of the MEC.

the full dorsal-ventral extent of the MEC (Figure 9A), the frequency with which spikes occur within clusters follows a gradient (Figure 9B). In contrast, in our whole-cell recordings clustering appears to be reduced in more ventral cells, leading to the emergence of a gradient in the probability that an action potential is part of a clustered firing pattern (Figure 9A). The number of spikes per cluster is independent of location and recording method.
This organization of the pattern of spike firing by stellate cells can be explained by differences in the density of current through HCN channels (Garden et al., 2008) and their influence on the action potential AHP (Nolan et al., 2007; Dudman and Nolan, 2009). Because activation of HCN channels accelerates repolarization following the peak of the action potential AHP (Nolan et al., 2007; Dudman and Nolan, 2009), the greater density of HCN channels in dorsal cells will reduce the duration of the AHP and, therefore, increase the intra-cluster spike frequency. Consistent with this explanation, the AHP duration is shortest in more dorsal cells and increases for cells at progressively more ventral locations (Figure 9D) (Boehlen et al., 2010). The dependence of the probability of clustered firing on location found with whole-cell recordings is consistent with a strong correlation between these parameters in whole-cell recording conditions (Nolan et al., 2007; Fernandez and White, 2008). The absence of this gradient from perforated-patch recordings suggests wash-out during whole-cell recordings of a factor that promotes clustered patterns of activity in ventral cells. One possible explanation is that wash-out during whole-cell recordings causes a shift in the voltage dependence of HCN channels that reduces their influence on the AHP.

The threshold-current required to initiate action potential firing also follows a dorsal-ventral organization (Garden et al., 2008; Boehlen et al., 2010). This organization is maintained in perforated patch-clamp and whole-cell recordings (Figure 9E). The organization of the threshold-current is at least in part explained by the lower input resistance of dorsal compared to more ventral neurons, resulting from their higher density of leak potassium and HCN channels (Garden et al., 2008).

**Synaptic integration**

The waveforms of evoked and spontaneous excitatory synaptic potentials recorded from stellate cells follow a dorsal-ventral organization that is also explained by the density of currents through HCN and leak potassium channels. For neurons located more dorsally, EPSPs are shorter than for neurons located more ventrally (Garden et al., 2008). One consequence is that gamma frequency inputs to stellate cells at ventral locations sum more effectively relative to cells at more dorsal locations (Garden et al., 2008). Another is that dorsal cells have narrower time windows for detection of coincident synaptic activity than ventral cells and may, therefore, be better at temporal discrimination (Garden et al., 2008). Importantly, either leak potassium channels or HCN channels alone can support the gradient in integrative properties, although block of either channel increases the duration of EPSPs (Garden et al., 2008). Simulations suggest that gradients in the amplitude of the currents that organize synaptic integration are sufficient to also account for the dorsal-ventral organization of peri-threshold theta-frequency activity (Dodson et al., 2011).

**FROM ION CHANNELS TO GRID FIRING FIELDS**

The electrophysiological investigations of stellate cells described above provide a foundation to begin addressing questions about cellular mechanisms for spatial computations carried out within the MEC. For example, what is the relationship between computational properties evaluated with in vitro experiments and generation of grid firing fields? What are the roles in generation of grid firing fields of particular ion channels expressed by stellate cells? Does modulation of the intrinsic electrophysiological properties of stellate cells play roles in firing during spatial behaviors? Answering these and related questions will require integration of cellular data with predictive models for computation carried out during spatial behaviors, and testing of these models using recordings from behaving animals and specific manipulation of cellular properties of stellate cells.

**Constraining models based on intrinsic electrophysiological properties**

Two general classes of abstract model have been proposed to account for generation of grid firing fields (Burgess and O’Keefe, 2011; Giocomo et al., 2011b). In one, grid fields are produced by interference between oscillations with frequency sensitive to an animal’s velocity (Burgess et al., 2007; Giocomo et al., 2007; Blair et al., 2008; Burgess, 2008; Hasselmo, 2008). In the second, grid fields result from bumps in an attractor network that are shifted by velocity inputs (Fuhs and Touretzky, 2006; McNaughton et al., 2006; Burak and Fiete, 2009; Navratilova et al., 2012). We consider below how known intrinsic electrophysiological properties of stellate cells constrain the ways in which each class of abstract model can be implemented, and how knowledge of ion channel function in stellate cells can be used to distinguish between different models.

Models that compute location through oscillatory interference rely on periodic oscillators that are sensitive to velocity. Initial versions of these models proposed that theta-frequency activity of stellate cells reflected one or more oscillations of this kind (Burgess et al., 2007; Giocomo et al., 2007). However, several cellular properties argue against this mechanism. First, while theta-frequency activity can be interpreted as the output of a periodic oscillator (Dickson et al., 2000a; Fransén et al., 2004), the data discussed above suggests it is more likely to reflect filtered stochastic noise (White et al., 1998; Nolan et al., 2007; Dodson et al., 2011). This is a substantial problem as filtered noise signals are not periodic and computation by oscillatory interference mechanisms is extremely sensitive to noise (Burak and Fiete, 2009). Second, models that require independent oscillators in different dendrites appear to be impossible to implement given known cable properties of dendrites (Remme et al., 2010). Third, theta-frequency activity is abolished by synaptic conductances that mimic activity that might occur in vivo (Fernandez and White, 2008).

Alternative implementations of oscillatory interference models assume that neurons that act as velocity-sensitive oscillators are located upstream of grid cells (Blair et al., 2008; Burgess, 2008; Hasselmo, 2008). Recent evidence suggests that firing of “theta cells” in the anterior thalamus, hippocampus and medial septum may perform this function (Welday et al., 2011). In these cells theta-frequency bursts of action potentials are modulated by both movement speed and movement direction. For further evaluation of this theory it will be important to determine if these cells synapse with stellate cells and how the intrinsic properties of stellate cells influence responses to these inputs. For example, dorsal-ventral gradients in theta-frequency resonance or synaptic
integration might contribute to decoding of inputs from velocity sensitive oscillators.

In models of grid cell firing that rely on network attractor states the influence of stellate cell integrative properties has received less direct attention. While these models are typically implemented using abstract neurons, in all cases tuning of neuronal gain is necessary for the network to produce attractor states (Fuhs and Touretzky, 2006; McNaughton et al., 2006; Burak and Fiete, 2009; Navratilova et al., 2012). Therefore models of this kind can also be used to make predictions for how intrinsic properties of stellate cells might influence grid firing fields. A recent attractor model predicts possible roles for the AHP and ADP, along with synaptic NMDA conductances, in grid cell firing, periodicity, and phase precession (Navratilova et al., 2012). This model assumes dorsal-ventral tuning of NMDA kinetics, but observations of dorsal-ventral organization of synaptic integration (Garden et al., 2008) would likely also predict experimentally observed dorsal-ventral spacing of grid firing fields using models of this kind (Navratilova et al., 2012). Dorsal-ventral gradients in the AHP waveform are also consistent with this model (Figure 9) (Boehlen et al., 2010).

**Ion channel manipulation in behaving animals**

One of the most promising approaches for understanding the contribution of the intrinsic electrical properties of hippocampal and entorhinal neurons to spatial behavior is using animals in which key ion channels are genetically deleted. So far most attention has been directed toward the HCN1 ion channel (Nolan et al., 2004; Giocomo et al., 2011a; Hussaini et al., 2011). When the HCN1 channel is deleted from forebrain neurons spatial memory is enhanced (Nolan et al., 2004). Removal of HCN1 also increases the amplitude of the theta rhythm recorded in the hippocampus (Nolan et al., 2004). Consistent with these results and the effects of HCN1 deletion on the intrinsic properties of stellate cells (Nolan et al., 2007), grid-like firing fields of neurons in layer II of the MEC and place fields of hippocampal CA1 neurons are maintained in mice with deletion of HCN1 from forebrain neurons (Giocomo et al., 2011a; Hussaini et al., 2011). Together, these observations provide strong evidence against models for neuronal representation of space that require functions attributed to HCN1, including models based on interference between intrinsic oscillations. These results also argue against the view that HCN channels in the entorhinal cortex are pacemakers of theta-frequency activity.

Recent experiments that record from grid cells in mice with deletion of HCN1 from forebrain neurons are of further interest because they suggest that HCN1 plays important regulatory roles in the encoding of space (Giocomo et al., 2011a; Hussaini et al., 2011). Most intriguingly, grid cells and place cells observed following deletion of HCN1 appear to be larger and are more stable than in control animals (Giocomo et al., 2011a; Hussaini et al., 2011). The spacing between grid fields is also increased, while the dorsal-ventral organization of grid fields is maintained, so that on average firing fields at all locations are larger and further apart in HCN1 knockout compared to control mice (Giocomo et al., 2011a). While a potential caveat in interpretation of these observations is that deletion of HCN1 is throughout the forebrain, including neurons upstream of grid cells, they nevertheless raise important new questions. For example, why does deletion of HCN1 make grid fields larger and wider? One possibility is that this reflects the increased time-window for synaptic integration or summation of synaptic responses in stellate cells (Garden et al., 2008). Another is that it reflects place field expansion in CA1 (Hussaini et al., 2011), possibly in turn resulting from increased synaptic plasticity of CA1 pyramidal cells (Nolan et al., 2004). What mechanisms maintain the dorsal-ventral organization of grid field properties in the absence of HCN1? A possibility is that control of synaptic integration by the dorsal-ventral organization of leak potassium channels is sufficient to coordinate these properties of grid fields (Garden et al., 2008). What aspects of HCN1 channel function are critical to the behavioral roles of the channel? On the one hand, these behavioral roles may reflect control of sub-threshold resonant or oscillatory dynamics of stellate cells through voltage-dependent gating of HCN1 (Nolan et al., 2007; Giocomo and Hasselmo, 2009). On the other hand, the pattern of spike output from stellate cells may be independent of sub-threshold resonant properties (Fernandez and White, 2008), and instead the key function of HCN1 may be simply through its contribution to the resting membrane potential and conductance (Nolan et al., 2007). Further experimental investigation will be required to address these and other questions about the relationship between HCN1 channels, leak potassium channels and grid firing fields.

A further possible caveat in interpretation of behavioral results from knockout mice is that adaptation may mask roles of the deleted protein. For example, while data from knockout mice indicates that HCN1 is not required for generation of grid firing fields, this does not rule out the possibility that in wild-type animals HCN1 plays a central role that can nevertheless be compensated for when the channel is deleted. Comparison of intrinsic electrophysiological properties of stellate cells from control and HCN1 deletion mice during block of Ik does not reveal evidence of adaptation (Nolan et al., 2007). However, in pyramidal neurons from the somatosensory cortex there is strong evidence for up-regulation of a tonic GABA_A receptor conductance following deletion of HCN1 (Chen et al., 2010). This up-regulation is not found in hippocampal neurons (Nolan et al., 2004; Chen et al., 2010), and is not consistent with data from stellate neurons (Nolan et al., 2007), suggesting it may be specific to neocortical neurons. Nevertheless, these results highlight the need for additional approaches to corroborate results from knockout mice. New pharmacological tools that discriminate between HCN sub-units may be of use [e.g., (Del Lungo et al., 2011; McClure et al., 2011)], although interpretation of in vivo effects of existing HCN channel blockers is difficult because of their off-target effects (Chevaleyre and Castillo, 2002). Viruses that knockdown HCN1 in mature mice may in the future also have important experimental roles (White et al., 2011).

**Stellate cells across species**

Comparison of stellate cell properties between species may give additional insights into their roles in generation of grid firing fields. Grid fields have so far been directly observed in neurons from layer II of the MEC of rats, mice and bats (Hafting et al., 2008). Recent experiments that record from grid cells in mice with deletion of HCN1 from forebrain neurons are of further interest because they suggest that HCN1 plays important regulatory roles in the encoding of space (Giocomo et al., 2011a; Hussaini et al., 2011). Most intriguingly, grid cells and place cells observed following deletion of HCN1 appear to be larger and are more stable than in control animals (Giocomo et al., 2011a; Hussaini et al., 2011). The spacing between grid fields is also increased, while the dorsal-ventral organization of grid fields is maintained, so that on average firing fields at all locations are larger and further apart in HCN1 knockout compared to control mice (Giocomo et al., 2011a). While a potential caveat in interpretation of these observations is that deletion of HCN1 is throughout the forebrain, including neurons upstream of grid cells, they nevertheless raise important new questions. For example, why does deletion of HCN1 make grid fields larger and wider? One possibility is that this reflects the increased time-window for synaptic integration or summation of synaptic responses in stellate cells (Garden et al., 2008). Another is that it reflects place field expansion in CA1 (Hussaini et al., 2011), possibly in turn resulting from increased synaptic plasticity of CA1 pyramidal cells (Nolan et al., 2004). What mechanisms maintain the dorsal-ventral organization of grid field properties in the absence of HCN1? A possibility is that control of synaptic integration by the dorsal-ventral organization of leak potassium channels is sufficient to coordinate these properties of grid fields (Garden et al., 2008). What aspects of HCN1 channel function are critical to the behavioral roles of the channel? On the one hand, these behavioral roles may reflect control of sub-threshold resonant or oscillatory dynamics of stellate cells through voltage-dependent gating of HCN1 (Nolan et al., 2007; Giocomo and Hasselmo, 2009). On the other hand, the pattern of spike output from stellate cells may be independent of sub-threshold resonant properties (Fernandez and White, 2008), and instead the key function of HCN1 may be simply through its contribution to the resting membrane potential and conductance (Nolan et al., 2007). Further experimental investigation will be required to address these and other questions about the relationship between HCN1 channels, leak potassium channels and grid firing fields.

A further possible caveat in interpretation of behavioral results from knockout mice is that adaptation may mask roles of the deleted protein. For example, while data from knockout mice indicates that HCN1 is not required for generation of grid firing fields, this does not rule out the possibility that in wild-type animals HCN1 plays a central role that can nevertheless be compensated for when the channel is deleted. Comparison of intrinsic electrophysiological properties of stellate cells from control and HCN1 deletion mice during block of Ik does not reveal evidence of adaptation (Nolan et al., 2007). However, in pyramidal neurons from the somatosensory cortex there is strong evidence for up-regulation of a tonic GABA_A receptor conductance following deletion of HCN1 (Chen et al., 2010). This up-regulation is not found in hippocampal neurons (Nolan et al., 2004; Chen et al., 2010), and is not consistent with data from stellate neurons (Nolan et al., 2007), suggesting it may be specific to neocortical neurons. Nevertheless, these results highlight the need for additional approaches to corroborate results from knockout mice. New pharmacological tools that discriminate between HCN sub-units may be of use [e.g., (Del Lungo et al., 2011; McClure et al., 2011)], although interpretation of in vivo effects of existing HCN channel blockers is difficult because of their off-target effects (Chevaleyre and Castillo, 2002). Viruses that knockdown HCN1 in mature mice may in the future also have important experimental roles (White et al., 2011).
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tuning of synaptic integration for grid cell computation (Garden
tesults and previous cellular data suggests the importance of
these results and previous cellular data suggests the importance of
We suggest that the influence of these ion channels on integration
be a secondary consequence of tuning ion channels to control
Nevertheless, more systematic study of this issue and compari-
on species it is likely that cellular mechanisms critical for representa-
representation of space.
It demonstrates that forebrain HCN1 channels are not required
for generation or dorsal-ventral organization of grid firing fields,
but instead may be critical modulators of grid field size, spacing,
and stability (Giocomo et al., 2011a). Correspondence between
these results and previous cellular data suggests the importance of
tuning of synaptic integration for grid cell computation (Garden
ef al., 2008). It also highlights the significance of ion channels in
addition to HCN1, for example the K2P ion channels, in the topo-
graphic organization of stellate cell integrative properties (Garden
ef al., 2008). These studies point toward the exciting prospect of
a cellular and molecular understanding of the neural represen-
tation of space.

MATERIALS AND METHODS

SLICE PREPARATION AND MAINTENANCE

Sagittal brain slices with thickness of 400 µm were prepared from
30 to 53 day old mice using previously described proced-
dures (Garden et al., 2008; Dodson et al., 2011). Slices were
maintained in physiologically artificial cerebrospinal fluid of the
following composition (mM): NaCl 124, NaH2PO4 1.2, KCl 2.5,
NaHCO3 25, Glucose 20, CaCl2 2, MgCl2 1. Patch electrodes were
filled with an intracellular solution of the following composition
(mM): K Gluconate 130; KCl 10, HEPES 10, MgCl2 2, EGTA 0.1,
Na2ATP 2, Na2GTP 0.3, NaPhosphoCreatine 10. Excitatory gluta-
mate receptor-mediated and inhibitory GABA receptor-mediated
synaptic transmission was blocked with NBQX (5 µM), D-APV
(50 µM), picrotoxin (50 µM) and CGP55845 (1 µM).

DATA COLLECTION

All recordings were made at 34–37°C. Recording electrodes had a
tip resistance of 3–7 MΩ. All membrane seals had resistance of
>2 GΩ as calculated by the current response to 10 mV pulses
in voltage-clamp. A liquid junction potential of 8.1 mV (bath
relative to pipette) was not corrected for. Stellate cells were iden-
tified in the superficial part of layer II by their large soma,
the presence of multiple similar sized primary dendrites and
their characteristic electrophysiological properties including
the presence of prominent sag potentials and bi-phasic spike after-
hyperpolarization (AHP).

For the perforated patch recordings the electrodes were front-
filled with standard intracellular solution and then back-filled
with an intracellular solution containing the antibiotic ampho-
terin B (final concentration 0.075–0.15 mg/ml) dissolved in
DMSO (final concentration 3–6 µl/ml). For these recordings,
after rapidly establishing a cell-attached configuration, rather
than rupturing the membrane to achieve a whole-cell configura-
tion, we waited for the antibiotic to permeabilize the membrane.
We confirmed that spontaneous break-in did not occur by includ-
ing a fluorescent label (Alexa 488) in the intracellular solu-
tion. Using this technique, access resistances of 24–85 MΩ were
achieved after a delay lasting from 20 min to an hour after seal
formation. Whole-cell patch recordings had access resistances
<30 MΩ. All access resistances were fully compensated for using
the bridge-balance technique. All signals were recorded in current
clamp using a Multiclamp 700B amplifier (Molecular Devices),
low pass filtered at 10 kHz (two pole Bessel filter), converted with
an ADC (National Instruments ITC-18) and sampled at 20 kHz.

DIAA ANALYSIS

Data analysis was performed as previously described (Nolan et al.,
2007; Garden et al., 2008; Dodson et al., 2011) using Axograph
and custom written MATLAB routines. Statistical calculations
were done using R (www.r-project.org). All results are reported
as mean ± standard error.

Time frequency analysis

To examine the persistence of spectral components across time we
used Lomb analysis to identify significant frequencies in contigu-
ous 3 segments perithreshold recordings as previously described
(Dodson et al., 2011). Spectrograms (in Figure 4) had windows
of 1 s or 6.5 s, with 32768 or 131072 points and had 90% or

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95% overlap. The first second of each 20 s trace was ignored to eliminate high frequency transients associated with current injection steps.

**Spike clustering**

We used a previously described algorithm (Nolan et al., 2007) to calculate the spike clustering coefficient, \( P_c \), the proportion of spikes occurring in a cluster. Briefly, a group of spikes was considered a cluster if no inter-spike interval in the group exceeded 250 ms and if the inter-cluster interval was at least 300 ms.

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Conflict of Interest Statement: The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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