A nano-patterned photonic crystal laser with a dye-doped liquid crystal


Cavendish Laboratory, University of Cambridge, JJ Thomson Avenue, Cambridge CB3 0HE, United Kingdom
Department of Engineering, University of Cambridge, 9 JJ Thomson Avenue, Cambridge CB3 0FA, United Kingdom

(Received 10 April 2013; accepted 4 June 2013; published online 29 July 2013)

Covering a nano-patterned titanium dioxide photonic crystal (PC) within a well-oriented film of dye-doped liquid crystal (LC), a distributed feedback laser is constructed whereby the emission characteristics can be manipulated in situ using an electric field. This hybrid organic-inorganic structure permits simultaneous selectivity of both the beam pattern and laser wavelength by electrical addressing of the LC director. In addition, laser emission is obtained both in the plane and normal to the PC. Along with experimental data, a theoretical model is presented that is based upon an approximate calculation of the band structure of this birefringent, tuneable laser device. © 2013 AIP Publishing LLC. [http://dx.doi.org/10.1063/1.4816243]

Laser devices formed from two dimensional photonic crystals (PCs), based upon both inorganic1–5 and organic6–8 materials, have been fabricated demonstrating a range of remarkable emission properties.9–11 Control of the laser emission characteristics can be achieved at the fabrication stage by pre-selecting the periodicity and the geometrical arrangement of the lattice. However, many potential applications of PC lasers would benefit from in situ control of the emission properties. Tuneability can be achieved by infiltrating liquid crystals (LCs) into PCs so as to vary the spectral position of the photonic band gap through a range of external stimuli.12–23 However, it is often difficult to align the LC uniformly in the zero-field state resulting in limited tuning capabilities.20,21

In this letter, we report on a compact laser that consists of a nano-patterned structure. A thin, nano-patterned titanium dioxide (TiO2) layer is covered with a nematic, dye-doped LC. Uniform alignment of the LC director is provided by gentle rubbing of the TiO2-layer and by covering the device with a glass plate coated with both polyimide (PI) and indium tin oxide (ITO). With the rubbing directions of the PI and the TiO2 layers being parallel, a uniform director field parallel to the x-y-plane (Figure 1) is created. The device shows distributed feedback lasing that can be characterized by the band structure of the thin PC layer, even though the PC layer possesses a finite thickness (expansion in z-direction) even smaller than the laser wavelength.24 Additionally, the emission characteristics of the device can be manipulated in situ by using a quasi static electric field that can be applied across the device. Above a threshold of \( \approx 0.5 \text{ V/\mu m} \), the electric field induces realignment of the LC director parallel to the electric field lines.25 By addressing the LC with an electric field, both the beam pattern and laser wavelength can be controlled simultaneously.

The TiO2 nano-pattern was fabricated by embossing a TiO2 sol-gel with a perfluropolyether mould at 150°C for 2 h under pressure.26,27 After peeling off the mould, the TiO2 film was then annealed at 450°C for 2 h to increase the refractive index of the nano-patterned film. Measurement of the refractive index of flat TiO2 layers fabricated with this method yielded values of \( n_r \approx 1.8 \) before and \( n_r \approx 2.0 \) after annealing, respectively, at a wavelength of 589 nm. As the active medium, a nematic LC25 (E49, Merck KGaA) doped with 1 wt. % laser dye (4-(Dicyanomethylene)-2-methyl-6-(4-dimethylaminostyryl)-4H-pyran, Exciton) was used. The refractive indices of this mixture were found to be \( n_r = 1.77 \) and \( n_c = 1.53 \). As shown by scanning electron microscope images (Figure 1), the PC consisted of a hexagonal array of pillars with a protrusion of 130 nm, a pitch of \( a = 400 \text{ nm} \), and a radius of \( r = 0.80 \text{ nm} = 0.2a \).

The LC alignment was investigated with polarized optical microscopy. Although the feature size of the nano-patterned TiO2 layer is beyond the optical resolution, a non-uniform alignment at the TiO2-LC interface would lead to intensity variations in the microscope field of view. In contrast, a defect free, uniform field of view was observed. Thus, a uniform alignment of the confined nematic LC can be

---

**FIG. 1.** SEM images of the PC layer and schematic of the TiO2 PC/LC laser device. The SEM images show the cross-section and top view of the PC layer. All scale bars are 500 nm. Transparent electrodes on both the top and bottom substrates enabled an electrical field to be applied perpendicular to the plane of the PC.

---

\[ \text{Center for Opto-Electronic Convergence Systems, Korea Institute of Science and Technology, Seoul, Republic of Korea.} \]
\[ \text{E-mail: smm56@cam.ac.uk} \]
\[ \text{E-mail: rhf10@cam.ac.uk} \]
assumed. The thickness of the LC layer was defined to be 10 μm using spacer beads, which was chosen to ensure a large dye reservoir above the nano-patterned layer that could be replenished through thermal diffusion.

To confirm that distributed feedback of the thin PC-layer was responsible for the laser emission rather than leaky modes from within either the dye-doped nematic layer or random amplification as a result of multiple scattering, a device with a non-patterned TiO2 layer was also studied. The TiO2-LC samples were excited normal to the glass substrates by the second harmonic of an Nd:YAG laser (Polaris II, New Wave Research) at an excitation rate of 1 Hz and pulse width of 3–4 ns. The pump beam was focused to a spot of 100 μm. Emission was collected along the normal to the substrates for the 1st-order Bragg mode (surface emission) by a high numerical aperture objective and focused into an optical fiber that was connected to a broadband spectrometer with a resolution of 0.3 nm. In the device with the non-patterned TiO2 layer, fluorescence but not lasing was observed (Figure 2(a)). For all excitation energies used, the threshold for amplified spontaneous emission was found to be greater than 10 000 nJ/pulse (130 mJ/cm2). On the contrary, a thin PC layer provides distributed feedback at certain locations in the band structure where three dimensional lasing can be expected at specific points of the Brillouin-zone boundary corresponding to where the bands cross and split. In the nano-patterned device, laser emission at a wavelength of 617 nm with a line-width of 0.6 nm was observed (Figure 2(b)).

Incorporating the dye-doped LC film leads to the device becoming sensitive to the plane of polarization of the input beam. A half-wave plate was used to control the orientation of the plane of linear polarization relative to the LC director. Results for the input-output characteristics for two orthogonal polarizations are presented in Figure 3, which shows a notable change in the excitation threshold from $E_{th} = 89 \pm 5$ nJ/pulse (1.1 mJ/cm2) to $E_{th} = 299 \pm 15$ nJ/pulse (3.8 mJ/cm2). The large difference in the threshold is considered to be due to the change in the absorption of the dye through the dye-guest host effect; the orientation of the transition dipole moment of the dye aligns preferentially with the director of the LC and therefore the absorption is maximized when the input polarization was parallel to the director.

FIG. 2. Comparison of non-patterned and nano-patterned TiO2-LC devices. (a) Fluorescence from the non-patterned TiO2 device when photo-excited along a direction normal to the plane of the device. (b) Laser emission from the nano-patterned TiO2-LC device. Δλ represents the full-width at half maximum. The cell gap in both cases was 10 μm. The insets of the figures represent the two different device structures. Note that in (b) the height of the pillars is not to scale and is simply to indicate the nano-patterned array.

FIG. 3. Emission characteristics of the nano-patterned TiO2-LC laser device. (a) The output intensity as a function of excitation energy for the two orthogonal linear polarization states of the TiO2-LC laser when photo-excited normal to the plane of the device: parallel to the LC director (blue squares) and perpendicular to the director (orange triangles). (b) Output as a function of excitation energy for a chiral nematic LC laser with the same film thickness and dye. (c) Laser emission in the no-field state (black line) and with an applied electric field (5 kHz) with a field strength of 12 V/μm (red line). (d) Laser emission in three dimensions consisting of in-plane (x, y) and surface (z) emission.
The lowest recorded laser threshold for this device compares favourably with values that have been reported recently for other 2D surface-emitting PC lasers. For comparison, the input-output curve of a conventional dye-doped chiral nematic LC that has a 1D periodic structure is shown in Figure 3(b), where the periodicity is perpendicular to the plane of the device. The 1D chiral nematic laser consisted of the same laser dye and nematic host as that of the TiO$_2$-LC laser. Furthermore, to ensure optimum conditions, the laser wavelength was matched to the gain maximum of the dye. Additionally, to ensure optimum conditions, the laser wavelength was matched to the gain maximum of the dye. For all factors being equal, there is almost a 3-fold decrease in the power at the output of the laser.

Wavelength tuning of the TiO$_2$-LC laser device is shown in Figure 3(c) where it is observed that the emission wavelength can be blue-shifted by 12 nm and was found to be fully reversible. Upon removal of the applied electric field, the laser emission red-shifted back to the initial wavelength of 617 nm. The recovery of the initial wavelength was found to occur within the timescale set by the repetition rate of the pump laser (i.e., 1 s). The shift that is recorded here is larger than those observed previously in LC-infiltrated PC laser devices. At the boundaries to the glass cover plate and the nematic TiO$_2$, a finite coherence length of the elastic deformation of the LC may still be present. However, this coherence length decreases with the electric field strength. Therefore, a uniform director field where the LC director is parallel to the z-direction can be assumed to be in the field-on state, which was confirmed using polarizing optical microscopy. In addition to the laser emission in the z-direction, laser emission was also detected in the x-, y-direction (Figure 3(d)), as expected. The present experiments indicate that addressing of the LC with an electric field causes a shift of the laser wavelength of the three-dimensional emission.

There is also a simultaneous change in the far-field beam profile from an elliptical shape to a circular profile when an electric field is applied, Figure 4. The far-field pattern is determined by the Fourier transform of the internal electromagnetic field distribution, which can be controlled by changing the refractive index profile of the medium between the pillars by adjusting the alignment of the LC. The change from an asymmetric to symmetric beam shape reflects the variation from an anisotropic refractive index profile within the plane of the pillars to a single refractive index value when the electric field is applied. Adjusting the polarization of the pump beam does not alter the far-field emission profile for both cases. The control of the beam profile using this configuration is an important development as it is generally only achievable through the appropriate design of the unit cell of the PC.

The band structure of layered structures with a central, thin PC layer can be simulated in a two-stepped approach where first, the fundamental mode of the layered structure is simulated and the effective refractive index $n_{\text{eff}}$ and a confinement factor $f_c$ in the PC layer are then extracted. Subsequently, the band diagram of a two-dimensional PC with infinite expansion in the z-direction, but with unchanged two-dimensional geometry ($p = 400 \text{ nm}$, $r = 0.2p$) is calculated. To account for the finite thickness of the PC-layer, effective material constants are considered in this second step. An effective average dielectric constant $\tilde{\varepsilon}$ is approximated as

$$\tilde{\varepsilon} = \frac{\varepsilon_a + \varepsilon_b}{2}. \quad (1)$$

The effective dielectric contrast $\Delta \varepsilon$ and the effective dielectric constants $\varepsilon_a$ (effective background material) and $\varepsilon_b$ (in z-direction infinitely expanded cylinders) are calculated by using $f_c$ and the dielectric constants of the background material $\varepsilon_1$ and of the pillars $\varepsilon_2$

$$\Delta \varepsilon = \varepsilon_b - \varepsilon_a = f_c \cdot (\varepsilon_2 - \varepsilon_1). \quad (2)$$

First, a modal analysis of a cross section of the structure along the $\Gamma$-J-direction was conducted by using a full vectorial finite element method-based algorithm (COMSOL

![Image](https://example.com/image.png)

**FIG. 4.** Modification of the far-field beam profile using electric fields. No electric field: (a) illustration of the LC alignment, (b) 2D plot of the beam profile, and (c) 3D plot of the beam profile. (d), (e), and (f) are equivalent figures when the sample is subjected to an electric field (12 V/µm at 5 kHz). These results were recorded on a Spiricon Beam Profiler.
The results are listed in Table I. As expected, the calculated value \( \Delta \varepsilon_x \) (x-polarized mode) is smaller than the calculated value \( \Delta \varepsilon_z \) (z-polarized mode). An effective dielectric constant of the infinitely extended cylinders (\( \varepsilon_b = 3.616 \)) was approximated from \( \varepsilon \) obtained for the x-polarized mode because it is a close approximation of the fundamental mode.

### TABLE I. Polarization and dielectric properties of two simulated modes of a two dimensional cross section of the layered structure with nano-patterned TiO\(_2\) layer covered by a birefringent LC.

<table>
<thead>
<tr>
<th>Director field</th>
<th>Polarization</th>
<th>( \varepsilon_1 )</th>
<th>( \varepsilon_2 )</th>
<th>( \varepsilon )</th>
<th>( \Delta \varepsilon )</th>
</tr>
</thead>
<tbody>
<tr>
<td>Uniaxial x</td>
<td>X</td>
<td>( n_x^2 )</td>
<td>2.0(^2)</td>
<td>3.415</td>
<td>0.403</td>
</tr>
<tr>
<td>Uniaxial y</td>
<td>Z</td>
<td>( n_y^2 )</td>
<td>2.0(^2)</td>
<td>3.152</td>
<td>0.581</td>
</tr>
<tr>
<td>Uniaxial z</td>
<td>X</td>
<td>( n_z^2 )</td>
<td>2.0(^2)</td>
<td>3.301</td>
<td>0.516</td>
</tr>
<tr>
<td>Uniaxial z</td>
<td>Z</td>
<td>( n_z^2 )</td>
<td>2.0(^2)</td>
<td>3.171</td>
<td>0.407</td>
</tr>
</tbody>
</table>

The band diagram was calculated with MIT photonic bands (MPB),\(^{30}\) where a diagonal dielectric tensor (\( \varepsilon_{xx} = \varepsilon_{yy} = \varepsilon_{zz} = \varepsilon_b - \Delta \varepsilon \)) was considered to model the birefringent background material (uniaxial with optic axis along the \( \Gamma-J \)-direction). The calculated band diagram is shown (Fig. 5(c), light colored symbols). A point of the Brillouin-zone boundary where six bands couple is indicated by an arrow. Laser emission perpendicular to the layers of the structure can be expected at this point.\(^{24}\) Accordingly, the simulation results in six possible laser modes at wavelengths of \( \approx 618 \pm 8 \) nm, which is in good agreement with the experiments, where lasing at a wavelength of 617 nm was observed. The simulation was repeated by considering a birefringent medium with optic axis parallel to the z-direction to model the case that the LC director was uniformly aligned with the application of an electric field. The simulation resulted in modes with varying confinement factors (20.6\% for the x-polarized mode and 23.9\% for the z-polarized mode). From the dielectric data (Table I), values were obtained for use in the diagonal dielectric tensor (\( \varepsilon_{xx} = \varepsilon_{yy} = \varepsilon_{zz} = \varepsilon_b - \Delta \varepsilon \)) that was considered for the background material in the band diagram calculation. The result is shown (Figure 5(e), dark colored symbols). Additionally, one mode with linear polarization of the electric field is shown [Fig. 5(e), inset]. This mode is blue shifted from 620 nm (first simulation) to 607 nm (second simulation). A blue-shift ranging from 4 to 14 nm in magnitude was predicted for all simulated modes, respectively. The presented model cannot predict the exact behaviour of the real device, since it is only two-dimensional\(^{24}\) and alignment effects of the LC cannot be considered. Nevertheless, it is in good qualitative agreement with the experiments.

In summary, these results demonstrate a laser device that has the capacity to “tune on demand” a number of emission properties besides the laser wavelength. This is in contrast to the tuning method of passive systems which involve a “tuning by design” approach to change factors such as the lattice constant, filling fraction, and symmetry of the PC. The presented results are encouraging for the fabrication of distributed feedback lasers using thin PC layers and other LC phases, which may result in alternative beam profiles providing a greater variety. Further work on the role of the device thickness on the performance of the laser would also be of interest.

The authors gratefully acknowledge the financial support of the Engineering and Physical Sciences Research Council (UK) through the COSMOS Technology Translation award (EP/H046658/1). One of the authors (S.M.M.) acknowledges The Royal Society for financial support and one of the authors (A.L.) acknowledges financial support by the German Research Council (DFG, LO 1922/1-1). Use of NCN/nanohub.org computational resources [MIT Photonic Bands, Ref. 30] is gratefully acknowledged.

---